

QUANTITATIVE ESTIMATES ON THE BINDING ENERGY FOR HYDROGEN IN NON-RELATIVISTIC QED

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ABSTRACT. In this paper, we determine the exact expression for the hydrogen binding energy in the Pauli-Fierz model up to the order $O(\alpha^5 \log \alpha^{-1})$, where α denotes the finestructure constant, and prove rigorous bounds on the remainder term of the order $o(\alpha^5 \log \alpha^{-1})$. As a consequence, we prove that the binding energy is not a real analytic function of α , and verify the existence of logarithmic corrections to the expansion of the ground state energy in powers of α , as conjectured in the recent literature.

1. INTRODUCTION

For a hydrogen-like atom consisting of an electron interacting with a static nucleus of charge eZ , described by the Schrödinger Hamiltonian $-\Delta - \frac{\alpha Z}{|x|}$,

$$\inf \sigma(-\Delta) - \inf \sigma(-\Delta - \frac{\alpha Z}{|x|}) = \frac{(Z\alpha)^2}{4}.$$

corresponds to the binding energy necessary to remove the electron to spatial infinity.

The interaction of the electron with the quantized electromagnetic field is accounted for by adding to $-\Delta - \frac{\alpha Z}{|x|}$ the photon field energy operator H_f , and an operator $I(\alpha)$ which describes the coupling of the electron to the quantized electromagnetic field; the small parameter α is the fine structure constant. Thereby, one obtains the Pauli-Fierz Hamiltonian described in detail in Section 2. The systematic study of this operator, in a more general case involving more than one electron, was initiated by Bach, Fröhlich and Sigal [5, 3, 4].

In this case, determining the binding energy

$$(1) \quad \inf \sigma(-\Delta + H_f + I(\alpha)) - \inf \sigma(-\Delta - \frac{\alpha Z}{|x|} + H_f + I(\alpha))$$

is a very hard problem. A main obstacle emerges from the fact that the ground state energy is not an isolated eigenvalue of the Hamiltonian, and can not be determined with ordinary perturbation theory. Furthermore, the photon form factor in the quantized electromagnetic vector potential occurring in the interaction term $I(\alpha)$ contains a critical frequency space singularity that is responsible for the infamous infrared problem in quantum electrodynamics. As a consequence, quantities such as the ground state energy do not a priori exist as a convergent power series in the fine structure constant α with coefficients independent of α .

In recent years, several rigorous results addressing the computation of the binding energy have been obtained, [16, 15, 10]. In particular, the coupling to the photon field has been shown to increase the binding energy of the electron to the nucleus, and that up to normal ordering, the leading term is $\frac{(\alpha Z)^2}{4}$ [14, 16, 10].

Moreover, for a model with scalar bosons, the binding energy is determined in [15], in the first subleading order in powers of α , up to α^3 , with error term $\alpha^3 \log \alpha^{-1}$. This result has inspired the question of a possible emergence of logarithmic terms in the expansion of the binding energy; however, this question has so far remained open.

In [2], a sophisticated rigorous renormalization group analysis is developed in order to determine the ground state energy (and the renormalized electron mass) up to any arbitrary precision in powers of α , with an expansion of the form

$$(2) \quad \varepsilon_0 + \sum_{k=1}^{2N} \varepsilon_k(\alpha) \alpha^{k/2} + o(\alpha^N) ,$$

(for any given N) where the coefficients $\varepsilon_k(\alpha)$ diverge as $\alpha \rightarrow 0$, but are smaller in magnitude than any positive power of α^{-1} . The recursive algorithms developed in [2] are highly complex, and explicitly computing the ground state energy to any subleading order in powers of α is an extensive task. While it is expected that the rate of divergence of these coefficient functions is proportional to a power of $\log \alpha^{-1}$, this is not explicitly exhibited in the current literature; for instance, it can a priori not be ruled out that terms involving logarithmic corrections cancel mutually.

The choice of atomic units in [2], inherited from [3], and motivated by the necessity to keep the α -dependence only in the interaction term but not in the Coulomb term, introduces a dependence between the ultraviolet cutoff and the fine structure constant. This makes it challenging to compare the result derived for (2) to our estimate of the binding energy stated in Theorem 2.1, since we employ different units with an ultraviolet cutoff independent of α . However, such comparison would require to apply the procedure used in this paper to the model described in [2].

The goal of the current paper is to develop an alternative method (as a continuation of [7]) that determines the binding energy up to several subleading orders in powers of α , with rigorous error bounds, and proving the presence of terms logarithmic in α .

The main result established in the present paper (for $Z = 1$) states that the binding energy can be estimated as

$$(3) \quad \frac{\alpha^2}{4} + e^{(1)}\alpha^3 + e^{(2)}\alpha^4 + e^{(3)}\alpha^5 \log \alpha^{-1} + o(\alpha^5 \log \alpha^{-1}) ,$$

where $e^{(i)}$ ($i = 1, 2, 3$) are independent of α , $e^{(1)} > 0$, and $e^{(3)} \neq 0$. Their explicit values are given in Theorem 2.1.

As a consequence, we conclude that the binding energy is *not analytic* in α . In addition, our proof clarifies how the logarithmic factor in $\alpha^5 \log \alpha^{-1}$ is linked to the infrared singularity of the photon form factor in the interaction term $I(\alpha)$. We note that for some models with a mild infrared behavior, [12], the ground state energy is proven to be analytic in α .

Notice that on the basis of the present computations, it is impossible to determine whether the $\alpha^5 \log \alpha^{-1}$ term comes from $\inf \sigma(-\Delta + H_f + I(\alpha))$, or $\inf \sigma(-\Delta - \frac{\alpha Z}{|x|} + H_f + I(\alpha))$ or both. However, this question has been recently answered in [6].

Organization of the proof. The strategy consists mainly in an iteration procedure based on variational estimates. The derivation of the main result (3) is accomplished in two main steps, with first an estimate up to the order α^3 with an

error of the order $\mathcal{O}(\alpha^4)$, and then up to the order $\alpha^5 \log \alpha$. Each of these steps requires both an upper and a lower bound on the binding energy. From the knowledge of an approximate ground state for $-\Delta + H_f + I(\alpha)$ (see [7]), the remaining hard task consists mainly in establishing a lower bound for the ground state energy for H yielding an upper bound for the binding energy. A lower bound for the binding energy can be obtained by choosing a bona fide trial state.

It is easy to see that the first term in the expansion for the binding energy is not smaller than the Coulomb energy $\frac{\alpha^2}{4}$ by taking a trial state, which is a product state of the electronic ground state u_α (15) of the Schrödinger-Coulomb operator $-\Delta - \frac{\alpha}{|x|}$ and the photon ground state of the self-energy operator given by (7) with zero total momentum. Indeed, as shown in [10] there is an increase of the binding energy. However the first term is exactly the Coulomb energy [15]

The proof of the upper bound, up to the order α^3 and to the order $\alpha^5 \log \alpha^{-1}$ reduces to a minimization process.

Because of the soft photon problem, the set of minimizing sequences does not have a clear structure which helps to find the minimizer, or approximate minimizers. Therefore, prior to the above two steps, we have to appropriately restrict the set of states which we are looking at. For that purpose, we first establish, in Lemma 3.1, an a priori bound on the expected photon number in the ground state Ψ^{GS} of the Pauli-Fierz operator $(-\Delta - \frac{\alpha}{|x|} + H_f + I(\alpha))$, namely $\langle \Psi^{\text{GS}}, N_f \Psi^{\text{GS}} \rangle = \mathcal{O}(\alpha^2 \log \alpha^{-1})$, where N_f is the photon number operator. Moreover, in Theorem 4.1, we estimate the contribution to the binding energy stemming from states orthogonal to the ground state u_α of the Schrödinger operator. It enables us to show that the projection of the ground state Ψ^{GS} onto the subspace of functions orthogonal to u_α is small in several norms.

Estimate up to the order α^3 : To find a lower bound for the binding energy (Theorem 5.1), we pick a trial function Φ^{trial} which is the sum of two states. The first state is the product of the ground state u_α of the Schrödinger-Coulomb operator with the approximate ground state Ψ_0 of the self-energy operator (given by (7) below) with zero total momentum which was derived in [7]. The second state $\Phi_\#^{u_\alpha}$ (42) is a term orthogonal to u_α and has been chosen to minimize, up to the order α^3 , the sum of the cross term $\langle (-\Delta - \frac{\alpha}{|x|} + H_f + I(\alpha))u_\alpha \Psi_0, \Phi_\#^{u_\alpha} \rangle$ and the term $\langle (H_f + P_f^2 - \Delta - \frac{\alpha}{|x|} + \frac{\alpha^2}{4})\Phi_\#^{u_\alpha}, \Phi_\#^{u_\alpha} \rangle$ stemming from the quadratic form for H on the trial state $\Phi^{\text{trial}} = u_\alpha \Psi_0 + \Phi_\#^{u_\alpha}$.

At that stage, we emphasize that the contribution of these two states yields not only an α^3 term, but also an $\alpha^5 \log \alpha^{-1}$ term, which is in fact the only contribution to the binding energy for this order as the detailed proof of the next order estimate shall show.

To recover the upper bound, we take an arbitrary state Ψ satisfying the condition on the expected photon number. Then we consider $u_\alpha \Psi^{u_\alpha}$, the projection of this state onto u_α , and Ψ^\perp the projection onto the subspace orthogonal to u_α . We estimate in Lemma 7.1 the quadratic form of H on $u_\alpha \Psi^{u_\alpha}$, in Lemma 7.4 the quadratic form for H on Ψ^\perp , and in Lemmata 7.2-7.3 the cross term. Collecting this estimates in (82) below yields that the function Ψ can be an approximate minimizer up to the order α^3 , with error term $\mathcal{O}(\alpha^4 \log \alpha^{-1})$ of the functional $\langle H\Psi, \Psi \rangle$ only if the difference R between this function Ψ and the previous trial state Φ^{trial} satisfies the condition $\|H_f^{\frac{1}{2}} R\|^2 = \mathcal{O}(\alpha^4 \log \alpha^{-1})$ (see (87))

According to Lemma 7.5, the last estimate itself allows to improve the expected photon number estimate for the ground state of H given by Lemma 3.1, replacing there $\alpha^2 \log \alpha^{-1}$ by α^2 . Then, repeating the above steps for the upper bound with this improved expected photon number yields the upper bound of the binding energy with an error $\mathcal{O}(\alpha^4)$ instead of $\mathcal{O}(\alpha^4 \log \alpha^{-1})$.

Estimate up to the order $\alpha^5 \log \alpha^{-1}$: In order to derive the lower bound for the binding energy, we improve the previous trial function by adding several terms that were irrelevant for the estimate up to the order α^3 . Due to the improved estimates in Lemma 7.5, the expected photon number for the difference of the true ground state and this new trial function can not exceed $\mathcal{O}(\alpha^{\frac{33}{16}})$. Furthermore, we infer from Theorem 5.2 that this difference satisfies several smallness conditions. Using these conditions we estimate once more the quadratic form and minimize it with respect to this difference in a similar way to what is done in section 5 for the previous step.

The paper is organized as follows: In section 2, we give a detailed definition of the model and state the main result. In section 3 we establish an a priori bound on the expected photon number in the ground state of the Pauli-Fierz operator. In section 4 we estimate the contribution to the binding energy stemming from states orthogonal to the ground state of the Schrödinger operator. The section 5 is devoted to the statements of the lower and the upper bounds for the binding energy up to the order α^3 , as well as to the proof of the lower bound. The difficult part of the proof, namely the upper bound on the binding energy, is presented in four parts and postponed to section 7. In subsections 7.1-7.3, we estimate separately the terms according to the splitting of the variational state. We collect these results in subsection 7.4 and establish then the proof of the upper bound up to the order α^3 . In section 6, we establish the main steps of the proof of the upper bound for the binding energy up to the order $\alpha^5 \log \alpha^{-1}$. We start with some useful definitions, the statement for the upper bound, and two propositions (proved in appendices subsections A and B), that gives estimates of the contributions to the binding energy according to a refined splitting of the variational state. Details of the remaining straightforward computations are given in section 8. Finally, subsection 9 provide the proof for a lower bound on the binding energy. In appendix C, we provide some technical lemmata whose proof are straightforward but rather long.

2. THE MODEL

We study a scalar electron interacting with the quantized electromagnetic field in the Coulomb gauge, and with the electrostatic potential generated by a fixed nucleus. The Hilbert space accounting for the Schrödinger electron is given by $\mathfrak{H}_{el} = L^2(\mathbb{R}^3)$. The Fock space of photon states is given by

$$\mathfrak{F} = \bigoplus_{n \in \mathbb{N}} \mathfrak{F}_n,$$

where the n -photon space $\mathfrak{F}_n = \bigotimes_s^n (L^2(\mathbb{R}^3) \otimes \mathbb{C}^2)$ is the symmetric tensor product of n copies of one-photon Hilbert spaces $L^2(\mathbb{R}^3) \otimes \mathbb{C}^2$. The factor \mathbb{C}^2 accounts for the two independent transversal polarizations of the photon. On \mathfrak{F} , we introduce creation- and annihilation operators $a_\lambda^*(k)$, $a_\lambda(k)$ satisfying the distributional commutation relations

$$[a_\lambda(k), a_{\lambda'}^*(k')] = \delta_{\lambda, \lambda'} \delta(k - k') \quad , \quad [a_\lambda^\sharp(k), a_{\lambda'}^\sharp(k')] = 0,$$

where a_λ^\sharp denotes either a_λ or a_λ^* . There exists a unique unit ray $\Omega_f \in \mathfrak{F}$, the Fock vacuum, which satisfies $a_\lambda(k)\Omega_f = 0$ for all $k \in \mathbb{R}^3$ and $\lambda \in \{1, 2\}$.

The Hilbert space of states of the system consisting of both the electron and the radiation field is given by

$$\mathfrak{H} := \mathfrak{H}_{el} \otimes \mathfrak{F}.$$

We use units such that $\hbar = c = 1$, and where the mass of the electron equals $m = 1/2$. The electron charge is then given by $e = \sqrt{\alpha}$, where the fine structure constant α will here be considered as a small parameter.

Let $x \in \mathbb{R}^3$ be the position vector of the electron and let $y_i \in \mathbb{R}^3$ be the position vector of the i -th photon.

We consider the normal ordered Pauli-Fierz Hamiltonian on \mathfrak{H} for Hydrogen,

$$(4) \quad : (i\nabla_x \otimes I_f - \sqrt{\alpha}A(x))^2 : + V(x) \otimes I_f + I_{el} \otimes H_f.$$

where $: \dots :$ denotes normal ordering, corresponding to the subtraction of a normal ordering constant $c_{n.o.}\alpha$, with $c_{n.o.}I_f := [A^-(x), A^+(x)]$ is independent of x .

The electrostatic potential $V(x)$ is the Coulomb potential for a static point nucleus of charge $e = \sqrt{\alpha}$ (i.e., $Z = 1$)

$$V(x) = -\frac{\alpha}{|x|}.$$

We will describe the quantized electromagnetic field by use of the Coulomb gauge condition.

The operator that couples an electron to the quantized vector potential is given by

$$A(x) = \sum_{\lambda=1,2} \int_{\mathbb{R}^3} \frac{\chi_\Lambda(|k|)}{2\pi|k|^{1/2}} \varepsilon_\lambda(k) \left[e^{ikx} \otimes a_\lambda(k) + e^{-ikx} \otimes a_\lambda^*(k) \right] \mathbf{k}$$

where by the Coulomb gauge condition, $\text{div} A = 0$.

The vectors $\varepsilon_\lambda(k) \in \mathbb{R}^3$ are the two orthonormal polarization vectors perpendicular to k ,

$$\varepsilon_1(k) = \frac{(k_2, -k_1, 0)}{\sqrt{k_1^2 + k_2^2}} \quad \text{and} \quad \varepsilon_2(k) = \frac{k}{|k|} \wedge \varepsilon_1(k).$$

The function χ_Λ implements an *ultraviolet cutoff* on the momentum k . We assume χ_Λ to be of class C^1 , with compact support in $\{|k| \leq \Lambda\}$, $\chi_\Lambda \leq 1$ and $\chi_\Lambda = 1$ for $|k| \leq \Lambda - 1$. For convenience, we shall write

$$A(x) = A^-(x) + A^+(x),$$

where

$$A^-(x) = \sum_{\lambda=1,2} \int_{\mathbb{R}^3} \frac{\chi_\Lambda(|k|)}{2\pi|k|^{1/2}} \varepsilon_\lambda(k) e^{ikx} \otimes a_\lambda(k) \, dk$$

is the part of $A(x)$ containing the annihilation operators, and $A^+(x) = (A^-(x))^*$.

The photon field energy operator H_f is given by

$$H_f = \sum_{\lambda=1,2} \int_{\mathbb{R}^3} |k| a_\lambda^*(k) a_\lambda(k) \, dk.$$

We will, with exception of our discussion in Section 3, study the unitarily equivalent Hamiltonian

$$(5) \quad H = U \left(: (i\nabla_x \otimes I_f - \sqrt{\alpha}A(x))^2 : + V(x) \otimes I_f + I_{el} \otimes H_f \right) U^*,$$

where the unitary transform U is defined by

$$U = e^{iP_f \cdot x},$$

and

$$P_f = \sum_{\lambda=1,2} \int k a_{\lambda}^*(k) a_{\lambda}(k) \mathbb{k}$$

is the photon momentum operator. We have

$$U i \nabla_x U^* = i \nabla_x + P_f \quad \text{and} \quad U A(x) U^* = A(0) .$$

In addition, the Coulomb operator V , the photon field energy H_f , and the photon momentum P_f remain unchanged under the action of U . Therefore, in this new system of variables, the Hamiltonian reads as follows

$$(6) \quad H = : ((i \nabla_x \otimes I_f - I_{el} \otimes P_f) - \sqrt{\alpha} A(0))^2 : + H_f - \frac{\alpha}{|x|} ,$$

where $: \dots :$ denotes again the normal ordering. Notice that the operator H can be rewritten, taking into account the normal ordering and omitting, by abuse of notations, the operators I_{el} and I_f ,

$$\begin{aligned} H = & (-\Delta_x - \frac{\alpha}{|x|}) + (H_f + P_f^2) - 2 \operatorname{Re} (i \nabla_x \cdot P_f) \\ & - 2 \sqrt{\alpha} (i \nabla_x - P_f) \cdot A(0) + 2 \alpha A^+(0) \cdot A^-(0) + 2 \alpha \operatorname{Re} (A^-(0))^2 . \end{aligned}$$

For a free spinless electron coupled to the quantized electromagnetic field, the self-energy operator T is given by

$$T = : (i \nabla_x \otimes I_f - \sqrt{\alpha} A(x))^2 : + I_{el} \otimes H_f .$$

We note that this system is translationally invariant; that is, T commutes with the operator of total momentum

$$P_{tot} = p_{el} \otimes I_f + I_{el} \otimes P_f ,$$

where p_{el} and P_f denote respectively the electron and the photon momentum operators.

Let $\mathbb{C} \otimes \mathfrak{F}$ be the fiber Hilbert space corresponding to conserved total momentum $p \in \mathbb{R}^3$.

For fixed value p of the total momentum, the restriction of T to the fibre space is given by (see e.g. [8])

$$(7) \quad T(p) = : (p - P_f - \sqrt{\alpha} A(0))^2 : + H_f$$

where by abuse of notation, we again dropped all tensor products involving the identity operators I_f and I_{el} . Henceforth, we will write

$$A^{\pm} \equiv A^{\pm}(0) .$$

Moreover, we denote

$$\Sigma_0 = \inf \sigma(T) \quad \text{and} \quad \Sigma = \inf \sigma(H) = \inf \sigma(T + V) .$$

It is proven in [1, 8] that $\Sigma_0 = \inf \sigma(T(0))$ is an eigenvalue of the operator $T(0)$.

Our main result is the following theorem.

Theorem 2.1. *The binding energy fulfills*

$$(8) \quad \Sigma_0 - \Sigma = \frac{1}{4}\alpha^2 + e^{(1)}\alpha^3 + e^{(2)}\alpha^4 + e^{(3)}\alpha^5 \log \alpha^{-1} + o(\alpha^5 \log \alpha^{-1}),$$

where

$$e^{(1)} = \frac{2}{\pi} \int_0^\infty \frac{\chi_\Lambda^2(t)}{1+t} dt,$$

$$\begin{aligned} e^{(2)} &= \frac{2}{3} \operatorname{Re} \sum_{i=1}^3 \langle (A^-)^i (H_f + P_f^2)^{-1} A^+ . A^+ \Omega_f, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle \\ &+ \frac{1}{3} \sum_{i=1}^3 \| (H_f + P_f^2)^{-\frac{1}{2}} \left(2A^+ . P_f (H_f + P_f^2)^{-1} (A^+)^i - P_f^i (H_f + P_f^2)^{-1} A^+ . A^+ \right) \Omega_f \|^2 \\ &- \frac{2}{3} \sum_{i=1}^3 \| A^- (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \|^2 + 4a_0^2 \| Q_1^\perp (-\Delta - \frac{1}{|x|} + \frac{1}{4})^{-\frac{1}{2}} \Delta u_1 \|^2, \end{aligned}$$

$$a_0 = \int \frac{k_1^2 + k_2^2}{4\pi^2 |k|^3} \frac{2}{|k|^2 + |k|} \chi_\Lambda(|k|) k_1 k_2 k_3,$$

$$e^{(3)} = -\frac{1}{3\pi} \| (-\Delta - \frac{1}{|x|} + \frac{1}{4})^{\frac{1}{2}} \nabla u_1 \|^2,$$

and Q_1^\perp is the projection onto the orthogonal complement to the ground state u_1 of the Schrödinger operator $-\Delta - \frac{1}{|x|}$ (for $\alpha = 1$).

3. BOUNDS ON THE EXPECTED PHOTON NUMBER

Lemma 3.1. *Let*

$$K = (i\nabla - \sqrt{\alpha}A(x))^2 + H_f - \frac{\alpha}{|x|} = v^2 + H_f - \frac{\alpha}{|x|},$$

be the Pauli-Fierz operator defined without normal ordering, where $v = i\nabla - \sqrt{\alpha}A(x)$. Let $\Psi \in \mathcal{H}$ denote the ground state of K ,

$$K \Psi = E \Psi,$$

normalized by

$$\|\Psi\| = 1.$$

Let

$$N_f := \sum_{\lambda=1,2} \int a_\lambda^*(k) a_\lambda(k) k$$

denote the photon number operator. Then, there exists a constant c independent of α , such that for any sufficiently small $\alpha > 0$, the estimate

$$\langle \Psi, N_f \Psi \rangle \leq c \alpha^2 \log \alpha^{-1}$$

is satisfied.

Proof. Using

$$[a_\lambda(k), H_f] = |k|, \quad \text{and} \quad [a_\lambda(k), v] = \frac{\epsilon(k)}{2\pi |k|^{\frac{1}{2}}} \chi_\lambda(k) e^{ik \cdot x}$$

and the pull-through formula,

$$\begin{aligned} a_\lambda(k)E\Psi &= a_\lambda(k)K\Psi \\ &= \left[(H_f + |k|)a_\lambda(k) - \frac{1}{|x|}a_\lambda(k) + [v, a_\lambda(k)]v + v[v, a_\lambda(k)] \right] \Psi, \end{aligned}$$

we get

$$(9) \quad a_\lambda(k) \Psi = -\frac{\sqrt{\alpha}\chi_\Lambda(|k|)}{2\pi\sqrt{|k|}} \frac{2}{K+|k|-E} \left((i\nabla - \sqrt{\alpha}A(x)) \cdot \epsilon_\lambda(k) e^{ik \cdot x} \right) \Psi.$$

From (9), we obtain

$$\begin{aligned} \|a_\lambda(k)\Psi\|^2 &\leq \frac{\alpha\chi_\Lambda(|k|)}{\pi^2|k|} \left\| \frac{1}{K+|k|-E} \right\|^2 \left\| (i\nabla - \sqrt{\alpha}A(x))\Psi \right\|^2 \\ &\leq \frac{\alpha\chi_\Lambda(|k|)}{\pi^2|k|^3} \left[\langle \Psi, K\Psi \rangle + \langle \Psi, \frac{\alpha}{|x|}\Psi \rangle \right], \end{aligned}$$

since $K - E \geq 0$, and $(i\nabla - \sqrt{\alpha}A(x))^2 \leq (K + \frac{\alpha}{|x|})$.

Since $\langle \Psi, :K:\Psi \rangle = -\alpha^2/4 + o(\alpha^2)$ ([14, 16, 10]) and $:K:= K - c_{\text{n.o.}}\alpha$, we have

$$(10) \quad \langle \Psi, K\Psi \rangle \leq c_{\text{n.o.}}\alpha.$$

Moreover, for the normal ordered hamiltonian defined in (4), we have

$$\begin{aligned} :K: &= -\Delta_x - 4\sqrt{\alpha}\text{Re}(i\nabla_x \cdot A^-(x)) + \alpha :A(x)^2: + H_f - \frac{\alpha}{|x|} \\ (11) \quad &\geq (1 - c\sqrt{\alpha})(-\Delta_x) + \alpha :A(x)^2: + (1 - c\sqrt{\alpha})H_f - \frac{\alpha}{|x|} \\ &\geq -c\alpha^2 - \frac{1}{2}\Delta_x - \frac{\alpha}{|x|} = -c\alpha^2 + 2(-\frac{1}{4}\Delta_x - \frac{\alpha}{|x|}) + \frac{\alpha}{|x|}. \end{aligned}$$

where in the last inequality we used (see e.g. [13])

$$\alpha :A(x)^2: + (1 - c\sqrt{\alpha})H_f \geq -c\alpha^2.$$

Since

$$-\frac{1}{4}\Delta_x - \frac{\alpha}{|x|} \geq -4\alpha^2,$$

inequality (11) implies

$$(12) \quad \langle \Psi, \frac{\alpha}{|x|}\Psi \rangle \leq \langle -\Delta_x \Psi, \Psi \rangle + \frac{1}{4}\alpha^2 \|\Psi\|^2 \leq c\alpha^2 \|\Psi\|^2.$$

Collecting (10), (12) and (9) we find

$$(13) \quad \|a_\lambda(k)\Psi\| \leq \frac{c\alpha\chi_\Lambda(|k|)}{|k|^{\frac{3}{2}}}.$$

This a priori bound exhibits the L^2 -critical singularity in frequency space. It does not take into consideration the exponential localization of the ground state due to the confining Coulomb potential, and appears in a similar form for the free electron.

To account for the latter, we use the following two results from the work of Griesemer, Lieb, and Loss, [13]. Equation (58) in [13] provides the bound

$$\|a_\lambda(k)\Psi\| < \frac{c\sqrt{\alpha}\chi_\Lambda(|k|)}{|k|^{\frac{1}{2}}} \| |x|\Psi \|.$$

Moreover, Lemma 6.2 in [13] states that

$$\left\| \exp[\beta|x|]\Psi \right\|^2 \leq c \left[1 + \frac{1}{\Sigma_0 - E - \beta^2} \right] \|\Psi\|^2 ,$$

for any

$$\beta^2 < \Sigma_0 - E = O(\alpha^2) .$$

For the 1-electron case, Σ_0 is the infimum of the self-energy operator, and E is the ground state energy of K . Choosing $\beta = O(\alpha)$,

$$\begin{aligned} \| |x|\Psi \| &\leq \| |x|^4 \Psi \|^{1/4} \|\Psi\|^{3/4} \leq \frac{(4!)^{1/4}}{\beta} \left\| \exp[\beta|x|]\Psi \right\|^{1/4} \|\Psi\|^{3/4} \\ &\leq \frac{c}{\beta} \left[1 + \frac{1}{\Sigma_0 - E - \beta^2} \right]^{1/8} \|\Psi\| \\ &\leq c_1 \alpha^{-5/4} . \end{aligned}$$

Notably, this bound only depends on the binding energy of the potential.

Thus,

$$(14) \quad \left\| a_\lambda(k)\Psi \right\| < \frac{c\alpha^{-3/4}\chi_\Lambda(|k|)}{|k|^{1/2}} .$$

We see that binding to the Coulomb potential weakens the infrared singularity by a factor $|k|$, but at the expense of a large constant factor α^{-2} . For the free electron, this estimate does not exist.

Using (13) and (14), we find

$$\begin{aligned} \langle \Psi , N_f \Psi \rangle &= \int \mathbf{k} \left\| a_\lambda(k)\Psi \right\|^2 \\ &\leq \int_{|k|<\delta} \mathbf{k} \frac{c\alpha^{-3/2}}{|k|} + \int_{\delta \leq |k| \leq \Lambda} \mathbf{k} \frac{c\alpha^2}{|k|^3} \\ &\leq c\alpha^{-3/2}\delta^2 + c'\alpha^2 \log \frac{1}{\delta} \\ &\leq c\alpha^{3/4} + c''\alpha^2 \log \alpha^{-1} . \end{aligned}$$

for $\delta = \alpha^{15/8}$. This proves the lemma. \square

4. ESTIMATES ON THE QUADRATIC FORM FOR STATES ORTHOGONAL TO THE GROUND STATE OF THE SCHRÖDINGER OPERATOR

Throughout this paper, we will denote by π^n the projection onto the n -th photon sector (without distinction for the n -photon sector of \mathfrak{F} and the n -photon sector of \mathfrak{H}). We also define $\pi^{\geq n} = 1 - \sum_{j=0}^{n-1} \pi^j$.

Starting with this section, we study the Hamiltonian H defined in (6), written in relative coordinates. In particular, $i\nabla_x$ now stands for the operator unitarily equivalent to the operator of total momentum, which, by abuse of notation, will be denoted by P .

Let

$$(15) \quad u_\alpha(x) = \frac{1}{\sqrt{8\pi}} \alpha^{3/2} e^{-\alpha|x|/2}$$

be the normalized ground state of the Schrödinger operator

$$h_\alpha := -\Delta_x - \frac{\alpha}{|x|}.$$

We will also denote by $-e_0 = -\frac{\alpha^2}{4}$ and $-e_1 = -\frac{\alpha^2}{16}$ the two lowest eigenvalues of this operator.

Theorem 4.1. *Assume that $\Phi \in \mathfrak{H}$ fulfils $\langle \pi^k \Phi, u_\alpha \rangle_{L^2(\mathbb{R}^3, dx)} = 0$, for all $k \geq 0$. Then there exists $1 > \nu > 0$ and $\alpha_0 > 0$ such that for all $0 < \alpha < \alpha_0$*

$$(16) \quad \langle H\Phi, \Phi \rangle \geq (\Sigma_0 - e_0)\|\Phi\|^2 + \delta\|\Phi\|^2 + \nu\|H_f^{\frac{1}{2}}\Phi\|^2,$$

where $\delta = (e_0 - e_1)/2 = \frac{3}{32}\alpha^2$.

Remark 4.1. *All photons with momenta larger than the ultraviolet cutoff do not contribute to lower the energy. More precisely, due to the cutoff function $\chi_\Lambda(|k|)$ in the definition of $A(x)$, and since we have*

$$H = (i\nabla_x - P_f)^2 - 2\sqrt{\alpha}\text{Re}(i\nabla_x - P_f).A(0) + \alpha : A(0)^2 : + H_f - \frac{\alpha}{|x|},$$

it follows that for any given normalized state $\Phi \in \mathfrak{H}$, there exists a normalized state $\Phi_{\leq \Lambda}$ such that $\forall x \in \mathbb{R}^3$, for all $n \in \{1, 2, \dots\}$, for all $((k_1, \lambda_1), (k_2, \lambda_2), \dots, (k_n, \lambda_n)) \in ((\mathbb{R}^3 \setminus \{k, |k| \leq \Lambda + 1\}) \times \{1, 2\})^n$, we have

$$(17) \quad \pi^n \Phi_{\leq \Lambda}(x, (k_1, \lambda_1), (k_2, \lambda_2), \dots, (k_n, \lambda_n)) = 0$$

and

$$\langle \Phi_{\leq \Lambda}, H\Phi_{\leq \Lambda} \rangle \leq \langle \Phi, H\Phi \rangle \quad \text{and} \quad \langle \Phi_{\leq \Lambda}, T\Phi_{\leq \Lambda} \rangle \leq \langle \Phi, T\Phi \rangle.$$

A key consequence of this remark is that throughout the paper, all states will be implicitly assumed to fulfill condition (17). This is crucial for the proof of Corollary 4.2.

To prove Theorem 4.1, we first need the following Lemma.

Lemma 4.1. *There exists $c_0 > 0$ such that for all α small enough we have*

$$H - \frac{1}{2}(P - P_f)^2 - \frac{1}{2}H_f \geq -c_0\alpha^2.$$

A straightforward consequence of this lemma is the following result.

Corollary 4.1. *Let Ψ^{GS} be the normalized ground state of H . Then*

$$(18) \quad \langle H_f \Psi^{GS}, \Psi^{GS} \rangle \leq 2c_0\alpha^2 \|\Psi^{GS}\|^2$$

Proof. This follows from $\langle H\Psi^{GS}, \Psi^{GS} \rangle \leq (\Sigma_0 - e_0)\|\Psi^{GS}\|^2 < 0$. The last inequality holds since $e_0 = -\alpha^2/4$ and Σ_0 is the infimum for the normal ordered Hamiltonian for the free electron, and thus $\Sigma_0 \leq 0$. \square

Moreover, from Theorem 4.1 and Lemma 4.1, we obtain

Corollary 4.2. *Assume that $\Phi \in \mathfrak{H}$ is such that $\langle \pi^n \Phi, u_\alpha \rangle_{L^2(\mathbb{R}^3, dx)} = 0$ holds for all $n \geq 0$. Then, for ν and δ defined in Theorem 4.1, there exists $\zeta > 0$, and $\alpha_0 > 0$ such that for all $0 < \alpha < \alpha_0$ we have*

$$(19) \quad \langle H\Phi, \Phi \rangle \geq (\Sigma_0 - e_0)\|\Phi\|^2 + M[\Phi],$$

where

$$(20) \quad M[\Phi] := \frac{\delta}{2}\|\Phi\|^2 + \frac{\nu}{2}\|H_f^{\frac{1}{2}}\Phi\|^2 + \zeta\|(P - P_f)\Phi\|^2 + \zeta\|\pi^{n \leq 4}P\Phi\|^2.$$

Proof. According to Remark 4.1, there exists $\tilde{c} > 1$ such that the operator inequality $P_f^2 \pi^{n \leq 4} \leq \tilde{c} H_f \pi^{n \leq 4}$ holds on the set of states for which (17) is satisfied. The value of \tilde{c} only depends on the ultraviolet cutoff Λ . Thus,

$$\|\pi^{n \leq 4} P \Phi\|^2 \leq 2\|(P - P_f)\Phi\|^2 + 2\|\pi^{n \leq 4} P_f \Phi\|^2 \leq 2\|(P - P_f)g\|^2 + 2\tilde{c}\|H_f^{\frac{1}{2}} \pi^{n \leq 4} \Phi\|^2,$$

which yields

$$\|(P - P_f)\Phi\|^2 \geq \frac{1}{2}\|\pi^{n \leq 4} P \Phi\|^2 - \tilde{c}\|H_f^{\frac{1}{2}} \pi^{n \leq 4} \Phi\|^2.$$

Therefore, it suffices to prove Corollary 4.2 with $M[\Phi]$ replaced by

$$(21) \quad \widetilde{M}[\Phi] := \frac{\delta}{2}\|\Phi\|^2 + \frac{3}{4}\nu\|H_f^{\frac{1}{2}}\Phi\|^2 + 2\zeta\|(P - P_f)\Phi\|^2,$$

and ζ small enough so that $\tilde{c}\zeta < \frac{\nu}{4}$.

Now we consider two cases. Let $c_1 := \max\{8c_0, 8\delta/\alpha^2\}$.

If $\|(P - P_f)\Phi\|^2 \leq c_1\alpha^2\|\Phi\|^2$, Theorem 4.1 and the above remark imply (19).

If $\|(P - P_f)\Phi\|^2 > c_1\alpha^2\|\Phi\|^2$, Lemma 4.1 implies that

$$\begin{aligned} \langle H\Phi, \Phi \rangle &\geq \frac{1}{2}\langle (P - P_f)^2 \Phi, \Phi \rangle + \frac{1}{2}\langle H_f \Phi, \Phi \rangle - c_0\alpha^2\|\Phi\|^2 \\ &\geq \frac{1}{4}\langle (P - P_f)^2 \Phi, \Phi \rangle + \frac{1}{2}\langle H_f \Phi, \Phi \rangle + \frac{1}{8}c_1\alpha^2\|\Phi\|^2, \end{aligned}$$

which concludes the proof since $\Sigma_0 - e_0 < 0$. \square

Proof of Lemma 4.1. Recall the notation $A^\pm \equiv A^\pm(0)$. The following holds.

$$\begin{aligned} H - \frac{1}{2}(P - P_f)^2 - \frac{1}{2}H_f \\ = \frac{1}{2}(P - P_f)^2 - \frac{\alpha}{|x|} - 2\sqrt{\alpha}\text{Re}((P - P_f).A(0)) + 2\alpha\text{Re}(A^-)^2 + 2\alpha A^+.A^- + \frac{1}{2}H_f, \end{aligned}$$

$$(22) \quad \frac{1}{4}(P - P_f)^2 - \frac{\alpha}{|x|} \geq -4\alpha^2.$$

and

$$(23) \quad 2\sqrt{\alpha}|\langle (P - P_f).A(0)\psi, \psi \rangle| \leq 2\sqrt{\alpha}\|(P - P_f)\psi\|^2 + 2\sqrt{\alpha}\|A^-\psi\|^2.$$

By the Schwarz inequality, there exists c_1 independent of α such that

$$(24) \quad \|A^-\psi\|^2 \leq c_1\|H_f^{\frac{1}{2}}\psi\|^2.$$

Inequalities (23)-(24) imply that for small α ,

$$(25) \quad 2\sqrt{\alpha}|\langle (P - P_f).A(0)\psi, \psi \rangle| \leq \frac{1}{4}\|(P - P_f)\psi\|^2 + \frac{1}{4}\langle H_f \psi, \psi \rangle.$$

Moreover using (24) and

$$(26) \quad \|A^+\psi\|^2 \leq c_2\|\psi\|^2 + c_3\|H_f^{\frac{1}{2}}\psi\|^2,$$

we arrive at

$$\begin{aligned} \alpha\langle (A^-)^2 \psi, \psi \rangle &= \alpha\langle A^-\psi, A^+\psi \rangle \leq \epsilon\|A^-\psi\|^2 + \epsilon^{-1}\alpha^2\|A^+\psi\|^2 \\ (27) \quad &\leq \epsilon c_1\|H_f^{\frac{1}{2}}\psi\|^2 + \epsilon^{-1}\alpha^2(c_2\|\psi\|^2 + c_3\|H_f^{\frac{1}{2}}\psi\|^2). \end{aligned}$$

Collecting the inequalities (22), (25) and (27) with $\epsilon < 1/(8c_1)$ and α small enough, completes the proof. \square

Proof of Theorem 4.1: Let $\Phi := \Phi_1 + \Phi_2 := \chi(|P| < \frac{p_c}{2})\Phi + \chi(|P| \geq \frac{p_c}{2})\Phi$, where $P = i\nabla_x$ is the total momentum operator (due to the transformation (5)) and $p_c = \frac{1}{3}$ is a lower bound on the norm of the total momentum for which [8, Theorem 3.2] holds.

Since P commutes with the translation invariant operator $H + \frac{\alpha}{|x|}$, we have for all $\epsilon \in (0, 1)$,

$$(28) \quad \begin{aligned} \langle H\Phi, \Phi \rangle &= \langle H\Phi_1, \Phi_1 \rangle + \langle H\Phi_2, \Phi_2 \rangle - 2\operatorname{Re} \langle \frac{\alpha}{|x|} \Phi_1, \Phi_2 \rangle \\ &\geq \langle H\Phi_1, \Phi_1 \rangle + \langle H\Phi_2, \Phi_2 \rangle - \epsilon \langle \frac{\alpha}{|x|} \Phi_1, \Phi_1 \rangle - \epsilon^{-1} \langle \frac{\alpha}{|x|} \Phi_2, \Phi_2 \rangle. \end{aligned}$$

• First, we have the following estimate

$$\begin{aligned} &: (P - P_f - \sqrt{\alpha}A(0))^2 : + H_f \\ &= (P - P_f)^2 - 2\operatorname{Re} (P - P_f) \cdot \sqrt{\alpha}A(0) + \alpha : A(0)^2 : + H_f \\ &\geq (1 - \sqrt{\alpha})(P - P_f)^2 + (\alpha - \sqrt{\alpha}) : A(0)^2 : + H_f - c_{\text{n.o.}} \sqrt{\alpha} \\ &\geq (1 - \sqrt{\alpha})(P - P_f)^2 + (1 - \mathcal{O}(\sqrt{\alpha}))H_f - \mathcal{O}(\sqrt{\alpha}) \end{aligned}$$

where in the last inequality we used (24) and (26). Therefore

$$(29) \quad \begin{aligned} \langle (H - \epsilon^{-1} \frac{\alpha}{|x|}) \Phi_2, \Phi_2 \rangle &\geq \langle (\frac{1 - \sqrt{\alpha}}{2} (P - P_f)^2 - (1 + \epsilon^{-1}) \frac{\alpha}{|x|}) \Phi_2, \Phi_2 \rangle \\ &+ \langle \left(\frac{1 - \sqrt{\alpha}}{2} (P - P_f)^2 + (1 - \mathcal{O}(\sqrt{\alpha}))H_f - \mathcal{O}(\sqrt{\alpha}) \right) \Phi_2, \Phi_2 \rangle \end{aligned}$$

The lowest eigenvalue of the Schrödinger operator $-(1 - \mathcal{O}(\sqrt{\alpha}))\frac{\Delta}{2} - \frac{(1 + \epsilon^{-1})\alpha}{|x|}$ is larger than $-c_\epsilon \alpha^2$. Thus, using (29) and denoting

$$L := \frac{1 - \sqrt{\alpha}}{2} (P - P_f)^2 + (1 - \mathcal{O}(\sqrt{\alpha}))H_f - \mathcal{O}(\sqrt{\alpha}) - c_\epsilon \alpha^2,$$

we get

$$(30) \quad \langle H\Phi_2, \Phi_2 \rangle - \epsilon^{-1} \langle \frac{\alpha}{|x|} \Phi_2, \Phi_2 \rangle \geq \langle L\Phi_2, \Phi_2 \rangle.$$

Now we have the following alternative: Either $|P_f| < \frac{p_c}{3}$, in which case we have $\langle L\Phi_2, \Phi_2 \rangle \geq (\frac{p_c^2}{24} - \mathcal{O}(\sqrt{\alpha}))\|\Phi_2\|^2$, or $|P_f| \geq \frac{p_c}{3}$, in which case, using $\Phi_2 = \chi(|P| > \frac{p_c}{2})\Phi_2$ and $H_f \geq |P_f|$, we have $L \geq (\frac{p_c}{6} - \mathcal{O}(\sqrt{\alpha}))\|\Phi_2\|^2$. In both cases, for α small enough, this yields the bound

$$(31) \quad \langle L\Phi_2, \Phi_2 \rangle \geq \frac{p_c^2}{48} \|\Phi_2\|^2 \geq (\Sigma_0 - e_0 + \frac{7}{8}(e_0 - e_1)) \|\Phi_2\|^2$$

since, for α small enough, the right hand side tends to zero, whereas p_c is a constant independent of α . Inequalities (30) and (31) yield

$$(32) \quad \langle H\Phi_2, \Phi_2 \rangle - \epsilon^{-1} \langle \frac{\alpha}{|x|} \Phi_2, \Phi_2 \rangle \geq (\Sigma_0 - e_0 + \frac{7}{8}(e_0 - e_1)) \|\Phi_2\|^2.$$

• For $T(p)$ being the self-energy operator with fixed total momentum $p \in \mathbb{R}^3$ defined in (7), we have from [8, Theorem 3.1 (B)]

$$|\inf \sigma(T(p)) - p^2 - \inf \sigma(T(0))| \leq c_0 \alpha p^2.$$

Therefore

$$(33) \quad T(p) \geq (1 - o_\alpha(1))p^2 + \Sigma_0.$$

Case 1: If $\|\Phi_2\|^2 \geq 8\|\Phi_1\|^2$, we first do the following estimate, using (33),

$$(34) \quad \begin{aligned} & \langle H\Phi_1, \Phi_1 \rangle - \epsilon \langle \frac{\alpha}{|x|} \Phi_1, \Phi_1 \rangle \\ & \geq (1 - o_\alpha(1))(P^2\Phi_1, \Phi_1) - \langle (1 + \epsilon) \frac{\alpha}{|x|} \Phi_1, \Phi_1 \rangle + \Sigma_0 \|\Phi_1\|^2 \\ & \geq (\Sigma_0 - (1 + \mathcal{O}(\alpha) + \mathcal{O}(\epsilon))e_0) \|\Phi_1\|^2. \end{aligned}$$

Therefore, together with $\|\Phi_2\|^2 \geq 8\|\Phi_1\|^2$ and (32), for α and ϵ small enough this implies

$$(35) \quad \langle H\Phi, \Phi \rangle \geq (\Sigma_0 - e_0) \|\Phi\|^2 + \frac{3}{4}(e_0 - e_1) \|\Phi\|^2.$$

Case 2: If $\|\Phi_2\|^2 < 8\|\Phi_1\|^2$, we write the estimate

$$(36) \quad \begin{aligned} & \langle H\Phi_1, \Phi_1 \rangle - \epsilon \langle \frac{\alpha}{|x|} \Phi_1, \Phi_1 \rangle \\ & \geq (1 - o_\alpha(1))(P^2\Phi_1, \Phi_1) - \langle (1 + \epsilon) \frac{\alpha}{|x|} \Phi_1, \Phi_1 \rangle + \Sigma_0 \|\Phi_1\|^2 \\ & \geq (1 + o_\alpha(1) + \mathcal{O}(\epsilon)) \left(-e_0 \sum_{k=0}^{\infty} \|\langle \pi^k \Phi_1, u_\alpha \rangle_{L^2(\mathbb{R}^3, dx)}\|^2 - e_1 (\|\Phi_1\|^2 \right. \\ & \quad \left. - \sum_{k=0}^{\infty} \|\langle \pi^k \Phi_1, u_\alpha \rangle_{L^2(\mathbb{R}^3, dx)}\|^2) \right) + \Sigma_0 \|\Phi_1\|^2. \end{aligned}$$

Now, by orthogonality of Φ and u_α in the sense that for all k , $\langle \pi^k \Phi, u_\alpha \rangle_{L^2(\mathbb{R}^3, dx)} = 0$, we get

$$(37) \quad \begin{aligned} & \sum_{k=0}^{\infty} \|\langle u_\alpha, \pi^k \Phi_1 \rangle_{L^2(\mathbb{R}^3, dx)}\|^2 = \sum_{k=0}^{\infty} \|\langle u_\alpha, \pi^k \Phi_2 \rangle_{L^2(\mathbb{R}^3, dx)}\|^2 \\ & \leq \|\Phi_2\|^2 \|\chi(|P| \geq \frac{p_c}{2}) u_\alpha\|^2 \leq 8\|\Phi_1\|^2 \|\chi(|P| \geq \frac{p_c}{2}) u_\alpha\|^2 \rightarrow_{\alpha \rightarrow 0} 0 \end{aligned}$$

Thus, for α and ϵ small enough, (36), (37) and (32) imply also (35) in that case.

• Let $\tilde{c} = \max\{\delta, |c_0|\alpha^2\}$.

If $\langle H_f \Phi, \Phi \rangle < 8\tilde{c} \|\Phi\|^2$, (16) follows from (35) with $\nu = \delta/(16\tilde{c})$.

If $\langle H_f \Phi, \Phi \rangle \geq 8\tilde{c} \|\Phi\|^2$, using Lemma 4.1, we obtain

$$(38) \quad \begin{aligned} \langle H\Phi, \Phi \rangle & \geq \frac{1}{2} \langle H_f \Phi, \Phi \rangle - c_0 \alpha^2 \|\Phi\|^2 \geq \frac{1}{4} \langle H_f \Phi, \Phi \rangle + \tilde{c} \|\Phi\|^2 \\ & \geq \frac{1}{4} \langle H_f \Phi, \Phi \rangle + \delta \|\Phi\|^2 + (\Sigma_0 - e_0) \|\Phi\|^2, \end{aligned}$$

since $\Sigma_0 - e_0 \leq 0$, which yields (16) with $\nu = \frac{1}{4}$.

This concludes the proof of (16).

5. ESTIMATE OF THE BINDING ENERGY UP TO α^3 TERM

Definition 5.1. Let u_α be the normalized ground state of the Schrödinger operator h_α , as defined in (15).

We define the projection $\Psi^{u_\alpha} \in \mathfrak{F}$ of the normalized ground state Ψ^{GS} of H , onto u_α as follows

$$\Psi^{GS} = u_\alpha \Psi^{u_\alpha} + \Psi^\perp,$$

where for all $k \geq 0$,

$$(39) \quad \langle u_\alpha, \pi^k \Psi^\perp \rangle_{L^2(\mathbb{R}^3, dx)} = 0.$$

Remark 5.1. The definition implies that for all m

$$(\pi^m \Psi^{u_\alpha})(k_1, \lambda_1; k_2, \lambda_2; \dots; k_m, \lambda_m) = \int_{\mathbb{R}^3} (\pi^m \Psi^{GS})(y; k_1, \lambda_1; \dots; k_m, \lambda_m) \overline{u_\alpha(y)} y.$$

Definition 5.2. Let

$$\begin{aligned} \Phi_*^2 &:= -(H_f + P_f^2)^{-1} A^+ \cdot A^+ \Omega_f \\ \Phi_*^3 &:= -(H_f + P_f^2)^{-1} P_f \cdot A^+ \Phi_*^2 \\ \Phi_*^1 &:= -(H_f + P_f^2)^{-1} P_f \cdot A^- \Phi_*^2 \end{aligned}$$

where evidently, the state Φ_*^i contains i photons.

Definition 5.3. On \mathfrak{F} , we define the positive bilinear form

$$(40) \quad \langle v, w \rangle_* := \langle v, (H_f + P_f^2) w \rangle,$$

and its associated semi-norm $\|v\|_* = \langle v, v \rangle_*^{1/2}$.

We will also use the same notation for this bilinear forms on \mathfrak{F}_n , \mathfrak{H} and \mathfrak{H}_n .

Similarly, we define the bilinear form $\langle \cdot, \cdot \rangle_\#$ on \mathfrak{H} as

$$\langle u, v \rangle_\# := \langle u, (H_f + P_f^2 + h_\alpha + e_0) v \rangle$$

and its associated semi-norm $\|v\|_\# = \langle v, v \rangle_\#^{1/2}$.

Definition 5.4. Let

$$(41) \quad \Phi_*^{u_\alpha} := 2\alpha^{\frac{1}{2}} \nabla u_\alpha \cdot (H_f + P_f^2)^{-1} A^+ \Omega_f,$$

and

$$(42) \quad \Phi_\#^{u_\alpha} := 2\alpha^{\frac{1}{2}} (H_f + P_f^2 + h_\alpha + e_0)^{-1} A^+ \cdot \nabla u_\alpha \Omega_f.$$

Remark 5.2. The function $\Phi_*^{u_\alpha}$ is not a vector in the Hilbert space \mathfrak{H} because of the infrared singularity of the photon form factor. However, in the rest of the paper, we only used the vectors $H_f^\gamma \Phi_*^{u_\alpha}$ or $P_f^\gamma \Phi_*^{u_\alpha}$, with some $\gamma > 0$, which are always well defined. In particular, all expressions involving $(H_f + P_f^2)^{-1}$ are always well-defined in the sequel.

The next theorem gives an upper bound on the binding energy up to the term α^3 with an error term $\mathcal{O}(\alpha^4)$.

Theorem 5.1 (Lower bound on the binding energy). *We have*

$$(43) \quad \Sigma \leq \Sigma_0 - e_0 - \|\Phi_\#^{u_\alpha}\|_\#^2 + \mathcal{O}(\alpha^4)$$

Proof. Using the trial function in \mathfrak{H}

$$\Phi^{\text{trial}} := u_\alpha (\Omega_f + 2\alpha^{\frac{3}{2}} \Phi_*^1 + \alpha \Phi_*^2 + 2\alpha^{\frac{3}{2}} \Phi_*^3) + \Phi_\#^{u_\alpha},$$

and from [7, Theorem 3.1] which states

$$\Sigma_0 = -\alpha^2 \|\Phi_*^2\|_*^2 + \alpha^3 (2\|A^- \Phi_*^2\|^2 - 4\|\Phi_*^1\|^2 - 4\|\Phi_*^3\|^2) + \mathcal{O}(\alpha^4),$$

the result follows straightforwardly. \square

We decompose the function Ψ^{u_α} defined in Definition 5.1 as follows.

Definition 5.5. Let η_1, η_2, η_3 and $\Delta_*^{u_\alpha}$ be defined by

$$\Psi^{u_\alpha} =: \pi^0 \Psi^{u_\alpha} + 2\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \eta_2 \alpha \Phi_*^2 + 2\eta_3 \alpha^{\frac{3}{2}} \Phi_*^3 + \Delta_*^{u_\alpha},$$

with the conditions $\pi^0 \Delta_*^{u_\alpha} = 0$ and $\langle \Phi_*^i, \pi^i \Delta_*^{u_\alpha} \rangle_* = 0$ ($i=1,2,3$), where $\Phi_*^1, \Phi_*^2, \Phi_*^3$ are given in Definition 5.2.

We further decompose Ψ^\perp into two parts.

Definition 5.6. Let

$$\Psi^\perp =: \kappa_1 \Phi_\#^{u_\alpha} + \Delta_\#^\perp,$$

be defined by $\langle \Phi_\#^{u_\alpha}, \pi^1 \Delta_\#^\perp \rangle_\# = 0$.

The following theorem provides an upper bound of the binding energy with an error term of the order $\mathcal{O}(\alpha^4)$. Together with Theorem 5.1, it establishes an estimate up to the order α^3 with an error term $\mathcal{O}(\alpha^4)$.

Theorem 5.2 (Upper bound on the binding energy). 1) Let $\Sigma = \inf \sigma(H)$. Then

$$(44) \quad \Sigma \geq \Sigma_0 - e_0 - \|\Phi_*^{u_\alpha}\|^2 + \mathcal{O}(\alpha^4),$$

$$\Sigma_0 = -\alpha^2 \|\Phi_*^2\|^2 + \alpha^3 (2\|A^- \Phi_*^2\|^2 - 4\|\Phi_*^1\|^2 - 4\|\Phi_*^3\|^2) + \mathcal{O}(\alpha^4),$$

and $\Phi_*^{u_\alpha}$ defined by (41).

2) For the components $\Delta_\#^\perp, \Psi^{u_\alpha}, \Delta_*^{u_\alpha}$ of the ground state Ψ^{GS} , and the coefficients η_1, η_2, η_3 and κ_1 defined in Definitions 5.1-5.5, holds

$$(45) \quad \|\Delta_\#^\perp\|^2 = \mathcal{O}(\alpha^{\frac{33}{16}}), \quad \|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 = \mathcal{O}(\alpha^4), \quad \|(P - P_f) \Delta_\#^\perp\|^2 = \mathcal{O}(\alpha^4),$$

$$(46) \quad |\pi^0 \Psi^{u_\alpha}|^2 \geq 1 - c\alpha^2,$$

$$(47) \quad \|\Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^4), \quad \|\Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^{\frac{33}{16}}),$$

$$(48) \quad |\eta_{1,3} - 1|^2 \leq c\alpha, \quad |\eta_2 - 1|^2 \leq c\alpha^2, \quad |\kappa_1 - 1|^2 \leq c\alpha.$$

To prove this theorem, we will compute $\langle H \Psi^{GS}, \Psi^{GS} \rangle$ according to the decomposition of Ψ^{GS} introduced in Definitions 5.1 to 5.6. Using

$$H = (P^2 - \frac{\alpha}{|x|}) + : (P_f + \sqrt{\alpha} A(0))^2 : + H_f - 2\text{Re } P \cdot (P_f + \sqrt{\alpha} A(0)),$$

and due to the orthogonality of u_α and Ψ^\perp , we obtain

$$(49) \quad \langle H \Psi^{GS}, \Psi^{GS} \rangle = \langle H u_\alpha \Psi^{u_\alpha}, u_\alpha \Psi^{u_\alpha} \rangle + \langle H \Psi^\perp, \Psi^\perp \rangle - 4\text{Re} \langle P \cdot (P_f + \sqrt{\alpha} A(0)) u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle.$$

We will estimate separately each term in (49) in subsections 7.1-7.3. These estimates will be used to establish in subsection 7.4 the proof of Theorem 5.2.

6. ESTIMATE OF THE BINDING ENERGY UP TO $o(\alpha^5 \log \alpha^{-1})$ TERM

We develop here the proof of the difficult part in Theorem 2.1 which is the upper bound in (8), and which is stated in Theorem 6.1 below for convenience.

Some technical aspects of this proof are detailed in section 8 and appendices A and B.

Theorem 6.1 (Upper bound up to the order $\alpha^5 \log \alpha^{-1}$ for the binding energy).
For α small enough, we have

$$(50) \quad \Sigma_0 - \Sigma \geq \frac{1}{4}\alpha^2 + e^{(1)}\alpha^3 + e^{(2)}\alpha^4 + e^{(3)}\alpha^5 \log \alpha^{-1} + o(\alpha^5 \log \alpha^{-1}),$$

where $e^{(1)}$, $e^{(2)}$ and $e^{(3)}$ are defined in Theorem 2.1.

In order to prove this result, we need to refine the splitting of the function Ψ^\perp orthogonal to $u_\alpha \Psi^{u_\alpha}$ defined in Definition 5.1. Therefore, we consider the following decomposition

Definition 6.1. 1) Pick

$$\kappa_2 = \begin{cases} \alpha^{-1} \frac{\langle \Psi^\perp, \Phi_*^2 \pi^0 \Psi^\perp \rangle_\#}{\langle \Phi_*^2 \pi^0 \Psi^\perp, \Phi_*^2 \pi^0 \Psi^\perp \rangle_\#} & \text{if } \|\pi^0 \Psi^\perp\| > \alpha^{\frac{3}{2}}, \\ 0 & \text{if } \|\pi^0 \Psi^\perp\| \leq \alpha^{\frac{3}{2}}. \end{cases}$$

2) We split Ψ^\perp into Ψ_1^\perp and Ψ_2^\perp as follows: $\forall n \geq 0$, $\pi^n \Psi^\perp = \pi^n \Psi_1^\perp + \pi^n \Psi_2^\perp$ and for $n = 0$,

$$\pi^0 \Psi_1^\perp = \pi^0 \Psi^\perp \quad \text{and} \quad \pi^0 \Psi_2^\perp = 0,$$

for $n = 1$,

$$\pi^1 \Psi_1^\perp = \kappa_1 \Phi_\#^{u_\alpha} \quad \text{and} \quad \langle \pi^1 \Psi_2^\perp, \Phi_\#^{u_\alpha} \rangle_\# = 0,$$

for $n = 2$,

$$\pi^2 \Psi_1^\perp = \alpha \kappa_2 \Phi_*^2 \pi^0 \Psi_1^\perp + \sum_{i=1}^3 \alpha \kappa_{2,i} (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i},$$

$$\text{with } W_i = P_f^i \Phi_*^2 - 2A^+ . P_f (H_f + P_f^2)^{-1} (A^+)^i \Omega_f,$$

$$\pi^2 \Psi_2^\perp = \pi^2 \Psi^\perp - \pi^2 \Psi_1^\perp, \quad \text{with } \langle \pi^2 \Psi_2^\perp, (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \rangle_\# = 0 \quad (i = 1, 2, 3),$$

$$\text{and } \langle \pi^2 \Psi_2^\perp, \Phi_*^2 \pi^0 \Psi_1^\perp \rangle_\# = 0,$$

for $n = 3$,

$$\pi^3 \Psi_1^\perp = \alpha \kappa_3 (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha},$$

$$\pi^3 \Psi_2^\perp = \pi^3 \Psi^\perp - \pi^3 \Psi_1^\perp, \quad \langle \pi^3 \Psi_2^\perp, (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha} \rangle_\# = 0,$$

and for $n \geq 4$,

$$\pi^n \Psi_1^\perp = 0 \quad \text{and} \quad \pi^n \Psi_2^\perp = \pi^n \Psi^\perp.$$

The next step consists in establishing, in the next lemma, some a priori estimates concerning the function Ψ^\perp that give additional information to those obtained in (45) of Theorem 5.2.

Lemma 6.1. The following estimates hold

$$\kappa_1 = 1 + \mathcal{O}(\alpha^{\frac{1}{2}}),$$

$$|\kappa_2| \|\pi^0 \Psi_1^\perp\| = \mathcal{O}(\alpha),$$

$$\kappa_{2,i} = \mathcal{O}(1), \quad (i = 1, 2, 3),$$

$$\|\pi^0 \Psi_1^\perp\| = \mathcal{O}(\alpha),$$

$$\|P \pi^0 \Psi_1^\perp\| = \mathcal{O}(\alpha^2).$$

Proof. To derive these estimates, we use Theorem 5.2.

The first equality is a consequence of (48).

To derive the next two estimates, we first notice that (45) yields

$$\begin{aligned} \|P\pi^2\Delta_{\#}^{\perp}\|^2 &\leq 2\|(P - P_f)\Delta_{\#}^{\perp}\|^2 + 2\|P_f\pi^2\Delta_{\#}^{\perp}\|^2 \\ &\leq 2\|(P - P_f)\Delta_{\#}^{\perp}\|^2 + 2c\|H_f^{\frac{1}{2}}\pi^2\Delta_{\#}^{\perp}\|^2 = \mathcal{O}(\alpha^4), \end{aligned}$$

therefore, using again (45), we obtain

$$\|(h_{\alpha} + e_0)^{\frac{1}{2}}\pi^2\Delta_{\#}^{\perp}\|^2 \leq \|P\pi^2\Delta_{\#}^{\perp}\|^2 + e_0\|\pi^2\Delta_{\#}^{\perp}\|^2 = \mathcal{O}(\alpha^4),$$

and thus, using from (45) that $\|H_f^{\frac{1}{2}}\pi^2\Delta_{\#}^{\perp}\|^2 = \mathcal{O}(\alpha^4)$, we get

$$(51) \quad \|\pi^2\Delta_{\#}^{\perp}\|_{\#}^2 = \|\pi^2\Psi^{\perp}\|_{\#}^2 = \mathcal{O}(\alpha^4).$$

We then write, using (51) and the $\langle \cdot, \cdot \rangle_{\#}$ -orthogonality of $\pi^2\Psi_1^{\perp}$ and $\pi^2\Psi_2^{\perp}$,

$$(52) \quad \|\pi^2\Psi_1^{\perp}\|_{\#}^2 \leq \|\pi^2\Psi^{\perp}\|_{\#}^2 = \mathcal{O}(\alpha^4).$$

Since $\|\pi^2\Psi_1^{\perp}\|_{\#} \leq \|\pi^2\Psi_1^{\perp}\|_{*}$, and using (206) of Lemma C.2, we obtain

$$\begin{aligned} (53) \quad \mathcal{O}(\alpha^4) &= \|\pi^2\Psi_1^{\perp}\|_{*}^2 = \|\alpha\kappa_2\Phi_{*}^2\pi^0\Psi_1^{\perp}\|_{*}^2 + \|\alpha\sum_i \kappa_{2,i}(H_f + P_f^2)^{-1}W_i\frac{\partial u_{\alpha}}{\partial x_i}\|_{*}^2 \\ &= \alpha^2\|\Phi_{*}^2\|_{*}^2|\kappa_2|^2\|\pi^0\Psi_1^{\perp}\|_{*}^2 + \frac{\alpha^2}{3}\|\nabla u_{\alpha}\|^2\sum_i |\kappa_{2,i}|^2\|(H_f + P_f^2)^{-1}W_i\|_{*}^2. \end{aligned}$$

which gives

$$\kappa_{2,i} = \mathcal{O}(1) \quad \text{and} \quad |\kappa_2|\|\pi^0\Psi_1^{\perp}\| = \mathcal{O}(\alpha).$$

The last two estimates are consequences of (45). \square

Eventually, to derive the lower bound on the quadratic form of $\langle H(u_{\alpha}\Psi^{u_{\alpha}} + \Psi^{\perp}), u_{\alpha}\Psi^{u_{\alpha}} + \Psi^{\perp} \rangle$, yielding the upper bound (50) of Theorem 6.1, we will follow the same strategy as in Section 5, the only difference being that now we have better a priori estimates on $\Psi^{u_{\alpha}}$ and Ψ^{\perp} .

The two main results to achieve this are stated below, with the computation of the contribution to the ground state energy of the cross term (Proposition 6.1) and of the direct term $\langle H\Psi^{\perp}, \Psi^{\perp} \rangle$ (Proposition 6.2). Equipped with this two propositions, and using Theorem 5.2 and Lemma 6.1, the proof of Theorem 6.1 is only a straightforward computation which is detailed in section 8.

Proposition 6.1. *We have*

$$\begin{aligned} (54) \quad 2\text{Re} \langle H\Psi^{\perp}, u_{\alpha}\Psi^{u_{\alpha}} \rangle &\geq -\frac{1}{3}\alpha^4\text{Re} \sum_{i=1}^3 \overline{\kappa_{2,i}} \langle (H_f + P_f^2)^{-1}P_f^i\Phi_{*}^2, W_i \rangle \\ &\quad - 4\alpha\text{Re} \langle \nabla u_{\alpha} \cdot P_f\Phi_{*}^2, \pi^2\Psi_2^{\perp} \rangle - 2\text{Re} \overline{\kappa_1}\pi^0\Psi^{u_{\alpha}}\|\Phi_{\#}^{u_{\alpha}}\|_{\#}^2 \\ &\quad - \frac{2}{3}\alpha^4\text{Re} \sum_{i=1}^3 \langle (H_f + P_f^2)^{-\frac{1}{2}}(A^{-})^i\Phi_{*}^2, (H_f + P_f^2)^{-\frac{1}{2}}(A^{+})^i\Omega_f \rangle \\ &\quad - \epsilon\alpha^2\|(\Psi_1^{\perp})^a\|^2 - \frac{5}{8}M[\Psi_2^{\perp}] - \epsilon\alpha^5\log \alpha^{-1} - \epsilon\alpha^5|\kappa_3|^2 - |\kappa_1 - 1|c\alpha^4 + \mathcal{O}(\alpha^5). \end{aligned}$$

The proof of this Proposition is detailed in Appendix A

Proposition 6.2.

$$\begin{aligned}
\langle H\Psi^\perp, \Psi^\perp \rangle &\geq (\Sigma_0 - e_0) \|\Psi^\perp\|^2 - 4\alpha \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 \\
&+ |\kappa_1|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 + 2\alpha \|A^- \Phi_*^{u_\alpha}\|^2 + \frac{\alpha^4}{12} \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1} W_i\|_*^2 \\
(55) \quad &+ \frac{2}{3} \alpha^4 \operatorname{Re} \sum_{i=1}^3 \kappa_{2,i} \langle P_f A^- (H_f + P_f^2)^{-1} W_i, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle \\
&+ 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle \pi^2 \Psi_2^\perp, A^+ P_f \Phi_*^{u_\alpha} \rangle + M_1[\Psi^\perp],
\end{aligned}$$

where

$$\begin{aligned}
(56) \quad M_1[\Psi^\perp] &= (1 - c_0 \alpha) \|(h_\alpha + e_0)^{\frac{1}{2}} \pi^0 (\Psi_1^\perp)^a\|^2 + \frac{|\kappa_3 + 1|^2}{2} \alpha^2 \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 \\
&- |\kappa_1 - 1| c \alpha^4 + \frac{3}{4} M[\Psi_2^\perp] + o(\alpha^5 \log \alpha^{-1}),
\end{aligned}$$

and Q_α^\perp is the projection onto the orthogonal complement to the ground state u_α of the Schrödinger operator $h_\alpha = -\Delta - \frac{\alpha}{|x|}$

The proof of this Proposition is detailed in Appendix B.

7. PROOF OF THEOREM 5.2

We prove Theorem 5.2 by bounding the individual terms in the expression for the binding energy.

7.1. Estimate of the term $\langle H u_\alpha \Psi^{u_\alpha}, u_\alpha \Psi^{u_\alpha} \rangle$.**Lemma 7.1.**

$$\begin{aligned}
\langle H u_\alpha \Psi^{u_\alpha}, u_\alpha \Psi^{u_\alpha} \rangle &\geq -e_0 \|\Psi^{u_\alpha}\|^2 - \alpha^2 \|\Phi_*^2\|_*^2 \|\pi^0 \Psi^{u_\alpha}\|^2 + \alpha^2 |\eta_2 - \pi^0 \Psi^{u_\alpha}|^2 \|\Phi_*^2\|_*^2 \\
&+ \alpha^3 |\eta_2|^2 (2 \|A^- \Phi_*^2\|^2 - 4 \|\Phi_*^1\|_*^2 - 4 \|\Phi_*^3\|_*^2) \\
&+ 4\alpha^3 (|\eta_1 - \eta_2|^2 \|\Phi_*^1\|_*^2 + |\eta_3 - \eta_2|^2 \|\Phi_*^3\|_*^2) \\
&+ c \alpha^4 \log \alpha^{-1} (|\eta_1|^2 + |\eta_2|^2 + |\eta_3|^2 + \|\pi^0 \Psi^{u_\alpha}\|^2) + \frac{1}{2} \|\Delta_*^{u_\alpha}\|_*^2.
\end{aligned}$$

Proof. The proof is a trivial modification of the one given for the lower bound in [7, Theorem 3.1]. The only modification is that we have a slightly weaker estimate in Lemma 3.1 on the photon number for the ground state. This is accounted for by replacing the term of order α^4 in [7, Theorem 3.1] by a term of order $\alpha^4 \log \alpha^{-1}$. In addition, we need to use the equality $\langle P.(P_f + \sqrt{\alpha}A(0))u_\alpha \Psi^{u_\alpha}, u_\alpha \Psi^{u_\alpha} \rangle = 0$, due to the symmetry of u_α . \square

7.2. Estimates for the cross term $-4\operatorname{Re} \langle P.(P_f + \sqrt{\alpha}A(0))u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle$.**Lemma 7.2** ($-4\operatorname{Re} \langle P.P_f u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle$ term).

$$-4\operatorname{Re} \langle P.P_f u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle \geq -c \alpha^4 (|\eta_1|^2 + |\eta_2|^2 + |\eta_3|^2) - \epsilon \|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 - c \alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2.$$

Proof. • For $n = 1$ photon,

$$(57) \quad \langle P.P_f \pi^1 u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle = \langle P.P_f (\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \pi^1 \Delta_*^{u_\alpha}) u_\alpha, \kappa_1 \Phi_\#^{u_\alpha} + \pi^1 \Delta_\#^\perp \rangle.$$

Obviously

(58)

$$|\langle P.P_f(\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \pi^1 \Delta_*^{u_\alpha}) u_\alpha, \pi^1 \Delta_\#^\perp \rangle| \leq \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Delta_\#^\perp\|^2 + c\alpha^5 |\eta_1|^2 + c\alpha^2 \|H_f^{\frac{1}{2}} \pi^1 \Delta_*^{u_\alpha}\|^2.$$

Due to Lemma C.5 holds $\|H_f^{\frac{1}{2}}(\Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha})\|^2 = \mathcal{O}(\alpha^5)$, which implies

(59)

$$\begin{aligned} |\langle P.P_f(\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \pi^1 \Delta_*^{u_\alpha}) u_\alpha, \kappa_1 \Phi_\#^{u_\alpha} \rangle| &\leq |\kappa_1|^2 c\alpha^5 + c|\eta_1|^2 \alpha^5 + c\alpha^2 \|H_f^{\frac{1}{2}} \pi^1 \Delta_*^{u_\alpha}\|^2 \\ &+ |\langle P.P_f(\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \pi^1 \Delta_*^{u_\alpha}) u_\alpha, \kappa_1 \Phi_*^{u_\alpha} \rangle|. \end{aligned}$$

For the last term on the right hand side of (59), due to the orthogonality of $\frac{\partial u_\alpha}{\partial x_i}$ and $\frac{\partial u_\alpha}{\partial x_j}$, $i \neq j$, and the equality $\|\frac{\partial u_\alpha}{\partial x_i}\| = \|\frac{\partial u_\alpha}{\partial x_j}\|$, holds

(60)

$$\begin{aligned} |\langle P.P_f(\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \pi^1 \Delta_*^{u_\alpha}) u_\alpha, \kappa_1 \Phi_*^{u_\alpha} \rangle| &= \sum_{i=1}^3 \left\| \frac{\partial u_\alpha}{\partial x_i} \right\|^2 |\langle \eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \pi^1 \Delta_*^{u_\alpha}, \kappa_1 P_f^i(A^+ \Omega_f)^i \rangle| \\ &= c |\langle \eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \pi^1 \Delta_*^{u_\alpha}, \kappa_1 A^+ . P_f \Omega_f \rangle| = 0. \end{aligned}$$

- For $n = 2$ photons,

$$\begin{aligned} |\langle P.P_f \pi^2 u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle| &= |\langle P.P_f(\eta_2 \alpha \Phi_*^2 + \pi^2 \Delta_*^{u_\alpha}) u_\alpha, \pi^2 \Delta_\#^\perp \rangle| \\ &\leq c\alpha^4 |\eta_2|^2 + \epsilon \|H_f^{\frac{1}{2}} \pi^2 \Delta_\#^\perp\|^2 + \alpha \|H_f^{\frac{1}{2}} \pi^2 \Delta_*^{u_\alpha}\|^2. \end{aligned}$$

- For $n = 3$ photons, a similar estimate yields

$$\begin{aligned} |\langle P.P_f \pi^3 u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle| &= |\langle P.P_f(\eta_3 \alpha^{\frac{3}{2}} \Phi_*^3 + \pi^3 \Delta_*^{u_\alpha}) u_\alpha, \pi^3 \Delta_\#^\perp \rangle| \\ &\leq c\alpha^4 |\eta_3|^2 + \epsilon \|H_f^{\frac{1}{2}} \pi^3 \Delta_\#^\perp\|^2 + \alpha \|H_f^{\frac{1}{2}} \pi^3 \Delta_*^{u_\alpha}\|^2. \end{aligned}$$

- For $n \geq 4$ photons,

$$|\langle P.P_f \pi^{n \geq 4} u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle| \leq \epsilon \|H_f^{\frac{1}{2}} \pi^{n \geq 4} \Delta_\#^\perp\|^2 + \alpha \|H_f^{\frac{1}{2}} \pi^{n \geq 4} \Delta_*^{u_\alpha}\|^2.$$

□

Lemma 7.3 ($-4\text{Re}\langle \sqrt{\alpha} P.A(0) u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle$ term).

$$\begin{aligned} -4\text{Re}\langle \sqrt{\alpha} P.A(0) u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle &\geq -2\text{Re}\overline{\kappa_1} \pi^0 \Psi^{u_\alpha} \|\Phi_\#^{u_\alpha}\|_\#^2 - \epsilon \|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 - \epsilon \alpha^2 \|\Delta_\#^\perp\|^2 \\ &- c\alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 - c\alpha^4 (|\eta_1|^2 + |\eta_2|^2 + |\eta_3|^2 + |\kappa_1|^2) + \mathcal{O}(\alpha^5 \log \alpha^{-1}). \end{aligned}$$

Proof. We first estimate the term

$$\text{Re} \alpha^{\frac{1}{2}} \langle P.A^+ u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle = \alpha \text{Re} \sum_{n=0}^{\infty} \langle P.A^+ u_\alpha \pi^n \Psi^{u_\alpha}, \pi^{n+1} \Psi^\perp \rangle.$$

- For $n = 0$ photon, using the orthogonality $\langle \Phi_\#^{u_\alpha}, \pi^1 \Delta_\#^\perp \rangle_\# = 0$, yields

(61)

$$\begin{aligned} -\text{Re} \alpha^{\frac{1}{2}} \langle P.A^+ \pi^0 u_\alpha \Psi^{u_\alpha}, \kappa_1 \Phi_\#^{u_\alpha} + \pi^1 \Delta_\#^\perp \rangle &= -\frac{1}{2} \text{Re} \left((\pi^0 \Psi^{u_\alpha}) \langle \Phi_\#^{u_\alpha}, \kappa_1 \Phi_\#^{u_\alpha} + \pi^1 \Delta_\#^\perp \rangle_\# \right) \\ &= -\frac{1}{2} \text{Re} \left(\overline{\kappa_1} \pi^0 \Psi^{u_\alpha} \langle \Phi_\#^{u_\alpha}, \Phi_\#^{u_\alpha} \rangle_\# \right). \end{aligned}$$

- For $n \geq 1$ photons,

$$(62) \quad \begin{aligned} & \left| \sum_{n \geq 2} \operatorname{Re} \alpha^{\frac{1}{2}} \langle P.A^+ \pi^n u_\alpha \Delta_*^{u_\alpha}, \pi^{n+1} \Delta_\#^\perp \rangle \right| \leq c\alpha^3 \|\Delta_*^{u_\alpha}\|^2 + \epsilon \|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 \\ & \leq \epsilon \|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 + c\alpha^5 (|\eta_1|^2 + |\eta_2|^2 + |\eta_3|^2) + \mathcal{O}(\alpha^5 \log \alpha^{-1}), \end{aligned}$$

where we used from Lemma 3.1 that $\|\Delta_*^{u_\alpha}\|^2 \leq \mathcal{O}(\alpha^2 \log \alpha^{-1}) + c\alpha^3 (|\eta_1|^2 + |\eta_3|^2) + c\alpha^2 |\eta_2|^2$. We also have

$$(63) \quad \begin{aligned} & |\operatorname{Re} \alpha^{\frac{1}{2}} \langle P(2\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \eta_2 \alpha \Phi_*^2 + 2\eta_3 \alpha^{\frac{3}{2}} \Phi_*^3) u_\alpha, A^- \Delta_\#^\perp \rangle| \\ & \leq \epsilon \|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 + c(\alpha^6 |\eta_1|^2 + \alpha^5 |\eta_2|^2 + \alpha^6 |\eta_3|^2). \end{aligned}$$

We next estimate the term $2\operatorname{Re} \alpha^{\frac{1}{2}} \langle P.A^- u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle$. We first get

$$(64) \quad |\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A^- \Delta_*^{u_\alpha} u_\alpha, \Delta_\#^\perp \rangle| \leq \epsilon \alpha^2 \|\Delta_\#^\perp\|^2 + c\alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2.$$

Then we write

$$(65) \quad |\alpha^{\frac{1}{2}} \operatorname{Re} \langle (2\eta_1 \alpha^{\frac{3}{2}} A^- \Phi_*^1 + 2\eta_3 \alpha^{\frac{3}{2}} A^- \Phi_*^3) \nabla u_\alpha, \Delta_\#^\perp \rangle| \leq \epsilon \alpha^2 \|\Delta_\#^\perp\|^2 + c\alpha^4 (|\eta_1|^2 + |\eta_3|^2).$$

We also have

$$(66) \quad \begin{aligned} & |\alpha^{\frac{1}{2}} \operatorname{Re} \langle \eta_2 \alpha A^- \Phi_*^2 \cdot \nabla u_\alpha, \kappa_1 \Phi_\#^{u_\alpha} \rangle| \\ & = |\alpha^{\frac{1}{2}} \operatorname{Re} \langle \eta_2 \alpha H_f^{-\frac{1}{2}} A^- \Phi_*^2 \cdot \nabla u_\alpha, H_f^{\frac{1}{2}} \kappa_1 \Phi_\#^{u_\alpha} \rangle| \leq c\alpha^4 (|\eta_2|^2 + |\kappa_1|^2), \end{aligned}$$

since $H_f^{-\frac{1}{2}} A^- \Phi_*^2 \in L^2(\mathbb{R}^3)$ and $\|H_f^{\frac{1}{2}} \Phi_\#^{u_\alpha}\| = \mathcal{O}(\alpha^{\frac{3}{2}})$. Finally we get

$$(67) \quad \begin{aligned} & |\alpha^{\frac{1}{2}} \operatorname{Re} \langle \eta_2 \alpha A^- \Phi_*^2 \cdot \nabla u_\alpha, \pi^1 \Delta_\#^\perp \rangle| \\ & = |\alpha^{\frac{1}{2}} \operatorname{Re} \langle \eta_2 \alpha H_f^{-\frac{1}{2}} A^- \Phi_*^2 \cdot \nabla u_\alpha, H_f^{\frac{1}{2}} \pi^1 \Delta_\#^\perp \rangle| \leq \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Delta_\#^\perp\|^2 + c\alpha^5 |\eta_2|^2. \end{aligned}$$

Collecting (61) to (67) concludes the proof. \square

7.3. Estimate for the term $\langle H\Psi^\perp, \Psi^\perp \rangle$.

Lemma 7.4.

$$(68) \quad \langle H\Psi^\perp, \Psi^\perp \rangle \geq (\Sigma_0 - e_0) \|\Delta_\#^\perp\|^2 - c\alpha^4 |\kappa_1|^2 + \frac{1}{2} M[\Delta_\#^\perp] + |\kappa_1|^2 \|\Phi_\#^{u_\alpha}\|_\#^2,$$

where $M[\cdot]$ is defined in Corollary 4.2

Proof. Recall that

$$\begin{aligned} H &= (P^2 - \frac{\alpha}{|x|}) + (H_f + P_f^2) - 2\operatorname{Re} (P.P_f) \\ &\quad - 2\sqrt{\alpha}(P - P_f).A(0) + 2\alpha A^+.A^- + 2\alpha \operatorname{Re} (A^-)^2. \end{aligned}$$

Due to the orthogonality $\langle \Phi_{\#}^{u\alpha}, \Delta_{\#}^{\perp} \rangle_{\#} = 0$ we get

$$\begin{aligned}
 (69) \quad & \langle H\Psi^{\perp}, \Psi^{\perp} \rangle \\
 &= \langle Hr, r \rangle + |\kappa_1|^2 \|\Phi_{\#}^{u\alpha}\|_{\#}^2 - e_0 |\kappa_1|^2 \|\Phi_{\#}^{u\alpha}\|^2 - e_0 \kappa_1 \langle \Delta_{\#}^{\perp}, \Phi_{\#}^{u\alpha} \rangle + 2\alpha |\kappa_1|^2 \|A^{-} \Phi_{\#}^{u\alpha}\|^2 \\
 &\quad - 2\operatorname{Re} \langle P.P_f \Phi_{\#}^{u\alpha}, \Phi_{\#}^{u\alpha} \rangle + 2\alpha \operatorname{Re} \langle A^{-}.A^{-} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle - 4\sqrt{\alpha} \operatorname{Re} \langle P.A^{-} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle \\
 &\quad - 4\sqrt{\alpha} \operatorname{Re} \langle P.A^{+} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle + 4\operatorname{Re} \sqrt{\alpha} \langle P_f.A(0) \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle + 4\alpha \operatorname{Re} \langle A^{+}.A^{-} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle \\
 &\quad - 4\operatorname{Re} \langle P.P_f \Phi_{\#}^{u\alpha}, \Delta_{\#}^{\perp} \rangle.
 \end{aligned}$$

For the first term on the right hand side of (69), we have, from Corollary 4.2

$$(70) \quad \langle H\Delta_{\#}^{\perp}, \Delta_{\#}^{\perp} \rangle \geq (\Sigma_0 - e_0) \|\Delta_{\#}^{\perp}\|^2 + M[\Delta_{\#}^{\perp}].$$

According to Lemma C.4, we obtain

$$(71) \quad -e_0 \kappa_1 \langle \Delta_{\#}^{\perp}, \Phi_{\#}^{u\alpha} \rangle - e_0 |\kappa_1|^2 \|\Phi_{\#}^{u\alpha}\|^2 \geq -\epsilon \alpha^2 \|\Delta_{\#}^{\perp}\|^2 - c |\kappa_1|^2 \alpha^5 \log \alpha^{-1}.$$

The next term, namely $2\alpha |\kappa_1|^2 \|A^{-} \Phi_{\#}^{u\alpha}\|^2$, is positive.

Due to the symmetry in x -variable,

$$(72) \quad \langle P.P_f \Phi_{\#}^{u\alpha}, \Phi_{\#}^{u\alpha} \rangle = 0.$$

The term $2\alpha \operatorname{Re} \langle A^{-}.A^{-} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle$ is estimated as

$$(73) \quad 2\alpha \operatorname{Re} \langle A^{-}.A^{-} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle \geq -c\alpha^2 \|\kappa_1 \Phi_{\#}^{u\alpha}\|^2 - \frac{1}{4} \nu \|H_f^{\frac{1}{2}} \Delta_{\#}^{\perp}\|^2 = -\frac{1}{4} \nu \|H_f^{\frac{1}{2}} \Delta_{\#}^{\perp}\|^2 - c |\kappa_1|^2 \alpha^5 \log \alpha^{-1}.$$

Due to Lemma C.4 and [13, Lemma A4],

$$(74) \quad |\alpha^{\frac{1}{2}} \langle A^{-} \Delta_{\#}^{\perp}, P\kappa_1 \Phi_{\#}^{u\alpha} \rangle| \leq \frac{\nu}{8} \|H_f^{\frac{1}{2}} \Delta_{\#}^{\perp}\|^2 + c |\kappa_1|^2 \alpha^6 \log \alpha^{-1}.$$

The next term we have to estimate fulfils

$$(75) \quad |\alpha^{\frac{1}{2}} \langle P.A^{+} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle| \leq \epsilon \|P\pi^0 \Delta_{\#}^{\perp}\|^2 + c\alpha |\kappa_1|^2 \|A^{-} \Phi_{\#}^{u\alpha}\|^2 \leq \epsilon \|(P - P_f) \Delta_{\#}^{\perp}\|^2 - c |\kappa_1|^2 \alpha^4.$$

We have

$$\begin{aligned}
 (76) \quad & \operatorname{Re} \sqrt{\alpha} \langle P_f.A(0) \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle = \operatorname{Re} \sqrt{\alpha} \langle P_f.A^{-} \pi^2 \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle + \operatorname{Re} \sqrt{\alpha} \langle P_f.A^{+} \pi^0 \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle \\
 &= \operatorname{Re} \sqrt{\alpha} \langle P_f.A^{-} \pi^2 \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle + \operatorname{Re} \sqrt{\alpha} \langle A^{+}.P_f \pi^0 \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle
 \end{aligned}$$

Obviously, $\langle A^{+}.P_f \pi^0 \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle = 0$, and the first term is bounded by

$$\begin{aligned}
 (77) \quad & |\operatorname{Re} \sqrt{\alpha} \langle P_f.A^{-} \pi^2 \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle| \leq \epsilon \|H_f^{\frac{1}{2}} \pi^2 \Delta_{\#}^{\perp}\|^2 + c\alpha |\kappa_1|^2 \|P_f \Phi_{\#}^{u\alpha}\|^2 \\
 &\leq \epsilon \|H_f^{\frac{1}{2}} \pi^2 \Delta_{\#}^{\perp}\|^2 + c |\kappa_1|^2 \alpha^4.
 \end{aligned}$$

For the term $\alpha \operatorname{Re} \langle A^{+}.A^{-} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle$ we obtain

$$(78) \quad \alpha \operatorname{Re} \langle A^{+}.A^{-} \Delta_{\#}^{\perp}, \kappa_1 \Phi_{\#}^{u\alpha} \rangle \geq -c\alpha^2 \|H_f^{\frac{1}{2}} \kappa_1 \Phi_{\#}^{u\alpha}\|^2 |\kappa_1|^2 - \epsilon \|H_f^{\frac{1}{2}} \Delta_{\#}^{\perp}\|^2 = -c\alpha^5 |\kappa_1|^2 - \epsilon \|H_f^{\frac{1}{2}} \Delta_{\#}^{\perp}\|^2.$$

According to (214) of Lemma C.4, the last term we have to estimate fulfils

$$\begin{aligned}
 (79) \quad & \operatorname{Re} \langle P.P_f \kappa_1 \Phi_{\#}^{u\alpha}, \Delta_{\#}^{\perp} \rangle \leq \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Delta_{\#}^{\perp}\|^2 + c \|P|P_f|^{\frac{1}{2}} \kappa_1 \Phi_{\#}^{u\alpha}\|^2 \leq \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Delta_{\#}^{\perp}\|^2 + c\alpha^5 |\kappa_1|^2.
 \end{aligned}$$

Collecting the estimates (69) to (79) yields (68). \square

7.4. Upper bound on the binding energy with error term $\mathcal{O}(\alpha^4)$. We first establish a lemma that we shall need in the proof of Theorem 5.2 in order to improve the error term from $\mathcal{O}(\alpha^4 \log \alpha^{-1})$ to $\mathcal{O}(\alpha^4)$.

Lemma 7.5. *If $\|\Delta_*^{u_\alpha}\|_*^2 = \mathcal{O}(\alpha^4 \log \alpha^{-1})$ and $\|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 = \mathcal{O}(\alpha^4 \log \alpha^{-1})$ hold then we have*

$$(80) \quad \|N_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^{\frac{33}{16}})$$

$$(81) \quad \|N_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 = \mathcal{O}(\alpha^{\frac{33}{16}})$$

Proof. We note that from Definition 5.1, we have

$$\|a_\lambda(k) \Psi^{u_\alpha}\|^2 \leq \|a_\lambda(k) \Psi^{\text{GS}}\|^2 \leq \frac{c\alpha^{-\frac{3}{2}} \chi_\Lambda(|k|)}{|k|},$$

where in the last inequality, we used (14). Taking into account that

$$\Delta_*^{u_\alpha} = \Psi^{u_\alpha} - 2\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 - \eta_2 \alpha \Phi_*^2 - 2\eta_3 \alpha^{\frac{3}{2}} \Phi_*^3,$$

where

$$\|a_\lambda(k) \Phi_*^1\|^2 \leq \frac{c\chi_\Lambda(|k|)}{|k|}, \quad \|a_\lambda(k) \Phi_*^2\|^2 \leq \frac{c\chi_\Lambda(|k|)(1 + |\log |k||)}{|k|}$$

$$\|a_\lambda(k) \Phi_*^3\|^2 \leq \frac{c\chi_\Lambda(|k|)}{|k|},$$

and using (85), we arrive at

$$\|a_\lambda(k) \Delta_*^{u_\alpha}\|^2 \leq \frac{c\alpha^{-\frac{3}{2}} \chi_\Lambda(|k|)(1 + |\log |k||)}{|k|}.$$

For the expected photon number of $\Delta_*^{u_\alpha}$ thus holds

$$\begin{aligned} \|N_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 &= \sum_\lambda \int \|a_\lambda(k) \Delta_*^{u_\alpha}\|^2 \mathfrak{k} \\ &\leq \sum_\lambda \int_{|k| \leq \alpha^{\frac{15}{8}}} \frac{c\alpha^{-\frac{3}{2}}(1 + |\log |k||)}{|k|} \mathfrak{k} + \int_{|k| > \alpha^{\frac{15}{8}}} |k|^{-1} |k| \|a_\lambda(k) \Delta_*^{u_\alpha}\|^2 \mathfrak{k} \\ &\leq c\alpha^{\frac{17}{8}} + c\alpha^{-\frac{15}{8}} \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 \leq c\alpha^{\frac{33}{16}}, \end{aligned}$$

using (87) in the last inequality. The relation (81) can be proved similarly, using

$$\|a_\lambda(k) \Phi_\#^{u_\alpha}\|^2 \leq c \frac{\alpha^{-1}}{|k|}.$$

This concludes the proof of the lemma. \square

7.5. Concluding the proof of Theorem 5.2. The proof of Theorem 5.2 is obtained in the following two steps. We first show that the estimate holds with an error term $\mathcal{O}(\alpha^4 \log \alpha^{-1})$. In a second step, using Lemma 7.5, we improve to an error term $\mathcal{O}(\alpha^4)$. Then, we derive the estimates (45)-(48) that shall be used in the next section for the computation of the binding energy to higher order.

• Step 1: We first show that (44) holds with an error estimate of the order $\alpha^4 \log \alpha^{-1}$.

Collecting Lemmata 7.1, 7.2, 7.3 and Lemma 7.4 yields

$$\begin{aligned}
 (82) \quad & \langle H\Psi^{\text{GS}}, \Psi^{\text{GS}} \rangle \\
 & \geq -e_0 \|\Psi^{u_\alpha}\|^2 - \alpha^2 \|\Phi_*^2\|_*^2 |\pi^0 \Psi^{u_\alpha}|^2 + \alpha^3 |\eta_2|^2 (2\|A^- \Phi_*^2\|^2 - 4\|\Phi_*^1\|_*^2 - 4\|\Phi_*^3\|_*^2) \\
 & \quad + \alpha^2 |\eta_2 - \pi^0 \Psi^{u_\alpha}|^2 \|\Phi_*^2\|_*^2 + 4\alpha^3 (|\eta_1 - \eta_2|^2 \|\Phi_*^1\|_*^2 + |\eta_3 - \eta_2|^2 \|\Phi_*^3\|_*^2) \\
 & \quad - c\alpha^4 \log \alpha^{-1} (|\eta_1|^2 + |\eta_2|^2 + |\eta_3|^2 + |\pi^0 \Psi^{u_\alpha}|^2) + \frac{1}{4} \|\Delta_*^{u_\alpha}\|_*^2 \\
 & \quad + |\kappa_1|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 + (\Sigma_0 - e_0) \|\Delta_\#^\perp\|^2 + \frac{1}{4} M[\Delta_\#^\perp] - c\alpha^4 |\kappa_1|^2 \\
 & \quad - 2\text{Re}(\kappa_1 \pi^0 \Psi^{u_\alpha}) \|\Phi_\#^{u_\alpha}\|_\#^2 + \mathcal{O}(\alpha^5 \log \alpha^{-1}).
 \end{aligned}$$

We first estimate

$$\begin{aligned}
 (83) \quad & |\kappa_1|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 - 2\text{Re}(\overline{\kappa_1} \pi^0 \Psi^{u_\alpha}) \|\Phi_\#^{u_\alpha}\|_\#^2 - c\alpha^4 |\kappa_1|^2 \\
 & \geq -\|\Phi_\#^{u_\alpha}\|_\#^2 |\pi^0 \Psi^{u_\alpha}|^2 + \frac{|\overline{\kappa_1} - \pi^0 \Psi^{u_\alpha}|^2}{2} \|\Phi_\#^{u_\alpha}\|_\#^2 + \mathcal{O}(\alpha^4).
 \end{aligned}$$

Moreover, since $|\pi^0 \Psi^{u_\alpha}| \leq 1$, we replace in (82) $-\alpha^2 \|\Phi_*^2\|_*^2 |\pi^0 \Psi^{u_\alpha}|^2$ by $-\alpha^2 \|\Phi_*^2\|_*^2$ and in (83) $-\|\Phi_\#^{u_\alpha}\|_\#^2 |\pi^0 \Psi^{u_\alpha}|^2$ by $-\|\Phi_\#^{u_\alpha}\|_\#^2$. In addition, using the inequalities

$$|\eta_2 - \eta_j|^2 \geq \frac{1}{2} |\eta_j - \pi^0 \Psi^{u_\alpha}|^2 - |\eta_2 - \pi^0 \Psi^{u_\alpha}|^2 \quad \text{and} \quad |\eta_j|^2 \leq 2|\eta_j - \pi^0 \Psi^{u_\alpha}|^2 + 2$$

for $j = 1, 2, 3$ yields that for some $c > 0$,

$$\begin{aligned}
 (84) \quad & \langle H\Psi^{\text{GS}}, \Psi^{\text{GS}} \rangle \\
 & \geq -e_0 \|\Psi^{u_\alpha}\|^2 - \alpha^2 \|\Phi_*^2\|_*^2 + \alpha^3 |\eta_2|^2 (2\|A^- \Phi_*^2\|^2 - 4\|\Phi_*^1\|_*^2 - 4\|\Phi_*^3\|_*^2) - \|\Phi_\#^{u_\alpha}\|_\#^2 \\
 & \quad + c\alpha^2 |\eta_2 - \pi^0 \Psi^{u_\alpha}|^2 \|\Phi_*^2\|_*^2 + c\alpha^3 (|\eta_1 - \pi^0 \Psi^{u_\alpha}|^2 + |\eta_3 - \pi^0 \Psi^{u_\alpha}|^2) + c\alpha^3 |\kappa_1 - \pi^0 \Psi^{u_\alpha}|^2 \\
 & \quad + \frac{1}{4} \|\Delta_*^{u_\alpha}\|_*^2 + (\Sigma_0 - e_0) \|\Delta_\#^\perp\|^2 + \frac{1}{4} M[\Delta_\#^\perp] + \mathcal{O}(\alpha^4 \log \alpha^{-1}).
 \end{aligned}$$

Comparing this expression with (43) of Theorem 5.1 gives

$$(85) \quad \max\{|\eta_1|, |\eta_2|, |\eta_3|\} \leq 2.$$

and

$$(86) \quad \Sigma = \Sigma_0 - e_0 - \|\Phi_\#^{u_\alpha}\|_\#^2 + \mathcal{O}(\alpha^4 \log \alpha^{-1}).$$

Finally, by Lemma C.5, we can replace $\|\Phi_\#^{u_\alpha}\|_\#$ by $\|\Phi_*^{u_\alpha}\|_*$ in the above equality, which proves (44) with an error term $\mathcal{O}(\alpha^4 \log \alpha^{-1})$.

• Step 2: We now show that the error term does not contain any $\log \alpha^{-1}$ term.

From (84) we obtain

$$(87) \quad \|\Delta_*^{u_\alpha}\|_*^2 = \mathcal{O}(\alpha^4 \log \alpha^{-1}) \quad \text{and} \quad \|H_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 = \mathcal{O}(\alpha^4 \log \alpha^{-1}).$$

According to Lemma 7.5, this implies

$$\|N_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^{\frac{33}{16}}) \quad \text{and} \quad \|N_f^{\frac{1}{2}} \Delta_\#^\perp\|^2 = \mathcal{O}(\alpha^{\frac{33}{16}}).$$

Thus, we have

$$\begin{aligned} & \|N_f^{\frac{1}{2}} \Psi^{u_\alpha}\|^2 \\ & \leq 4|\eta_1|^2 \alpha^3 \|N^{\frac{1}{2}} \Phi_*^1\|^2 + |\eta_2|^2 \alpha^2 \|N^{\frac{1}{2}} \Phi_*^2\|^2 + 4|\eta_3|^2 \alpha^3 \|N^{\frac{1}{2}} \Phi_*^3\|^2 + \|N^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 \\ & = \mathcal{O}(\alpha^2), \end{aligned}$$

which implies that in Lemma 7.1 we can replace the term $c\alpha^4 \log \alpha^{-1} (|\eta_1|^2 + |\eta_2|^2 + |\eta_3|^2 + |\pi^0 \Psi^{u_\alpha}|^2)$ with $c\alpha^4 (|\eta_1|^2 + |\eta_2|^2 + |\eta_3|^2 + |\pi^0 \Psi^{u_\alpha}|^2)$ and consequently in (84) and (86), the term $\mathcal{O}(\alpha^4 \log \alpha^{-1})$ can be replaced by $\mathcal{O}(\alpha^4)$. This proves (44). The estimates (45)-(48) follow from (44) and (84) with $\mathcal{O}(\alpha^4 \log \alpha^{-1})$ replaced with $\mathcal{O}(\alpha^4)$. We thus arrive at the proof of Theorem 5.2. \square

8. PROOF OF THEOREM 6.1

In this section, we prove the upper bound on the binding energy up to the order $\alpha^5 \log \alpha^{-1}$ provided in Theorem 6.1.

We have

$$(88) \quad \langle H\Psi^{\text{GS}}, \Psi^{\text{GS}} \rangle \geq (\Sigma_0 - e_0) \|\Psi^{u_\alpha}\|^2 + 2\text{Re} \langle H\Psi^\perp, u_\alpha \Psi^{u_\alpha} \rangle + \langle H\Psi^\perp, \Psi^\perp \rangle.$$

The estimates for the last two terms in (88) are given in Propositions 6.1 and 6.2. We will bound below this expression by considering separately the terms involving the parameters κ_1 , $\kappa_{2,i}$, and κ_3 .

• We first estimate the second term on the right hand side of (54) together with the seventh term on the right hand side of (55). We have

$$(89) \quad \begin{aligned} & -4\alpha \text{Re} \langle \nabla u_\alpha \cdot P_f \Phi_*^2, \pi^2 \Psi_2^\perp \rangle + 4\alpha^{\frac{1}{2}} \text{Re} \langle \pi^2 \Psi_2^\perp, A^+ \cdot P_f \Phi_*^{u_\alpha} \rangle \\ & = -4\alpha \text{Re} \langle \pi^2 \Psi_2^\perp, (\sum_{i=1}^3 P_f^i \Phi_*^2 - 2A^+ \cdot P_f (H_f + P_f^2)^{-1} (A^+)^i \Omega_f) \frac{\partial u_\alpha}{\partial x_i} \rangle \\ & = -4\alpha \sum_{i=1}^3 \text{Re} \langle \pi^2 \Psi_2^\perp, (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \rangle_*. \end{aligned}$$

Using the $\langle \cdot, \cdot \rangle_\#$ -orthogonality of $\pi^2 \Psi_2^\perp$ and $(H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i}$, the last expression can be estimated as

$$(90) \quad -4\alpha \sum_{i=1}^3 \text{Re} \langle \pi^2 \Psi_2^\perp, (h_\alpha + e_0)(H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \rangle.$$

By the Schwarz inequality, this term is bounded below by

$$(91) \quad -\epsilon \alpha^2 \|\pi^2 \Psi_2^\perp\|^2 - c \sum_{i=1}^3 \|(h_\alpha + e_0) \frac{\partial u_\alpha}{\partial x_i}\|^2 = -\epsilon \alpha^2 \|\pi^2 \Psi_2^\perp\|^2 - \mathcal{O}(\alpha^6).$$

• Next, we collect all the terms involving κ_1 in (54) and (55). This yields

$$(92) \quad \begin{aligned} & -2\text{Re} \overline{\kappa_1} \pi^0 \Psi^{u_\alpha} \|\Phi_\#^{u_\alpha}\|_\#^2 + |\kappa_1|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 - |\kappa_1 - 1| c \alpha^4 \\ & \geq -|\pi^0 \Psi^{u_\alpha}|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 + |\kappa_1 - \pi^0 \Psi^{u_\alpha}|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 - |\kappa_1 - 1| c \alpha^4 \end{aligned}$$

Notice that from Theorem 5.2 we have $|\pi^0 \Psi^{u_\alpha}|^2 = 1 + \mathcal{O}(\alpha^2)$; moreover, we have $|\pi^0 \Psi^{u_\alpha}| = 1 + \mathcal{O}(\alpha^2)$. This yields

$$(93) \quad \begin{aligned} & -2\text{Re} \overline{\kappa_1} \pi^0 \Psi^{u_\alpha} \|\Phi_\#^{u_\alpha}\|_\#^2 + |\kappa_1|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 - |\kappa_1 - 1| c \alpha^4 \\ & \geq -\|\Phi_\#^{u_\alpha}\|_\#^2 + |\kappa_1 - 1|^2 c' \alpha^3 - |\kappa_1 - 1| c \alpha^4 + \mathcal{O}(\alpha^5) \geq -\|\Phi_\#^{u_\alpha}\|_\#^2 + \mathcal{O}(\alpha^5). \end{aligned}$$

• We now collect and estimates the terms in (54) and (55) involving $\kappa_{2,i}$. We get

$$\begin{aligned}
 (94) \quad & -\frac{1}{3}\alpha^4 \text{Re} \sum_{i=1}^3 \kappa_{2,i} \langle (H_f + P_f^2)^{-1} W_i, P_f^i \Phi_*^2 \rangle + \frac{\alpha^4}{12} \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1} W_i\|_*^2 \\
 & + \frac{2}{3}\alpha^4 \text{Re} \sum_{i=1}^3 \kappa_{2,i} \langle (H_f + P_f^2)^{-1} W_i, P_f \cdot A^+ (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle \\
 & = -\frac{1}{3}\alpha^4 \text{Re} \sum_{i=1}^3 \kappa_{2,i} \|(H_f + P_f^2)^{-1} W_i\|_*^2 + \frac{\alpha^4}{12} \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1} W_i\|_*^2 \\
 & \geq -\frac{\alpha^4}{3} \sum_{i=1}^3 \|(H_f + P_f^2)^{-1} W_i\|_*^2
 \end{aligned}$$

• Collecting in (54) and (55) the terms containing κ_3 yields

$$(95) \quad \frac{|\kappa_3 + 1|^2}{2} \alpha^2 \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 - \epsilon \alpha^5 |\kappa_3|^2 \geq c_1 \alpha^5 \log \alpha^{-1} |\kappa_3 + 1|^2 - \epsilon \alpha^5 |\kappa_3|^2 \geq -c_2 \alpha^5,$$

where c_1 and c_2 are positive constants.

• The fifth term on the right hand side of (54) and the first term on the right hand side of (56) are estimated, for α small enough, as

$$(96) \quad (1 - c_0 \alpha) \|(h_\alpha + e_0)^{\frac{1}{2}} \pi^0 (\Psi_1^\perp)^a\|^2 - \epsilon \alpha^2 \|(\Psi_1^\perp)^a\|^2 \geq \left(\frac{\delta}{2} - \epsilon \alpha^2\right) \|(\Psi_1^\perp)^a\|^2 \geq 0,$$

with $\delta = \frac{3}{32} \alpha^2$, and where we used that $(\Psi_1^\perp)^a$ is orthogonal to u_α .

• Substituting the above estimates in (88) yields

$$\begin{aligned}
 (97) \quad & \langle H \Psi^{\text{GS}}, \Psi^{\text{GS}} \rangle \\
 & \geq (\Sigma_0 - e_0) \|\Psi^{u_\alpha}\|^2 + (\Sigma_0 - e_0) \|\Psi^\perp\|^2 - \|\Phi_\#^{u_\alpha}\|_\#^2 - \alpha^4 \sum_{i=1}^3 \|(H_f + P_f^2)^{-1} W_i\|_*^2 \\
 & - \frac{2}{3} \alpha^4 \text{Re} \sum_{i=1}^3 \langle (H_f + P_f^2)^{-\frac{1}{2}} (A^-)^i \Phi_*^2, (H_f + P_f^2)^{-\frac{1}{2}} (A^+)^i \Omega_f \rangle \\
 & - 4\alpha \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 + 2\alpha \|A^- \Phi_*^{u_\alpha}\|^2 + o(\alpha^5 \log \alpha^{-1}),
 \end{aligned}$$

where Q_α^\perp is the projection onto the orthogonal complement to the ground state u_α of the Schrödinger operator $h_\alpha = -\Delta - \frac{\alpha}{|x|}$.

To complete the proof of Theorem 6.1 we first note that

$$(98) \quad \|\Psi^{u_\alpha}\|^2 + \|\Psi^\perp\|^2 = \|\Psi^{\text{GS}}\|^2.$$

Moreover, according to Lemma C.5

$$(99) \quad -\|\Phi_\#^{u_\alpha}\|_\#^2 = -\|\Phi_*^{u_\alpha}\|_*^2 + \frac{1}{3\pi} \|(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1\|^2 \alpha^5 \log \alpha^{-1} + o(\alpha^5 \log \alpha^{-1}),$$

and

$$(100) \quad \|\Phi_*^{u_\alpha}\|_*^2 = \frac{\alpha^3}{2\pi} \int_0^\infty \frac{\chi_\Lambda(t)}{1+t} dt = e^{(1)} \alpha^3.$$

In addition, we have the following identities ($i = 1, 2, 3$)

$$(101) \quad \begin{aligned} & \| (H_f + P_f^2)^{-1} W_i \|_*^2 = \\ & \| (H_f + P_f^2)^{-\frac{1}{2}} \left(2A^+ . P_f (H_f + P_f^2)^{-1} (A^+)^i - P_f^i (H_f + P_f^2)^{-1} A^+ . A^+ \right) \Omega_f \|^2, \end{aligned}$$

and

$$(102) \quad \begin{aligned} & -\frac{2}{3} \alpha^4 \operatorname{Re} \sum_{i=1}^3 \langle (H_f + P_f^2)^{-\frac{1}{2}} A^- \Phi_*^2, (H_f + P_f^2)^{-\frac{1}{2}} (A^+)^i \Omega_f \rangle \\ & = -\frac{2}{3} \alpha^4 \operatorname{Re} \sum_{i=1}^3 \langle A^- (H_f + P_f^2)^{-1} A^+ . A^+ \Omega_f, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle. \end{aligned}$$

We also have

$$(103) \quad -4\alpha \| (h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha} \|^2 = -4\alpha^4 a_0^2 \| (-\Delta - \frac{1}{|x|} + \frac{1}{4})^{-\frac{1}{2}} Q_1^\perp \Delta u_1 \|^2,$$

with

$$a_0 = \int \frac{k_1^2 + k_2^2}{4\pi^2 |k|^3} \frac{2}{|k|^2 + |k|} \chi_\Lambda(|k|) \mathbf{k}_1 \mathbf{k}_2 \mathbf{k}_3,$$

and

$$(104) \quad 2\alpha \| A^- \Phi_*^{u_\alpha} \|^2 = \frac{2}{3} \alpha^4 \sum_{i=1}^3 \| A^- (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \|^2.$$

Substituting (98)-(104) into (97) finishes the proof of Theorem 6.1 and thus the proof of the upper bound in Theorem 2.1. \square

9. PROOF OF THEOREM 2.1: LOWER BOUND UP TO $o(\alpha^5 \log \alpha^{-1})$ FOR THE BINDING ENERGY

In this section, we prove a lower bound for $\Sigma_0 - \Sigma$ in Theorem 2.1 which coincides with the upper bound given in (97). To this end, it suffices to compute

$$\frac{\langle (H - \Sigma_0 + e_0) \tilde{\Phi}^{\text{trial}}, \tilde{\Phi}^{\text{trial}} \rangle}{\| \tilde{\Phi}^{\text{trial}} \|^2},$$

with the following trial function

$$\tilde{\Phi}^{\text{trial}} = u_\alpha \Psi_0 + \tilde{\Psi}_0^\perp,$$

where Ψ_0 is ground state of the operator $T(0)$ (defined in (7)), with the normalization $\pi^0 \Psi_0 = \Omega_f$, u_α is the normalized ground state of $h_\alpha = -\Delta - \frac{\alpha}{|x|}$, and $\tilde{\Psi}_0^\perp$ is defined by

$$(105) \quad \begin{aligned} \pi^0 \tilde{\Psi}_0^\perp &= 2\alpha^{\frac{1}{2}} (h_\alpha + e_0)^{-1} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}, \quad \pi^1 \tilde{\Psi}_0^\perp = \Phi_\#^{u_\alpha}, \\ \pi^2 \tilde{\Psi}_0^\perp &= \alpha \Phi_*^2 \pi^0 \tilde{\Psi}_0^\perp + \sum_{i=1}^3 2\alpha (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i}, \\ \pi^3 \tilde{\Psi}_0^\perp &= -\alpha (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha}. \end{aligned}$$

Where $\Phi_*^{u_\alpha}$, $\Phi_\#^{u_\alpha}$, Φ_*^2 and W_i are defined as in Sections 5 and 6. Technical Lemmata used in this proof are given in Appendix C.

We compute

$$(106) \quad \langle H\tilde{\Phi}^{\text{trial}}, \tilde{\Phi}^{\text{trial}} \rangle = \langle Hu_\alpha \Psi_0, u_\alpha \Psi_0 \rangle + 2\text{Re} \langle Hu_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle + \langle H\tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle,$$

and we recall

$$(107) \quad H = h_\alpha + (H_f + P_f^2) - 2\text{Re} P.P_f - 2\alpha^{\frac{1}{2}} P.A(0) + 2\alpha^{\frac{1}{2}} P_f.A(0) + 2\alpha A^+.A^- + 2\alpha(A^-)^2.$$

- For the first term in (106), a straightforward computation shows

$$(108) \quad \langle Hu_\alpha \Psi_0, u_\alpha \Psi_0 \rangle = (\Sigma_0 - e_0) \|u_\alpha\|^2 \|\Psi_0\|^2.$$

- We estimate the second term on the right hand side of (106) by computing each term that occurs in the decomposition (107).

◇ Using the symmetry of u_α , the only non zero terms in $2\text{Re} \langle Hu_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle$ are given by

$$(109) \quad 2\text{Re} \langle Hu_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle = -4\text{Re} \langle P.P_f u_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle - 4\text{Re} \langle P.A^+ u_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle - 4\text{Re} \langle P.A^- u_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle.$$

◇ The first term on the right hand side of (109) is estimated with similar arguments as in Lemma 7.2, and using $\|\Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^3)$ (Lemma C.1) and $\|\Delta_*^{u_\alpha}\|_*^2 = \mathcal{O}(\alpha^4)$ ([7, Theorem 3.2]). We obtain

$$(110) \quad -4\text{Re} \langle P.P_f u_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle = -\frac{2}{3}\alpha^4 \sum_{i=1}^3 \langle P_f^i \Phi_*^2, (H_f + P_f^2) W_i \rangle + \mathcal{O}(\alpha^5 \sqrt{\log \alpha^{-1}}).$$

◇ The second and third terms on the right hand side of (109) are estimated as in Lemma 7.3, and using again $\|\Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^2)$ and $\|\Delta_*^{u_\alpha}\|_*^2 = \mathcal{O}(\alpha^4)$. This yields

$$(111) \quad -4\text{Re} \langle P.A^+ u_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle = -2\|\Phi_\#^{u_\alpha}\|_\#^2 + o(\alpha^5 \log \alpha^{-1}),$$

and

$$(112) \quad \begin{aligned} & -4\text{Re} \langle P.A^- u_\alpha \Psi_0, \tilde{\Psi}_0^\perp \rangle \\ &= -\frac{2}{3}\alpha^4 \sum_{i=1}^3 \langle (A^-)^i \Phi_*^2, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle + \mathcal{O}(\alpha^5 \sqrt{\log \alpha^{-1}}). \end{aligned}$$

- Next, we estimate the third term on the right hand side of (106). For that sake, we also use the decomposition (107) for H .

◇ For the term involving h_α , using $\|(h_\alpha + e_0)\pi^0 \tilde{\Psi}_0^\perp\| = \mathcal{O}(\alpha^3)$ (since $\|P.A^- \Phi_*^{u_\alpha}\| = \mathcal{O}(\alpha^{\frac{5}{2}})$), and $\|(h_\alpha + e_0)\frac{\partial u_\alpha}{\partial x_i}\| = \mathcal{O}(\alpha^3)$, we directly obtain

$$(113) \quad \begin{aligned} \langle h_\alpha \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle &= \langle (h_\alpha + e_0) \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle - e_0 \|\tilde{\Psi}_0^\perp\|^2 \\ &= 4\alpha \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P.A^- \Phi_*^{u_\alpha}\|^2 + \langle (h_\alpha + e_0) \Phi_\#^{u_\alpha}, \Phi_\#^{u_\alpha} \rangle - e_0 \|\tilde{\Psi}_0^\perp\|^2 + \mathcal{O}(\alpha^5). \end{aligned}$$

◇ For the term with $H_f + P_f^2$, we use the estimate (222) of Lemma C.6, and the $\langle \cdot, \cdot \rangle_*$ -orthogonality (see (206) of Lemma C.2) of the two vectors $\alpha \Phi_*^2 \pi^0 g$ and $\sum_{i=1}^3 2\alpha (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i}$ that occur in $\pi^2 \tilde{\Psi}_0^\perp$. We therefore obtain

$$(114) \quad \begin{aligned} \langle (H_f + P_f^2) \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle &= \langle (H_f + P_f^2) \Phi_\#^{u_\alpha}, \Phi_\#^{u_\alpha} \rangle + \alpha^2 \|\Phi_*^2 \pi^0 \tilde{\Psi}_0^\perp\|_*^2 \\ &+ \left\| \sum_{i=1}^3 2\alpha (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \right\|_*^2 + \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 + o(\alpha^5 \log \alpha^{-1}). \end{aligned}$$

◇ Using the symmetry of u_α , all terms in $\text{Re} \langle P.P_f \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle$ are zero, except the expression $\text{Re} \langle P.P_f \pi^2 \tilde{\Psi}_0^\perp, \pi^2 \tilde{\Psi}_0^\perp \rangle$, which is estimated as follows

$$\text{Re} \langle P.P_f \pi^2 \tilde{\Psi}_0^\perp, \pi^2 \tilde{\Psi}_0^\perp \rangle = 2\text{Re} \langle P.P_f \alpha \Phi_*^2 \pi^0 \tilde{\Psi}_0^\perp, \sum_{i=1}^3 2\alpha(H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \rangle = \mathcal{O}(\alpha^5),$$

where we used Lemma C.2 in the first equality to prove that only the crossed term remains. Therefore, we obtain

$$(115) \quad -2\text{Re} \langle P.P_f \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle = \mathcal{O}(\alpha^5).$$

◇ The terms involving $-2\alpha^{\frac{1}{2}} \text{Re} \langle P.A(0) \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle$ is estimated as in the proof of Lemma B.6. This yields

$$(116) \quad \begin{aligned} & -2\alpha^{\frac{1}{2}} \langle P.A(0) \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle = -4\text{Re} \alpha^{\frac{1}{2}} \langle P.A^+ \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle \\ & = -4\text{Re} \langle P.A^+ \pi^0 \tilde{\Psi}_0^\perp, \pi^1 \tilde{\Psi}_0^\perp \rangle + \mathcal{O}(\alpha^5 \sqrt{\log \alpha^{-1}}) \\ & = -8\alpha \| (h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P.A^- \Phi_*^{u_\alpha} \|^2 + \mathcal{O}(\alpha^5 \sqrt{\log \alpha^{-1}}). \end{aligned}$$

◇ For $2\alpha^{\frac{1}{2}} \text{Re} \langle P_f.A(0) \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle$, we proceed as in the proof of Lemma B.7, and obtain

$$(117) \quad \begin{aligned} & 2\alpha^{\frac{1}{2}} \langle P_f.A(0) \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle = 4\text{Re} \alpha^{\frac{1}{2}} \langle P_f.A^- \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle \\ & = 4\alpha^{\frac{1}{2}} \langle P_f.A^- \alpha \sum_{i=1}^3 2(H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i}, \Phi_*^{u_\alpha} \rangle + \mathcal{O}(\alpha^5 \sqrt{\log \alpha^{-1}}). \end{aligned}$$

◇ Using the symmetry of u_α and $\pi^0 \tilde{\Psi}_0^\perp$, the term $2\alpha \text{Re} \langle A^- . A^- \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle$ is estimated as follows,

$$(118) \quad \begin{aligned} & 2\alpha \text{Re} \langle A^- . A^- \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle \\ & = 2\alpha \text{Re} \langle A^- . A^- (\alpha \Phi_*^2 \pi^0 \tilde{\Psi}_0^\perp + \sum_{i=1}^3 2\alpha(H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i}), \pi^0 \tilde{\Psi}_0^\perp \rangle \\ & + 2\alpha \text{Re} \langle A^- . A^- (-\alpha(H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha}), \Phi_\#^{u_\alpha} \rangle \\ & = -2\alpha^2 \|\pi^0 \tilde{\Psi}_0^\perp\|^2 \|\Phi_*^2\|_*^2 - 2\alpha^2 \|(H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha}\|^2 \\ & = -2\alpha^2 \|\pi^0 \tilde{\Psi}_0^\perp\|^2 \|\Phi_*^2\|_*^2 - 2\alpha^2 \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 + o(\alpha^5 \log \alpha^{-1}), \end{aligned}$$

where in the last inequality we used (222) of Lemma C.6.

◇ Finally, a straightforward computation yields

$$(119) \quad 2\alpha \text{Re} \langle A^- . A^+ \tilde{\Psi}_0^\perp, \tilde{\Psi}_0^\perp \rangle = 2\alpha \|A^- \Phi_\#^{u_\alpha}\|^2 + \mathcal{O}(\alpha^5) = 2\alpha \|A^- \Phi_*^{u_\alpha}\|^2 + \mathcal{O}(\alpha^5),$$

where in the last equality, we used Lemma C.5.

• Before collecting (109)-(119), we show that gathering some terms yield simpler expressions. Namely, we have

$$(120) \quad \begin{aligned} & -\frac{2}{3} \alpha^4 \sum_{i=1}^3 \langle P_f^i \Phi_*^2, (H_f + P_f^2) W_i \rangle + 4\alpha^{\frac{1}{2}} \langle P_f.A^- \alpha \sum_{i=1}^3 2(H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i}, \Phi_*^{u_\alpha} \rangle \\ & + \left\| \sum_{i=1}^3 2\alpha(H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \right\|_*^2 = -\frac{1}{3} \alpha^4 \sum_{i=1}^3 \|(H_f + P_f^2)^{-1} W_i\|_*^2. \end{aligned}$$

We also have, using $-\alpha^2 \|\Phi_*^2\|_*^2 = \Sigma_0 + \mathcal{O}(\alpha^3)$ (see e.g. [7])

$$(121) \quad \begin{aligned} & (\Sigma_0 - e_0) \|\Psi_0\|^2 - e_0 \|\tilde{\Psi}_0^\perp\|^2 - \alpha^2 \|\Phi_*^2\|_*^2 \|\pi^0 \tilde{\Psi}_0^\perp\|^2 - \alpha^2 \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 \\ & = (\Sigma_0 - e_0) (\|\Psi_0\|^2 + \|\tilde{\Psi}_0^\perp\|^2) + \mathcal{O}(\alpha^5). \end{aligned}$$

Therefore, collecting (109)-(119), and using the two equalities (120)-(121), we obtain

$$\begin{aligned} \langle H(u_\alpha \Psi_0 + \tilde{\Psi}_0^\perp), u_\alpha \Psi_0 + \tilde{\Psi}_0^\perp \rangle &= (\Sigma_0 - e_0) (\|\Psi_0\|^2 + \|\tilde{\Psi}_0^\perp\|^2) - \frac{1}{3} \alpha^4 \sum_{i=1}^3 \|(H_f + P_f^2)^{-1} W_i\|_*^2 \\ &\quad - \frac{2}{3} \alpha^4 \sum_{i=1}^3 \langle (A^-)^i \Phi_*^2, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle - \|\Phi_\#^{u_\alpha}\|_\#^2 - 4\alpha \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 \\ &\quad + 2\alpha \|A^- \Phi_*^{u_\alpha}\|^2 + o(\alpha^5 \log \alpha^{-1}). \end{aligned}$$

With the definition $e^{(1)}$, $e^{(2)}$, and $e^{(3)}$, of Theorem 2.1 this expression can be rewritten as

$$(122) \quad \langle (H - \Sigma_0 + e_0) \tilde{\Phi}^{\text{trial}}, \tilde{\Phi}^{\text{trial}} \rangle = e^{(1)} \alpha^3 + e^{(2)} \alpha^4 + e^{(3)} \alpha^5 \log \alpha^{-1} + o(\alpha^5 \log \alpha^{-1}).$$

Using Lemma C.1 yields $\|\Psi_0\|^2 = 1 + \mathcal{O}(\alpha^2)$, which implies, due to the orthogonality of $\tilde{\Psi}_0^\perp$ and u_α in $L^2(\mathbb{R}^3, \mathfrak{x})$,

$$\|\tilde{\Phi}^{\text{trial}}\|^2 = \|u_\alpha\|^2 \|\Psi_0\|^2 + \|\tilde{\Psi}_0^\perp\|^2 = 1 + \mathcal{O}(\alpha^2).$$

Therefore, together with (122), this gives

$$\frac{\langle (H - \Sigma_0 + e_0) \tilde{\Phi}^{\text{trial}}, \tilde{\Phi}^{\text{trial}} \rangle}{\|\tilde{\Phi}^{\text{trial}}\|^2} = e^{(1)} \alpha^3 + e^{(2)} \alpha^4 + e^{(3)} \alpha^5 \log \alpha^{-1} + o(\alpha^5 \log \alpha^{-1}).$$

which concludes the proof of the lower bound in Theorem 2.1. \square

APPENDIX A. PROOF OF PROPOSITION 6.1

In this Appendix, we provide proofs of results that have a high level of technicality. To begin with, we establish Proposition 6.1.

Lemma A.1. *The following holds*

(123)

$$\begin{aligned} -4\operatorname{Re}\langle P.P_f u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle &\geq -\frac{4}{3}\alpha^2 \operatorname{Re} \sum_{i=1}^3 \overline{\kappa_{2,i}} \|\nabla u_\alpha\|^2 \langle (H_f + P_f^2)^{-1} P_f^i \Phi_*^2, W_i \rangle \\ &\quad - 4\alpha \operatorname{Re} \langle \nabla u_\alpha, P_f \Phi_*^2, \pi^2 \Psi_2^\perp \rangle - \epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 - \epsilon \alpha^5 |\kappa_3|^2 + \mathcal{O}(\alpha^5) \end{aligned}$$

Proof. For $n \neq 2, 3$, with the estimates from the proof of Lemma 7.2 and using that due to Theorem 5.2 we have

$$\|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^4), \quad |\eta_1| = \mathcal{O}(1), \quad \text{and} \quad |\kappa_1| = \mathcal{O}(1),$$

and since $\pi^1 \Delta_\#^\perp = \pi^1 \Psi_2^\perp$ and $\pi^{n \geq 4} \Delta_\#^\perp = \pi^{n \geq 4} \Psi_2^\perp$, we obtain

$$(124) \quad \sum_{n \neq 2, 3} -4\operatorname{Re} \langle P.P_f u_\alpha \pi^n \Psi^{u_\alpha}, \pi^n \Psi^\perp \rangle \geq -\epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5).$$

For $n = 2$,

(125)

$$\begin{aligned} -4\operatorname{Re} \langle \nabla u_\alpha, (\alpha \eta_2 P_f \Phi_*^2 + P_f \pi^2 \Delta_*^{u_\alpha}), \pi^2 \Psi^\perp \rangle &\geq -4\operatorname{Re} \langle \nabla u_\alpha, \alpha \eta_2 P_f \Phi_*^2, \pi^2 \Psi_1^\perp \rangle \\ &\quad - 4\operatorname{Re} \langle \nabla u_\alpha, \alpha \eta_2 P_f \Phi_*^2, \pi^2 \Psi_2^\perp \rangle - c\alpha \|H_f^{\frac{1}{2}} \pi^2 \Delta_*^{u_\alpha}\|^2 - c\alpha \|H_f^{\frac{1}{2}} \pi^2 \Psi^\perp\|^2. \end{aligned}$$

Using Theorem 5.2, the last two terms on the right hand side of (125) can be estimated by $\mathcal{O}(\alpha^5)$. For the first term on the right hand side of (125), using from Lemma C.2 that $\langle P_f^i \Phi_*^2, \Phi_*^2 \rangle = 0$, from Theorem 5.2 that $\eta_2 = 1 + \mathcal{O}(\alpha)$, and from Lemma 6.1 that $\kappa_{2,i} = \mathcal{O}(1)$, holds

$$\begin{aligned} &-4\alpha \operatorname{Re} \langle \nabla u_\alpha, \eta_2 P_f \Phi_*^2, \pi^2 \Psi_1^\perp \rangle \\ &= -4\operatorname{Re} \alpha \langle \nabla u_\alpha, \eta_2 P_f \Phi_*^2, \alpha \kappa_2 \Phi_*^2 \pi^0 \Psi^\perp \rangle \\ &\quad - 4\operatorname{Re} \alpha \langle \nabla u_\alpha, \eta_2 P_f \Phi_*^2, \sum_i \alpha \kappa_{2,i} (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \rangle \\ (126) \quad &= -\frac{4}{3} \operatorname{Re} \alpha^2 \|\nabla u_\alpha\|^2 \sum_i \overline{\kappa_{2,i}} \langle (H_f + P_f^2)^{-1} P_f^i \Phi_*^2, W_i \rangle \\ &= -\frac{1}{3} \alpha^4 \operatorname{Re} \sum_{i=1}^3 \overline{\kappa_{2,i}} \langle (H_f + P_f^2)^{-1} P_f^i \Phi_*^2, W_i \rangle. \end{aligned}$$

We also used that $\langle \frac{\partial u_\alpha}{\partial x_i}, \frac{\partial u_\alpha}{\partial x_j} \rangle = 0$ for $i \neq j$ and $\|\frac{\partial u_\alpha}{\partial x_i}\| = \|\frac{\partial u_\alpha}{\partial x_j}\|$ for all i and j .

Finally, the second term on the right hand side of (125) gives the second term on the right hand side of (123) plus $\mathcal{O}(\alpha^5)$, using from Theorem 5.2 that $|\eta_2 - 1|^2 = \mathcal{O}(\alpha^2)$ and $\|H_f^{\frac{1}{2}} \pi^2 \Psi_2^\perp\| = \mathcal{O}(\alpha^2)$.

To complete the proof, we shall estimate now the term for $n = 3$,

$$\begin{aligned} (127) \quad &4\operatorname{Re} \langle P.P_f u_\alpha \pi^3 \Psi^{u_\alpha}, \pi^3 \Psi^\perp \rangle = 4\operatorname{Re} \alpha^{\frac{3}{2}} 2\eta_3 \langle P.P_f u_\alpha \Phi_*^3, \pi^3 \Psi_1^\perp \rangle \\ &\quad + 4\operatorname{Re} \alpha^{\frac{3}{2}} 2\eta_3 \langle P.P_f u_\alpha \Phi_*^3, \pi^3 \Psi_2^\perp \rangle + 4\operatorname{Re} \langle P.P_f u_\alpha \pi^3 \Delta_*^{u_\alpha}, \pi^3 \Psi^\perp \rangle. \end{aligned}$$

The inequalities $\|H_f^{\frac{1}{2}}\Delta_*^{u_\alpha}\| \leq c\alpha^2$ and $\|H_f^{\frac{1}{2}}\pi^3\Psi^\perp\| \leq c\alpha^2$ (see Theorem 5.2) imply that the last term on the right hand side of (127) is $\mathcal{O}(\alpha^5)$. For the second term on the right hand side of (127) holds

$$(128) \quad \operatorname{Re} \alpha^{\frac{3}{2}} \eta_3 \langle P.P_f u_\alpha \Phi_*^3, \pi^3 \Psi_2^\perp \rangle \geq -\epsilon \|H_f^{\frac{1}{2}} \pi^3 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5),$$

since from Theorem 5.2 we have $\eta_3 = \mathcal{O}(1)$.

Finally to estimate the first term on the right hand side of (127), we note that

$$|k_1|^{-\frac{1}{6}} |k_2|^{-\frac{1}{6}} |k_3|^{-\frac{1}{6}} \Phi_*^3(k_1, k_2, k_3) \in L^2(\mathbb{R}^9, \mathbb{C}^6),$$

and from Lemma C.6,

$$\| |k_1|^{\frac{1}{6}} |k_2|^{\frac{1}{6}} |k_3|^{\frac{1}{6}} (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha} \|^2 = \mathcal{O}(\alpha^3).$$

This implies, using again $|\eta_3| = \mathcal{O}(1)$, and the explicit expression of $\pi^3 \Psi_1^\perp$

$$(129) \quad |\alpha^{\frac{3}{2}} 2\eta_3 \langle P.P_f u_\alpha \Phi_*^3, \pi^3 \Psi_1^\perp \rangle| \leq c\alpha^{\frac{5}{2}} |\kappa_3| \alpha^{\frac{3}{2}} |\eta_3| \|Pu_\alpha\| \leq \epsilon \alpha^5 |\kappa_3|^2 + \mathcal{O}(\alpha^5).$$

Collecting (124)-(129) concludes the proof. \square

Lemma A.2. *The following estimate holds*

$$(130) \quad \begin{aligned} & -4\sqrt{\alpha} \operatorname{Re} \langle P.A^+ u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle \\ & = -2\operatorname{Re} \overline{\kappa_1} \pi^0 \Psi^{u_\alpha} \|\Phi_\#^{u_\alpha}\|_\#^2 - \frac{1}{4} M[\Psi_2^\perp] - \alpha^5 |\kappa_3|^2 + \mathcal{O}(\alpha^5), \end{aligned}$$

Proof. Obviously

$$(131) \quad \begin{aligned} & -4\operatorname{Re} \alpha^{\frac{1}{2}} \langle P.A^+ u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle \\ & = -4\operatorname{Re} \alpha^{\frac{1}{2}} \langle P.A^+ u_\alpha (\pi^0 \Psi^{u_\alpha} + 2\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \eta_2 \alpha \Phi_*^2 + 2\eta_3 \alpha^{\frac{3}{2}} \Phi_*^3 + \Delta_*^{u_\alpha}), \Psi^\perp \rangle. \end{aligned}$$

\diamond *Step 1* From (131), let us first estimate the term

$$(132) \quad -4\operatorname{Re} \alpha^{\frac{1}{2}} \langle P.A^+ u_\alpha \Delta_*^{u_\alpha}, \Psi^\perp \rangle = -4\alpha^{\frac{1}{2}} \operatorname{Re} \sum_{n=0}^{\infty} \langle P.A^+ \pi^n u_\alpha \Delta_*^{u_\alpha}, \pi^{n+1} \Psi^\perp \rangle.$$

For $n = 0$, the corresponding term vanishes since $\pi^0 \Delta_*^{u_\alpha} = 0$.

For $n > 2$, we can use (62) where the term $\mathcal{O}(\alpha^5 \log \alpha^{-1})$ can be replaced with $\mathcal{O}(\alpha^5)$ because we know from Theorem 5.2 that $\|\Delta_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^{\frac{33}{16}})$.

For $n = 1$, we have

$$(133) \quad \begin{aligned} & |4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A^+ \pi^1 u_\alpha \Delta_*^{u_\alpha}, \pi^2 \Psi_1^\perp + \pi^2 \Psi_2^\perp \rangle| \leq c\alpha^3 \|\pi^1 \Delta_*^{u_\alpha}\|^2 + \epsilon \|H_f^{\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 \\ & + |4\alpha^{\frac{1}{2}} \langle \nabla u_\alpha \pi^1 \Delta_*^{u_\alpha}, A^- \left(\alpha \kappa_2 \Phi_*^2 \pi^0 \Psi_1^\perp + \alpha \sum_{i=1}^3 \kappa_{2,i} (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \right) \rangle|. \end{aligned}$$

To estimate the last term on the right hand side we note that $H_f^{-\frac{1}{2}} A^- \Phi_*^2 \in L^2$ and $H_f^{-\frac{1}{2}} A^- (H_f + P_f^2)^{-1} W_i \in L^2$ which thus gives for this term the bound

$$(134) \quad c\alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 + \epsilon \alpha^4 |\kappa_2|^2 \|\pi^0 \Psi_1^\perp\|^2 + \epsilon \alpha^6 \sum_i |\kappa_{2,i}|^2 = \mathcal{O}(\alpha^5),$$

using Theorem 5.2 and Lemma 6.1. The inequalities (133) and (134) imply

$$(135) \quad |\operatorname{Re} \alpha^{\frac{1}{2}} \langle P.A^+ \pi^1 u_\alpha \Delta_*^{u_\alpha}, \pi^2 \Psi_1^\perp + \pi^2 \Psi_2^\perp \rangle| \leq \epsilon \|H_f^{\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5).$$

To complete the estimate of the term $4\alpha^{\frac{1}{2}}\text{Re}\langle P.A^+u_\alpha\Delta_*^{u_\alpha}, \Psi^\perp \rangle$ we have to estimate the term for $n = 2$ in (132), namely $4\text{Re}\alpha^{\frac{1}{2}}\langle P.A^+u_\alpha\pi^2\Delta_*^{u_\alpha}, \pi^3\Psi_1^\perp + \pi^3\Psi_2^\perp \rangle$. Obviously,

$$(136) \quad |\text{Re}\alpha^{\frac{1}{2}}\langle P.A^+u_\alpha\pi^2\Delta_*^{u_\alpha}, \pi^3\Psi_2^\perp \rangle| \leq \epsilon\|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5).$$

For the term involving $\pi^3\Psi_1^\perp$ we have

$$(137) \quad \begin{aligned} & |\text{Re}\alpha^{\frac{1}{2}}\langle Pu_\alpha\pi^2\Delta_*^{u_\alpha}, \alpha\kappa_3A^-(H_f + P_f^2)^{-1}A^+.A^+\Phi_\#^{u_\alpha} \rangle| \\ & \leq c\alpha^{3-\frac{1}{16}}\|\pi^2\Delta_*^{u_\alpha}\|^2 + \epsilon|\kappa_3|^2\alpha^{2+\frac{1}{16}}\|\Phi_\#^{u_\alpha}\|^2 = |\kappa_3|^2\alpha^5 + \mathcal{O}(\alpha^5), \end{aligned}$$

using Theorem 5.2 and Lemma C.6. Collecting (132)-(137) yields

$$(138) \quad |\text{Re}\alpha^{\frac{1}{2}}\langle P.A^+u_\alpha\Delta_*^{u_\alpha}, g \rangle| \leq \epsilon\|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + \alpha^5|\kappa_3|^2 + \mathcal{O}(\alpha^5).$$

◇ *Step 2* We next estimate in (131) the term $-4\text{Re}\alpha^{\frac{1}{2}}\langle P.A^+u_\alpha(\pi^0\Psi^{u_\alpha} + \alpha^{\frac{3}{2}}2\eta_1\Phi_*^1 + \alpha\eta_2\Phi_*^2 + \alpha^{\frac{3}{2}}2\eta_3\Phi_*^3), \Psi^\perp \rangle$. First using (62) yields

$$(139) \quad -4\text{Re}\alpha^{\frac{1}{2}}\langle P.A^+u_\alpha\pi^0\Psi^{u_\alpha}, \Psi^\perp \rangle = -2\text{Re}(\overline{\kappa_1}\pi^0\Psi^{u_\alpha})\|\Phi_\#^{u_\alpha}\|_\#^2.$$

We also have, using Theorem 5.2

$$(140) \quad \begin{aligned} & |\alpha^{\frac{1}{2}}\langle Pu_\alpha(\alpha^{\frac{3}{2}}2\eta_1\Phi_*^1 + \alpha^{\frac{3}{2}}2\eta_3\Phi_*^3), A^-\Psi^\perp \rangle| \\ & \leq \alpha\|H_f^{\frac{1}{2}}\pi^2\Psi^\perp\|^2 + \alpha\|H_f^{\frac{1}{2}}\pi^4\Psi^\perp\|^2 + \mathcal{O}(\alpha^5) = \mathcal{O}(\alpha^5), \end{aligned}$$

and

$$(141) \quad \begin{aligned} & |\alpha^{\frac{1}{2}}\langle Pu_\alpha\alpha\eta_2\Phi_*^2, A^-(\pi^3\Psi_1^\perp + \pi^3\Psi_2^\perp) \rangle| \\ & \leq \epsilon\|H_f^{\frac{1}{2}}\pi^3\Psi_2^\perp\|^2 + |\alpha^{\frac{3}{2}}\langle Pu_\alpha\eta_2|k_1|^{-\frac{1}{4}}|k_2|^{-\frac{1}{4}}\Phi_*^2, |k_1|^{\frac{1}{4}}|k_2|^{\frac{1}{4}}A^-\pi^3\Psi_1^\perp \rangle| + \mathcal{O}(\alpha^5) \\ & \leq \epsilon\|H_f^{\frac{1}{2}}\pi^3\Psi_2^\perp\|^2 + \epsilon|\kappa_3|^2\alpha^5 + \mathcal{O}(\alpha^5). \end{aligned}$$

Here we used $|k_1|^{-\frac{1}{4}}|k_2|^{-\frac{1}{4}}\Phi_*^2 \in L^2$ and $\||k_1|^{\frac{1}{4}}|k_2|^{\frac{1}{4}}A^-(H_f + P_f^2)^{-1}A^+.A^+\Phi_\#^{u_\alpha}\|^2 = \mathcal{O}(\alpha^3)$ (see Lemma C.6).

Collecting (138)-(141) yields

$$(142) \quad -4\text{Re}\alpha^{\frac{1}{2}}\langle P.A^+u_\alpha\Psi^{u_\alpha}, \Psi^\perp \rangle \geq -2\text{Re}\overline{\kappa_1}\pi^0\Psi^{u_\alpha}\|\Phi_\#^{u_\alpha}\|_\#^2 - \epsilon\|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 - \epsilon\alpha^5|\kappa_3|^2 + \mathcal{O}(\alpha^5).$$

□

Lemma A.3.

$$(143) \quad \begin{aligned} & -4\sqrt{\alpha}\text{Re}\langle P.A^-u_\alpha\Psi^{u_\alpha}, \Psi^\perp \rangle \geq \\ & -\frac{2}{3}\alpha^4\text{Re}\sum_{i=1}^3\langle (H_f + P_f^2)^{-\frac{1}{2}}(A^-)^i\Phi_*^2, (H_f + P_f^2)^{-\frac{1}{2}}(A^+)^i\Omega_f \rangle \\ & -\frac{1}{4}M[\Psi_2^\perp] - \epsilon\alpha^2\|(\Psi_1^\perp)^a\|^2 - \epsilon\alpha^5\log\alpha^{-1}(\alpha|\kappa_3|^2 + 1) - |\kappa_1 - 1|c\alpha^4 + \mathcal{O}(\alpha^5), \end{aligned}$$

where $(\Psi_1^\perp)^a(x) := (\pi^0\Psi_1^\perp(x) - \pi^0\Psi_1^\perp(-x))/2$ is the odd part of $\pi^0\Psi_1^\perp$.

Proof. Since from Lemma C.2 we have $\nabla u_\alpha \cdot A^- \Phi_*^1 = 0$, we have

$$(144) \quad \begin{aligned} 4\operatorname{Re} \alpha^{\frac{1}{2}} \langle P \cdot A^- u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle &= 4\alpha^{\frac{3}{2}} \operatorname{Re} \eta_2 \langle A^- \Phi_*^2 \cdot P u_\alpha, \pi^1 \Psi^\perp \rangle \\ &+ 4\alpha^2 \operatorname{Re} 2\eta_3 \langle A^- \Phi_*^3 \cdot P u_\alpha, \pi^2 \Psi^\perp \rangle + 4\alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot P u_\alpha, \Psi^\perp \rangle. \end{aligned}$$

For the first term on the right hand side of (144) we have

$$(145) \quad \begin{aligned} &4\alpha^{\frac{3}{2}} \operatorname{Re} \eta_2 \langle A^- \Phi_*^2 \cdot P u_\alpha, \pi^1 \Psi_2^\perp + \kappa_1 \Phi_\#^{u_\alpha} \rangle \\ &= 4\alpha^{\frac{3}{2}} \operatorname{Re} \eta_2 \langle H_f^{-\frac{1}{2}} A^- \Phi_*^2 \cdot P u_\alpha, H_f^{\frac{1}{2}} \pi^1 \Psi_2^\perp \rangle + 4\alpha^{\frac{3}{2}} \operatorname{Re} \eta_2 \langle A^- \Phi_*^2 \cdot P u_\alpha, \kappa_1 \Phi_\#^{u_\alpha} \rangle. \end{aligned}$$

The first term on the right hand side of (145) is bounded from below by $-\epsilon \|H_f^{\frac{1}{2}} \pi^1 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5)$.

Applying Lemma C.5, we can replace $\Phi_\#^{u_\alpha}$ in the second term of the right hand side of (145) by $\Phi_*^{u_\alpha}$, at the expense of $\mathcal{O}(\alpha^5)$. More precisely

$$\begin{aligned} &|\alpha^{\frac{3}{2}} \langle \eta_2 A^- \Phi_*^2 \cdot P u_\alpha, \kappa_1 (\Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha}) \rangle| \\ &\leq c\alpha^5 |\eta_2|^2 \| (H_f + P_f^2)^{-\frac{1}{2}} A^- \Phi_*^2 \|^2 + |\kappa_1|^2 \| \Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha} \|^2 = \mathcal{O}(\alpha^5). \end{aligned}$$

Moreover

$$(146) \quad \begin{aligned} &4\alpha^{\frac{3}{2}} \operatorname{Re} \eta_2 \overline{\kappa_1} \langle A^- \Phi_*^2 \cdot P u_\alpha, \Phi_*^{u_\alpha} \rangle \\ &= \frac{8}{3} \alpha^2 \|\nabla u_\alpha\|^2 \operatorname{Re} \eta_2 \overline{\kappa_1} \sum_{i=1}^3 \langle (A^-)^i \Phi_*^2, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle \\ &= \frac{2}{3} \alpha^4 \operatorname{Re} \eta_2 \overline{\kappa_1} \sum_{i=1}^3 \langle (A^-)^i \Phi_*^2, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle \\ &\geq \frac{2}{3} \alpha^4 \operatorname{Re} \sum_{i=1}^3 \langle (A^-)^i \Phi_*^2, (H_f + P_f^2)^{-1} (A^+)^i \Omega_f \rangle - |\kappa_1 - 1| c\alpha^4, \end{aligned}$$

where we used $\kappa_1 = \mathcal{O}(1)$ (Lemma 6.1) and $\eta_2 = 1 + \mathcal{O}(\alpha)$ (Theorem 5.2). Note that the right hand side of (146) is well defined since $(H_f + P_f^2)^{-\frac{1}{2}} A^+ \Omega_f \in \mathfrak{F}$ and $(H_f + P_f^2)^{-\frac{1}{2}} A^- \Phi_*^2 \in \mathfrak{F}$.

Collecting the estimates for the first and the second term in the right hand side of (145), we arrive at

$$(147) \quad \begin{aligned} &-4\alpha^{\frac{3}{2}} \operatorname{Re} \eta_2 \langle A^- \Phi_*^2 \cdot P u_\alpha, \pi^1 \Psi^\perp \rangle \\ &\geq -\frac{8}{3} \alpha^2 \|\nabla u_\alpha\|^2 \operatorname{Re} \overline{\kappa_1} \langle (A^-)^i \Phi_*^2, (H_f + P_f^2)^{-1} A^+ \Omega_f \rangle - \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5). \end{aligned}$$

Here we used also $\eta_2 = 1 + \mathcal{O}(\alpha)$.

As the next step, we return to (144) and estimate the second term on the right hand side as

$$(148) \quad 4\alpha^2 \operatorname{Re} 2\eta_3 \langle A^- \Phi_*^3 \cdot P u_\alpha, \pi^2 \Psi^\perp \rangle = 8\alpha^2 \operatorname{Re} \eta_3 \langle H_f^{-\frac{1}{2}} A^- \Phi_*^3 \cdot P u_\alpha, H_f^{\frac{1}{2}} \pi^2 \Psi^\perp \rangle = \mathcal{O}(\alpha^5),$$

where we used $H_f^{-\frac{1}{2}} A^- \Phi_*^3 \in L^2$ and $\|\pi^2 H_f^{\frac{1}{2}} \Delta_*^\perp\|^2 = \|\pi^2 H_f^{\frac{1}{2}} \Psi^\perp\|^2 = \mathcal{O}(\alpha^4)$ from Theorem 5.2. For the last term on the right hand side of (144), we have

$$\begin{aligned}
 & 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \Psi^\perp \rangle \\
 (149) \quad & = 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^0 \Psi_1^\perp \rangle + 4\alpha^{\frac{1}{2}} \operatorname{Re} \overline{\kappa_1} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \Phi_\#^{u_\alpha} \rangle \\
 & + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^1 \Psi_2^\perp \rangle + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^2 \Psi^\perp \rangle \\
 & + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^3 \Psi^\perp \rangle + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^{n \geq 4} \Psi^\perp \rangle.
 \end{aligned}$$

We write the function $\pi^0 \Psi_1^\perp = (\Psi_1^\perp)^s + (\Psi_1^\perp)^a$ where $(\Psi_1^\perp)^s$ (respectively $(\Psi_1^\perp)^a$) denotes the even (respectively odd) part of $\pi^0 \Psi_1^\perp$. Obviously, we have

$$(150) \quad |\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^0 \Psi_1^\perp \rangle| \leq c\alpha \|H_f^{\frac{1}{2}} R\|^2 + \epsilon\alpha^2 \|(\Psi_1^\perp)^a\|^2 = \epsilon\alpha^2 \|(\Psi_1^\perp)^a\|^2 + \mathcal{O}(\alpha^5).$$

The constant ϵ can be chosen small for large c .

For the second term on the right hand side of (149), we have

$$(151) \quad |\alpha^{\frac{1}{2}} \overline{\kappa_1} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \Phi_\#^{u_\alpha} \rangle| \leq \epsilon\alpha^5 \log \alpha^{-1} |\kappa_1|^2 + c\alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 = \epsilon\alpha^5 \log \alpha^{-1} + \mathcal{O}(\alpha^5).$$

For the third term on the right hand side of (149), we have, since $\delta = \frac{3}{32}\alpha^2$

$$(152) \quad |4\alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^1 \Psi_2^\perp \rangle| \leq \frac{\delta}{8} \|\pi^1 \Psi_2^\perp\|^2 + c\alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 = \frac{\delta}{8} \|\pi^1 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5).$$

Similarly

$$(153) \quad |4\alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot \nabla u_\alpha, \pi^{n \geq 4} \Psi^\perp \rangle| \leq \frac{\delta}{8} \|\pi^{n \geq 4} \Psi^\perp\|^2 + \mathcal{O}(\alpha^5).$$

To complete the estimate of the last term in (149), we have to estimate two terms: $-4\operatorname{Re} \alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot P u_\alpha, \pi^2 \Psi^\perp \rangle$ and $-4\operatorname{Re} \alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot P u_\alpha, \pi^3 \Psi^\perp \rangle$. For the first one we have

$$\begin{aligned}
 & |\operatorname{Re} \alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot P u_\alpha, \pi^2 \Psi^\perp \rangle| \leq c\alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 + \epsilon\alpha^2 \|\pi^2 \Psi_2^\perp\|^2 + \epsilon\alpha^4 |\kappa_2|^2 \|\pi^0 \Psi^\perp\|^2 \\
 & + \epsilon\alpha^6 \sum_{i=1}^3 |\kappa_{2,i}|^2 = \epsilon\alpha^2 \|\pi^2 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5).
 \end{aligned}$$

Similarly,

$$(154) \quad |\operatorname{Re} \alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot P u_\alpha, \pi^3 \Psi^\perp \rangle| \leq c\alpha \|H_f^{\frac{1}{2}} \Delta_*^{u_\alpha}\|^2 + \epsilon\alpha^2 \|\pi^3 \Psi_2^\perp\|^2 + \epsilon\alpha^6 \log \alpha^{-1} |\kappa_3|^2 + \mathcal{O}(\alpha^5).$$

Collecting the estimates (149)-(154) yields

$$(155) \quad |4\operatorname{Re} \alpha^{\frac{1}{2}} \langle A^- \Delta_*^{u_\alpha} \cdot P u_\alpha, \Psi^\perp \rangle| \leq \frac{\delta}{8} \|\Psi_2^\perp\|^2 + \epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 + \epsilon\alpha^5 \log \alpha^{-1} (1 + \alpha |\kappa_3|^2) + \mathcal{O}(\alpha^5).$$

Collecting (138), (142), (147), (148) and (155) concludes the proof. \square

A.1. Concluding the proof of Proposition 6.1. We can now prove the estimate on $\text{Re} \langle Hg, u_\alpha \Psi^{u_\alpha} \rangle$ asserted in Proposition 6.1.

Using the orthogonality (39) of u_α and Ψ^\perp , yields

$$2\text{Re} \langle H\Psi^\perp, \Psi^{u_\alpha} \rangle = -4\text{Re} \langle P.P_f u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle - 4\sqrt{\alpha}\text{Re} \langle P.A(0)u_\alpha \Psi^{u_\alpha}, \Psi^\perp \rangle.$$

Together with Lemmata A.1-A.3, this concludes the proof of Proposition 6.1. \square

APPENDIX B. PROOF OF PROPOSITION 6.2

In this section, we present the proof of Proposition 6.2.

To begin with, we establish the estimate

Proposition B.1. *We have*

$$\begin{aligned} & \langle H\Psi^\perp, \Psi^\perp \rangle \\ (156) \quad & \geq \langle H\Psi_1^\perp, \Psi_1^\perp \rangle + \langle H\Psi_2^\perp, \Psi_2^\perp \rangle + 2\alpha (\|A^-\Psi^\perp\|^2 - \|A^-\Psi_1^\perp\|^2 - \|A^-\Psi_2^\perp\|^2) \\ & - 4\text{Re} \langle P.P_f \Psi_2^\perp, \Psi_1^\perp \rangle - 4\alpha^{\frac{1}{2}}\text{Re} \langle P.A(0)\Psi_2^\perp, \Psi_1^\perp \rangle + 4\text{Re} \langle P_f.A(0)\Psi_2^\perp, \Psi_1^\perp \rangle \\ & - \epsilon M[\Psi_2^\perp] - c\alpha^6 \log \alpha^{-1} |\kappa_3|^2 - c_0\alpha \|(h_\alpha + e_0)^{\frac{1}{2}} \pi^0 \Psi_1^\perp\|^2 + \mathcal{O}(\alpha^5). \end{aligned}$$

Proof. Recall that

$$(157) \quad H = P^2 - \frac{\alpha}{|x|} + T(0) - 2\text{Re} P. \left(P_f + \alpha^{\frac{1}{2}} A(0) \right),$$

and

$$(158) \quad T(0) =: (P_f + \alpha^{\frac{1}{2}} A(0))^2 : + H_f.$$

Due to the orthogonality

$$\langle \pi^n \Psi_1^\perp, \pi^n \Psi_2^\perp \rangle_\# = 0, \quad n = 0, 1, \dots,$$

and (157), (158), we obtain

$$\begin{aligned} (159) \quad & \langle (H + e_0)\Psi^\perp, \Psi^\perp \rangle \\ & = \langle (H + e_0)\Psi_2^\perp, \Psi_2^\perp \rangle + \langle (H + e_0)\Psi_1^\perp, \Psi_1^\perp \rangle + \sum_{n=0}^3 2\alpha \text{Re} \langle A^-.A^-\pi^{n+2}\Psi_2^\perp, \pi^n \Psi_1^\perp \rangle \\ & + 2\alpha \text{Re} \langle A^-.A^-\pi^3\Psi_1^\perp, \pi^1\Psi_2^\perp \rangle + 2\alpha (\|A^-\Psi^\perp\|^2 - \|A^-\Psi_1^\perp\|^2 - \|A^-\Psi_2^\perp\|^2) \\ & - 4\text{Re} \langle P.P_f \Psi_2^\perp, \Psi_1^\perp \rangle - 4\alpha^{\frac{1}{2}}\text{Re} \langle P.A(0)\Psi_2^\perp, \Psi_1^\perp \rangle + 4\alpha^{\frac{1}{2}}\text{Re} \langle P_f.A(0)\Psi_2^\perp, \Psi_1^\perp \rangle. \end{aligned}$$

We have

$$\begin{aligned} (160) \quad & 2\alpha \text{Re} \langle A^-.A^-\pi^5\Psi_2^\perp, \pi^3\Psi_1^\perp \rangle \geq -\epsilon \|H_f^{\frac{1}{2}}\pi^5\Psi_2^\perp\|^2 - c\alpha^2 \|\pi^3\Psi_1^\perp\|^2 \\ & \geq -\epsilon \|H_f^{\frac{1}{2}}\pi^5\Psi_2^\perp\|^2 - c\alpha^7 \log \alpha^{-1} |\kappa_3|^2. \end{aligned}$$

Similarly,

$$\begin{aligned} & 2\alpha \text{Re} \langle A^-.A^-\pi^4\Psi_2^\perp, \pi^2\Psi_1^\perp \rangle \\ & \geq -\epsilon \|H_f^{\frac{1}{2}}\pi^4\Psi_2^\perp\|^2 - c\alpha^4 |\kappa_2|^2 \|\pi^0\Psi_1^\perp\|^2 - \sum_{i=1}^3 c\alpha^6 |\kappa_{2,i}|^2 \\ & \geq -\epsilon \|H_f^{\frac{1}{2}}\pi^4\Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5). \end{aligned}$$

To estimate the term $2\alpha \text{Re} \langle A^- . A^- \pi^3 \Psi_2^\perp, \pi^1 \Psi_1^\perp \rangle$ we rewrite it as $2\alpha \text{Re} \langle \pi^3 \Psi_2^\perp, A^+ . A^+ \kappa_1 \Phi_\#^{u_\alpha} \rangle$ and use that $\langle \pi^3 \Psi_2^\perp, (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha} \rangle_\# = 0$. This yields, using Lemma C.6

$$(161) \quad \begin{aligned} \alpha \text{Re} \langle A^- . A^- \pi^3 \Psi_2^\perp, \pi^1 \Psi_1^\perp \rangle &= -\alpha \text{Re} \langle \pi^3 \Psi_2^\perp, (h_\alpha + e_0)(H_f + P_f^2)^{-1} A^+ . A^+ \kappa_1 \Phi_\#^{u_\alpha} \rangle \\ &\geq -\epsilon \alpha^2 \|\pi^3 \Psi_2^\perp\|^2 + c\alpha^7 \log \alpha^{-1}. \end{aligned}$$

Similarly, if $\|\pi^0 \Psi_1^\perp\| > \alpha^{\frac{3}{2}}$,

$$(162) \quad \begin{aligned} 2\alpha \text{Re} \langle A^- . A^- \pi^2 \Psi_2^\perp, \pi^0 \Psi_1^\perp \rangle &= -2\alpha \text{Re} \langle \pi^2 \Psi_2^\perp, (h_\alpha + e_0)(H_f + P_f^2)^{-1} A^+ . A^+ \pi^0 \Psi_1^\perp \rangle \\ &\geq -c\alpha \|(h_\alpha + e_0)^{\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 - c\alpha \|(h_\alpha + e_0)^{\frac{1}{2}} (H_f + P_f^2)^{-1} A^+ . A^+ \pi^0 \Psi_1^\perp\|^2 \\ &\geq -c\alpha \|P \pi^2 \Psi_2^\perp\|^2 + c\alpha \|x|^{-\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 - c\alpha e_0 \|\pi^2 \Psi_2^\perp\|^2 - c_0 \alpha \|(h_\alpha + e_0)^{\frac{1}{2}} \pi^0 \Psi_1^\perp\|^2 \\ &\geq -c\alpha \|P \pi^2 \Psi_2^\perp\|^2 - \epsilon \alpha^2 \|\pi^2 \Psi_2^\perp\|^2 - c_0 \alpha \|(h_\alpha + e_0)^{\frac{1}{2}} \pi^0 \Psi_1^\perp\|^2. \end{aligned}$$

If $\|\pi^0 \Psi_1^\perp\| \leq \alpha^{\frac{3}{2}}$, we have instead

$$(163) \quad \begin{aligned} 2\alpha \text{Re} \langle A^- . A^- \pi^2 \Psi_2^\perp, \pi^0 \Psi_1^\perp \rangle &\geq -\epsilon \|H_f^{\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 - c\alpha^2 \|\pi^0 \Psi_1^\perp\|^2 \\ &\geq -\epsilon \|H_f^{\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5). \end{aligned}$$

Finally, using Lemma C.7 yields

$$(164) \quad 2\alpha \text{Re} \langle A^- . A^- \pi^3 \Psi_1^\perp, \pi^1 \Psi_2^\perp \rangle \geq -\epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5).$$

Collecting (159)-(164) concludes the proof of the proposition. \square

In the rest of this section, we estimate further terms in (156).

B.1. Estimate of crossed terms involving Ψ_1^\perp and Ψ_2^\perp .

Lemma B.1.

$$(165) \quad 2\alpha (\|A^- \Psi^\perp\|^2 - \|A^- \Psi_1^\perp\|^2 - \|A^- \Psi_2^\perp\|^2) \geq -\epsilon M[\Psi_2^\perp] + \mathcal{O}(\alpha^5).$$

Proof. Obviously, the left hand side of (165) is equal to

$$(166) \quad \begin{aligned} 4\alpha \text{Re} \langle A^- \Psi_1^\perp, A^- \Psi_2^\perp \rangle &\geq -c\alpha \sum_n \|H_f^{\frac{1}{2}} \pi^n \Psi_1^\perp\| \|H_f^{\frac{1}{2}} \pi^n \Psi_2^\perp\| \\ &\geq -\epsilon \|H_f^{\frac{1}{2}} \pi^1 \Psi_2^\perp\|^2 - c\alpha^2 |\kappa_1|^2 \|H_f^{\frac{1}{2}} \Phi_\#^{u_\alpha}\|^2 - \sum_{n \neq 1} c\alpha \left(3\|H_f^{\frac{1}{2}} \pi^n \Psi_2^\perp\|^2 + 2\|H_f^{\frac{1}{2}} \pi^n \Psi_1^\perp\|^2 \right) \\ &\geq -\epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5), \end{aligned}$$

where in the last inequality we used (210) of Lemma C.4, and (45) of Theorem 5.2. \square

Lemma B.2.

$$|\langle P.P_f \Psi_2^\perp, \Psi_1^\perp \rangle| \leq \epsilon M[\Psi_2^\perp] + c\alpha^7 \log \alpha^{-1} |\kappa_3|^2 + \mathcal{O}(\alpha^5)$$

Proof. We have

$$(167) \quad \langle P.P_f\Psi_2^\perp, \Psi_1^\perp \rangle = \langle P_f\Pi^1\Psi_2^\perp, P\Pi^1\Psi_1^\perp \rangle + \langle P_f\Pi^2\Psi_2^\perp, P\Pi^2\Psi_1^\perp \rangle + \langle P_f\Pi^3\Psi_2^\perp, P\Pi^3\Psi_1^\perp \rangle.$$

Obviously, using Lemma C.4 and the equality $\kappa_1 = \mathcal{O}(1)$ from Lemma 6.1, yields

$$(168) \quad |\langle P_f\Pi^1\Psi_2^\perp, P\Pi^1\Psi_1^\perp \rangle| \leq \epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + |\kappa_1|^2 \|P|P_f|^{\frac{1}{2}}\Phi_\#^{u_\alpha}\|^2 \leq \epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5).$$

We also have, by definition of $\pi^2\Psi_1^\perp$ and using the estimates $\kappa_{2,i} = \mathcal{O}(1)$ from Lemma 6.1,

$$(169) \quad |\langle P_f\Pi^2\Psi_2^\perp, P\Pi^2\Psi_1^\perp \rangle| \leq \epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + c|\kappa_2|^2\alpha^2\|P\Pi^0\Psi_1^\perp\|^2 + \mathcal{O}(\alpha^6).$$

We next bound the second term on the right hand side of (169). Notice that by definition of κ_2 , this term is nonzero only if $\|\pi^0\Psi_1^\perp\|^2 > \alpha^3$, which implies, with Lemma 6.1, that $\|P\Pi^0\Psi_1^\perp\|^2 \leq c\alpha\|\pi^0\Psi_1^\perp\|^2$. The inequality (169) can thus be rewritten as

$$(170) \quad |\langle P.P_f\Pi^2\Psi_2^\perp, \pi^2\Psi_1^\perp \rangle| \leq \epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + c|\kappa_2|^2\alpha^3\|\pi^0\Psi_1^\perp\|^2 + \mathcal{O}(\alpha^6) \leq \epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5),$$

using in the last inequality that $|\kappa_2|\|\pi^0\Psi_1^\perp\| = \mathcal{O}(\alpha)$ (see Lemma 6.1).

For the second term on the right hand side of (167), using (212) from Lemma C.4 yields

$$(171) \quad \langle \pi^3 P_f\Psi_2^\perp, \pi^3 P\Psi_1^\perp \rangle \leq \epsilon \|\pi^3 H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 + c\alpha^7 \log \alpha^{-1} |\kappa_3|^2.$$

The inequalities (167), (170) and (171) prove the lemma. \square

Lemma B.3.

$$-4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A^+\Psi_2^\perp, \Psi_1^\perp \rangle \geq -\epsilon M[\Psi_2^\perp] + \mathcal{O}(\alpha^5).$$

Proof. Since $\pi^{n>3}\Psi_1^\perp = 0$, $\pi^0\Psi_2^\perp = 0$ and

$$\|H_f^{\frac{1}{2}}\pi^{n \neq 1}\Psi_1^\perp\|^2 \leq 2\|H_f^{\frac{1}{2}}\pi^{n \neq 1}\Psi^\perp\|^2 + 2\|H_f^{\frac{1}{2}}\pi^{n \neq 1}\Psi_2^\perp\|^2 \leq cM[\Psi_2^\perp] + \mathcal{O}(\alpha^4),$$

(see (45) of Theorem 5.2) we obtain

$$4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A^+\Psi_2^\perp, \Psi_1^\perp \rangle \leq \epsilon \|\pi^{n \leq 2}P\Psi_2^\perp\|^2 + c\alpha \|\pi^{n \geq 2}H_f^{\frac{1}{2}}\Psi_1^\perp\|^2 \leq \epsilon M[\Psi_2^\perp] + \mathcal{O}(\alpha^5). \quad \square$$

Lemma B.4.

$$-4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A^-\Psi_2^\perp, \Psi_1^\perp \rangle \geq -\epsilon M[\Psi_2^\perp] + \mathcal{O}(\alpha^5).$$

Proof.

$$\begin{aligned} & -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A^-\Psi_2^\perp, \Psi_1^\perp \rangle \geq -\epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 - c\alpha \|P\Psi_1^\perp\|^2 \\ & \geq -\epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 - c\alpha (\|\pi^{n=0,2,3,4}P\Psi^\perp\|^2 + \|\pi^{n=0,2,3,4}P\Psi_2^\perp\|^2) - c\alpha |\kappa_1|^2 \|P\Phi_\#^{u_\alpha}\|^2 \\ & \geq -\epsilon \|H_f^{\frac{1}{2}}\Psi_2^\perp\|^2 - \epsilon \|\pi^{n \leq 4}P\Psi_2^\perp\|^2 + \mathcal{O}(\alpha^6 \log \alpha^{-1}) \geq -\epsilon M[\Psi_2^\perp] + \mathcal{O}(\alpha^6 \log \alpha^{-1}), \end{aligned}$$

using (212) of Lemma C.4. \square

Lemma B.5.

$$4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A(0)\Psi_2^\perp, \Psi_1^\perp \rangle \geq 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle \pi^2\Psi_2^\perp, P_f.A^+\Phi_\#^{u_\alpha} \rangle - \epsilon M[\Psi_2^\perp] + \mathcal{O}(\alpha^5).$$

Proof. We have

$$\begin{aligned}
& 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^- \Psi_2^\perp, \Psi_1^\perp \rangle \\
& \geq -\epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 - c\alpha \|\pi^{n \geq 2} P_f \Psi_1^\perp\|^2 + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^- \pi^2 \Psi_2^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle \\
& \geq -\epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 - c\alpha \|\pi^{n=2,3} H_f^{\frac{1}{2}} \Psi_1^\perp\|^2 + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^- \pi^2 \Psi_2^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle \\
& \geq -\epsilon M[\Psi_2^\perp] + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^- \pi^2 \Psi_2^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle + \mathcal{O}(\alpha^5).
\end{aligned}$$

We estimate the second term on the right hand side as follows,

$$\begin{aligned}
& 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^- \pi^2 \Psi_2^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle \\
& \geq 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^- \pi^2 \Psi_2^\perp, \Phi_\#^{u_\alpha} \rangle - \epsilon \|H_f^{\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 - c\alpha |\kappa_1 - 1|^2 \|P_f \Phi_\#^{u_\alpha}\|^2 \\
& \geq 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^- \pi^2 \Psi_2^\perp, \Phi_\#^{u_\alpha} \rangle - \epsilon \|H_f^{\frac{1}{2}} \pi^2 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5),
\end{aligned}$$

where we used $|\kappa_1 - 1|^2 = \mathcal{O}(\alpha)$ from Lemma 6.1, $\|\Phi_\#^{u_\alpha}\|_*^2 = \mathcal{O}(\alpha^3)$ from Lemma C.4, and $\|P_f(\Phi_\#^{u_\alpha} - \Phi_\#^{u_\alpha})\|^2 = \mathcal{O}(\alpha^4)$ from Lemma C.5.

We also have, using $P_f . A^+ = A^+ . P_f$,

$$\begin{aligned}
& 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A^+ \Psi_2^\perp, \Psi_1^\perp \rangle = 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle \pi^{n \leq 2} P_f \Psi_2^\perp, A^- \pi^{n \geq 2} \Psi_1^\perp \rangle \\
& \geq -\epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2 - c\alpha \|H_f^{\frac{1}{2}} \pi^{n \geq 2} \Psi_1^\perp\|^2 \geq -\epsilon M[\Psi_2^\perp] + \mathcal{O}(\alpha^5),
\end{aligned}$$

where $\|H_f^{\frac{1}{2}} \pi^{n \geq 2} \Psi_1^\perp\|$ has been estimated as $\|P \Psi_1^\perp\|$ in the proof of Lemma B.4. \square

B.2. Estimates of the term $\langle (H + e_0) \Psi_1^\perp, \Psi_1^\perp \rangle$.

Due to (157) and (158), one finds

$$\begin{aligned}
(172) \quad & \langle (H + e_0) \Psi_1^\perp, \Psi_1^\perp \rangle = \langle \Psi_1^\perp, \Psi_1^\perp \rangle_\# - 2 \operatorname{Re} \langle P . (P_f + \alpha^{\frac{1}{2}} A(0)) \Psi_1^\perp, \Psi_1^\perp \rangle \\
& + 2\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f . A(0) \Psi_1^\perp, \Psi_1^\perp \rangle + 2\alpha \|A^- \Psi_1^\perp\|^2 + 2\alpha \operatorname{Re} \langle A^- . A^- \Psi_1^\perp, \Psi_1^\perp \rangle
\end{aligned}$$

We estimate the terms in (172) below.

Lemma B.6. *We have*

$$\begin{aligned}
& -2 \operatorname{Re} \langle P . (P_f + \alpha^{\frac{1}{2}} A(0)) \Psi_1^\perp, \Psi_1^\perp \rangle \\
& \geq -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^- \Phi_\#^{u_\alpha}, P \pi^0(\Psi_1^\perp)^s \rangle - \epsilon M[\Psi_2^\perp] - |\kappa_1 - 1| c \alpha^4 + \mathcal{O}(\alpha^5),
\end{aligned}$$

where $\pi^0(\Psi_1^\perp)^s = \pi^0 \Psi_1^\perp - \pi^0(\Psi_1^\perp)^a$ is the even part of $\pi^0 \Psi_1^\perp$.

Proof. Using $\langle \Phi_\#^2, P_f^i \Phi_\#^2 \rangle = 0$ (see Lemma C.2), the symmetry of u_α , and $\langle P . P_f \Phi_\#^{u_\alpha}, \Phi_\#^{u_\alpha} \rangle = 0$, we obtain

$$\begin{aligned}
(173) \quad & |2 \langle P . P_f \Psi_1^\perp, \Psi_1^\perp \rangle| = |2 \langle P_f \alpha \kappa_2 \Phi_\#^2 \pi^0 \Psi_1^\perp, P \alpha \sum_{i=1}^3 \kappa_{2,i} (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \rangle| \\
& \leq c \alpha^3 |\kappa_2|^2 \|\pi^0 \Psi_1^\perp\|^2 + c \alpha^5 \sum_{i=1}^3 |\kappa_{2,i}|^2 = \mathcal{O}(\alpha^5),
\end{aligned}$$

where in the last inequality we used Lemma 6.1.

We also have, using again the symmetry of u_α ,

$$\begin{aligned}
& -2\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A(0)\Psi_1^\perp, \Psi_1^\perp \rangle = -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P.A^\perp \Psi_1^\perp, \Psi_1^\perp \rangle \\
& = -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \kappa_1 \Phi_\#^{u_\alpha}, P\pi^0(\Psi_1^\perp)^s \rangle - 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \kappa_2 \alpha \Phi_\#^2 \pi^0 \Psi_1^\perp, P\kappa_1 \Phi_\#^{u_\alpha} \rangle \\
& - 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \pi^3 \Psi_1^\perp, P\alpha \kappa_2 \Phi_\#^2 \pi^0 \Psi_1^\perp \rangle \\
(174) \quad & \geq -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \kappa_1 \Phi_\#^{u_\alpha}, P\pi^0(\Psi_1^\perp)^s \rangle - c\alpha \|H_f^{\frac{1}{2}} \pi^3 \Psi_1^\perp\|^2 - c\alpha^2 \|P\pi^0 \Psi_1^\perp\|^2 |\kappa_2|^2 \\
& - 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle |k|^{-\frac{1}{6}} A^\perp \kappa_2 \alpha \Phi_\#^2 \pi^0 \Psi_1^\perp, |k|^{\frac{1}{6}} \kappa_1 P\Phi_\#^{u_\alpha} \rangle \\
& \geq -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \kappa_1 \Phi_\#^{u_\alpha}, P\pi^0(\Psi_1^\perp)^s \rangle - c\alpha \|H_f^{\frac{1}{2}} \pi^3 \Psi_1^\perp\|^2 - c\alpha \|H_f^{\frac{1}{2}} \pi^3 \Psi_2^\perp\|^2 \\
& - c\alpha^5 - c\alpha^{\frac{3}{2}} |\kappa_2| \|\pi^0 \Psi_1^\perp\| \| |k|^{\frac{1}{6}} P\Phi_\#^{u_\alpha} \| \\
& \geq -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \kappa_1 \Phi_\#^{u_\alpha}, P\pi^0(\Psi_1^\perp)^s \rangle - \epsilon M[\Psi_2^\perp] - c\alpha^5,
\end{aligned}$$

where we used Theorem 5.2 and Lemma C.4.

Moreover, because $\|A^\perp \Phi_\#^{u_\alpha}\| = \mathcal{O}(\alpha^{\frac{3}{2}})$ and $\|P\pi^0 \Psi_1^\perp\| = \mathcal{O}(\alpha^2)$, we obtain

$$(175) \quad -4\operatorname{Re} \alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \kappa_1 \Phi_\#^{u_\alpha}, P\pi^0(\Psi_1^\perp)^s \rangle \geq -4\alpha^{\frac{1}{2}} \operatorname{Re} \langle A^\perp \Phi_\#^{u_\alpha}, P\pi^0(\Psi_1^\perp)^s \rangle - |\kappa_1 - 1|c\alpha^4.$$

This estimate, together with (173) and (174), proves the lemma. \square

Lemma B.7. *We have*

$$\begin{aligned}
& 2\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A(0)\Psi_1^\perp, \Psi_1^\perp \rangle \\
& \geq \frac{2}{3} \operatorname{Re} \sum_i \kappa_{2,i} \langle P_f.A^\perp (H_f + P_f^2)^{-1} W_i, (H_f + P_f^2)^{-1} (A^\perp)^i \Omega_f \rangle \\
& - |\kappa_1 - 1|c\alpha^4 + \mathcal{O}(\alpha^5).
\end{aligned}$$

Proof. The following holds

$$\begin{aligned}
(176) \quad & 2\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A(0)\Psi_1^\perp, \Psi_1^\perp \rangle = 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A^\perp \pi^3 \Psi_1^\perp, \pi^2 \Psi_1^\perp \rangle + 4\alpha \operatorname{Re} \langle P_f.A^\perp \pi^2 \Psi_1^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle \\
& = 4\alpha^{\frac{3}{2}} \operatorname{Re} \sum_{i=1}^3 \overline{\kappa_{2,i}} \langle |k_1|^{\frac{1}{6}} |k_2|^{\frac{1}{6}} A^\perp \pi^3 \Psi_1^\perp, P_f |k_1|^{-\frac{1}{6}} |k_2|^{-\frac{1}{6}} (H_f + P_f^2)^{-1} W_i \frac{\partial u_\alpha}{\partial x_i} \rangle \\
& + 4\alpha^{\frac{3}{2}} \operatorname{Re} \kappa_2 \langle |k_1|^{\frac{1}{6}} |k_2|^{\frac{1}{6}} A^\perp \pi^3 \Psi_1^\perp, |k_1|^{-\frac{1}{6}} |k_2|^{-\frac{1}{6}} P_f \Phi_\#^2 \pi^0 \Psi_1^\perp \rangle \\
& + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A^\perp \pi^2 \Psi_1^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle.
\end{aligned}$$

Applying the Schwarz inequality and the estimates $\||k_1|^{\frac{1}{6}} |k_2|^{\frac{1}{6}} A^\perp \pi^3 \Psi_1^\perp\|^2 = \mathcal{O}(\alpha^5)$ (Lemma C.6), $\kappa_{2,i} = \mathcal{O}(1)$ (Lemma 6.1), and $\|\nabla u_\alpha\|^2 = \mathcal{O}(\alpha^2)$, we see that the first term on the right hand side of (176) is $\mathcal{O}(\alpha^5)$. Applying also the estimate $|\kappa_2| \|\pi^0 \Psi_1^\perp\| = \mathcal{O}(\alpha)$ (Lemma 6.1), we obtain that the second term on the right hand side of (176) is also $\mathcal{O}(\alpha^5)$.

Finally, we estimate $4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A^\perp \pi^2 \Psi_1^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle$. The following inequality holds,

$$(177) \quad 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A^\perp \pi^2 \Psi_1^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle \geq 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A^\perp \pi^2 \Psi_1^\perp, \Phi_\#^{u_\alpha} \rangle - |\kappa_1 - 1|c\alpha^4,$$

whose proof is similar to the one of (175). Next we get

$$\begin{aligned}
(178) \quad & |4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A^\perp \pi^2 \Psi_1^\perp, \Phi_\#^{u_\alpha} \rangle - 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle P_f.A^\perp \pi^2 \Psi_1^\perp, \Phi_\#^{u_\alpha} \rangle| \\
& \leq \alpha^{\frac{1}{2}} \|H_f^{\frac{1}{2}} \pi^2 \Psi_1^\perp\| \|P_f(\Phi_\#^{u_\alpha} - \Phi_\#^{u_\alpha})\| = \mathcal{O}(\alpha^6 |\log \alpha|^{\frac{1}{2}}),
\end{aligned}$$

using $\|P_f(\Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha})\| = \mathcal{O}(\alpha^{\frac{7}{2}}|\log \alpha|^{\frac{1}{2}})$ (see Lemma C.5) and $\|H_f^{\frac{1}{2}}\pi^2\Psi_1^\perp\| = \mathcal{O}(\alpha^2)$. Moreover,

$$\begin{aligned}
& 4\alpha^{\frac{1}{2}}\text{Re}\langle P_f.A^-\pi^2\Psi_1^\perp, \Phi_*^{u_\alpha} \rangle \\
&= 4\alpha^{\frac{3}{2}}\text{Re}\langle P_f.A^-(\kappa_2\Phi_*^2\pi^0\Psi_1^\perp + \sum_i \kappa_{2,i}(H_f + P_f^2)^{-1}W_i \frac{\partial u_\alpha}{\partial x_i}), \Phi_*^{u_\alpha} \rangle \\
(179) \quad &= 4\alpha^{\frac{3}{2}}\text{Re}\langle P_f.A^-\sum_i \kappa_{2,i}(H_f + P_f^2)^{-1}W_i \frac{\partial u_\alpha}{\partial x_i}, \Phi_*^{u_\alpha} \rangle \\
&= \frac{2}{3}\alpha^4\text{Re}\sum_i \kappa_{2,i}\langle P_f.A^-(H_f + P_f^2)^{-1}W_i, (H_f + P_f^2)^{-1}(A^+)^i\Omega_f \rangle
\end{aligned}$$

where we used (207) of Lemma C.2 in the second equality.

Collecting (176)-(199) concludes the proof. \square

Lemma B.8. *We have*

$$\begin{aligned}
& \langle \Psi_1^\perp, \Psi_1^\perp \rangle_\# + 2\alpha\text{Re}\langle A^-.A^-\Psi_1^\perp, \Psi_1^\perp \rangle \geq \Sigma_0(\|\pi^0\Psi_1^\perp\|^2 + \|\pi^1\Psi_1^\perp\|^2) \\
& + \frac{\alpha^4}{12}\sum_{i=1}^3|\kappa_{2,i}|^2\|(H_f + P_f^2)^{-1}W_i\|_*^2 + \alpha^2|\kappa_3 + 1|^2\|\Phi_*^2\|_*^2\|\Phi_\#^{u_\alpha}\|^2 \\
& + \|\pi^1\Psi_1^\perp\|_\#^2 + \|\pi^0\Psi_1^\perp\|_\#^2 + o(\alpha^5\log \alpha^{-1}).
\end{aligned}$$

Proof. Obviously we have

$$(180) \quad \langle \Psi_1^\perp, \Psi_1^\perp \rangle_\# = \sum_{i=0}^3 \langle \pi^i\Psi_1^\perp, \pi^i\Psi_1^\perp \rangle_\#,$$

and using Lemma C.2

$$(181) \quad \langle \pi^2\Psi_1^\perp, \pi^2\Psi_1^\perp \rangle_\# = \alpha^2|\kappa_2|^2\|\Phi_*^2\pi^0\Psi_1^\perp\|_\#^2 + \alpha^2\sum_{i=1}^3|\kappa_{2,i}|^2\|(H_f + P_f^2)^{-1}W_i \frac{\partial u_\alpha}{\partial x_i}\|_\#^2,$$

Moreover, from the inequality $\|\Phi_*^2\pi^0\Psi_1^\perp\|_\#^2 > \|\Phi_*^2\|_*^2\|\pi^0\Psi_1^\perp\|^2$ we obtain,

$$\begin{aligned}
& \alpha^2|\kappa_2|^2\|\Phi_*^2\pi^0\Psi_1^\perp\|_\#^2 + 2\alpha\text{Re}\langle A^-.A^-\pi^2\Psi_1^\perp, \pi^0\Psi_1^\perp \rangle \\
& \geq \alpha^2|\kappa_2|^2\|\Phi_*^2\|_*^2\|\pi^0\Psi_1^\perp\|^2 + 2\alpha\text{Re}\langle A^-.A^-\pi^2\Psi_1^\perp, \pi^0\Psi_1^\perp \rangle \\
& = \alpha^2|\kappa_2|^2\|\Phi_*^2\|_*^2\|\pi^0\Psi_1^\perp\|^2 - 2\alpha^2\text{Re}\kappa_2\|\Phi_*^2\|_*^2\|\pi^0\Psi_1^\perp\|^2 \\
(182) \quad & + 2\alpha\text{Re}\sum_{i=1}^3\alpha\langle A^-.A^-\kappa_{2,i}(H_f + P_f^2)^{-1}W_i \frac{\partial u_\alpha}{\partial x_i}, \pi^0\Psi_1^\perp \rangle \\
& \geq \Sigma_0\|\pi^0\Psi_1^\perp\|^2 + |\kappa_2 - 1|^2\alpha^2\|\pi^0\Psi_1^\perp\|^2 + \mathcal{O}(\alpha^5) \geq \Sigma_0\|\pi^0\Psi_1^\perp\|^2 + \mathcal{O}(\alpha^5).
\end{aligned}$$

where we used $\Sigma_0 = -\alpha^2\|\Phi_*^2\|_*^2 + \mathcal{O}(\alpha^3)$ and $\langle A^-.A^-(H_f + P_f^2)^{-1}W_i, \pi^0\Psi_1^\perp \rangle = -\overline{\pi^0\Psi_1^\perp}\langle W_i, \Phi_*^2 \rangle = 0$.

Similarly $\|\pi^3\Psi_1^\perp\|_\# > \|\pi^3\Psi_1^\perp\|_*$ yields

(183)

$$\begin{aligned} & \|\pi^3\Psi_1^\perp\|_\#^2 + 2\alpha\text{Re}\langle A^-.A^-\pi^3\Psi_1^\perp, \pi^1\Psi_1^\perp \rangle \geq \alpha^2|\kappa_3|^2\|(H_f + P_f^2)^{-1}A^+.A^+\Phi_\#^{u_\alpha}\|_*^2 \\ & + 2\text{Re}\alpha^2\kappa_3\langle A^-.A^-(H_f + P_f^2)^{-1}A^+.A^+\Phi_\#^{u_\alpha}, \kappa_1\Phi_\#^{u_\alpha} \rangle \\ & \geq -\alpha^2\|(H_f + P_f^2)^{-1}A^+.A^+\kappa_1\Phi_\#^{u_\alpha}\|_*^2 + |\kappa_3 + 1|^2\alpha^2\|(H_f + P_f^2)^{-1}A^+.A^+\Phi_\#^{u_\alpha}\|_*^2 \\ & \geq \Sigma_0\|\pi^1\Psi_1^\perp\|^2 + \alpha^2|\kappa_3 + 1|^2\|\Phi_\#^2\|_*^2\|\Phi_\#^{u_\alpha}\|^2 + o(\alpha^5\log\alpha^{-1}), \end{aligned}$$

where in the last inequality, we used (222) from Lemma C.6, $\kappa_1 = 1 + \mathcal{O}(\alpha^{\frac{1}{2}})$ from Lemma 6.1, $\|\Phi_\#^{u_\alpha}\|^2 = \mathcal{O}(\alpha^3\log\alpha^{-1})$ from Lemma C.4 and $-\alpha^2\|\Phi_\#^2\|_*^2 = \Sigma_0 + \mathcal{O}(\alpha^3)$.

The second term on the right hand side of (181) is estimated as

(184)

$$\begin{aligned} & \alpha^2 \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1}W_i \frac{\partial u_\alpha}{\partial x_i}\|_\#^2 \geq \alpha^2 \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1}W_i \frac{\partial u_\alpha}{\partial x_i}\|_*^2 \\ & = \alpha^2 \frac{1}{3} \|\nabla u_\alpha\|^2 \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1}W_i\|_*^2 = \frac{\alpha^4}{12} \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1}W_i\|_*^2. \end{aligned}$$

Collecting (180)-(184) concludes the proof. \square

Proposition B.2. *We have*

$$\begin{aligned} & \langle (H + e_0)\Psi_1^\perp, \Psi_1^\perp \rangle \\ & \geq -4\alpha\|(h_\alpha + e_0)^{-\frac{1}{2}}Q_\alpha^\perp P.A^-\Phi_\#^{u_\alpha}\|^2 + \|(h_\alpha + e_0)^{\frac{1}{2}}\pi^0(\Psi_1^\perp)^a\|^2 \\ & + c_0\alpha\|(h_\alpha + e_0)^{\frac{1}{2}}\pi^0(\Psi_1^\perp)^s\|^2 + \Sigma_0(\|\pi^0\Psi_1^\perp\|^2 + \|\pi^1\Psi_1^\perp\|^2) \\ (185) \quad & + \frac{\alpha^4}{12} \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1}W_i\|_*^2 + 4\alpha^{\frac{1}{2}}\text{Re}\langle \pi^2\Psi_1^\perp, A^+.P_f\Phi_\#^{u_\alpha} \rangle \\ & + \|\kappa_1\Phi_\#^{u_\alpha}\|_\#^2 + 2\alpha\|A^-\Phi_\#^{u_\alpha}\|^2 + \alpha^2|\kappa_3 + 1|^2\|\Phi_\#^2\|_*^2\|\Phi_\#^{u_\alpha}\|^2 \\ & - \epsilon M[\Psi_2^\perp] - |\kappa_1 - 1|c\alpha^4 + o(\alpha^5\log\alpha^{-1}), \end{aligned}$$

where c_0 is the same positive constant as in Proposition B.1 and

- Q_α^\perp is the orthogonal projection onto $\text{Span}(u_\alpha)^\perp$,
- $(\Psi_1^\perp)^a$ is the odd part of $\pi^0\Psi_1^\perp$.

Proof. Collecting Lemmata B.6, B.7, and B.8 yields

(186)

$$\begin{aligned} & \langle (H + e_0)\Psi_1^\perp, \Psi_1^\perp \rangle \geq -4\alpha^{\frac{1}{2}}\text{Re}\langle A^-\Phi_\#^{u_\alpha}, P\pi^0(\Psi_1^\perp)^s \rangle + 4\alpha\text{Re}\langle \pi^2\Psi_1^\perp, A^+.P_f\Phi_\#^{u_\alpha} \rangle \\ & + \Sigma_0(\|\pi^0\Psi_1^\perp\|^2 + \|\pi^1\Psi_1^\perp\|^2) + \alpha^2|\kappa_3 + 1|^2\|\Phi_\#^2\|_*^2\|\Phi_\#^{u_\alpha}\|^2 \\ & + \|(h_\alpha + e_0)^{\frac{1}{2}}\pi^0\Psi_1^\perp\|^2 + \|\pi^1\Psi_1^\perp\|_\#^2 + 2\alpha\|A^-\kappa_1\Phi_\#^{u_\alpha}\|^2 - |\kappa_1 - 1|c\alpha^4 \\ & + \frac{\alpha^4}{12} \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1}W_i\|_*^2 - \epsilon M[\Psi_2^\perp] + o(\alpha^5\log\alpha^{-1}). \end{aligned}$$

Obviously,

$$(187) \quad \|(h_\alpha + e_0)^{\frac{1}{2}}\pi^0\Psi_1^\perp\|^2 = \|(h_\alpha + e_0)^{\frac{1}{2}}\pi^0(\Psi_1^\perp)^s\|^2 + \|(h_\alpha + e_0)^{\frac{1}{2}}\pi^0(\Psi_1^\perp)^a\|^2.$$

As before, we write $\pi^0 \Psi_1^\perp$ as the sum of its odd part $(\Psi_1^\perp)^a$ and its even part $(\Psi_1^\perp)^s$. Since $\pi^0 \Psi_1^\perp$ is orthogonal to u_α by definition of Ψ^\perp , and $(\Psi_1^\perp)^a$ is orthogonal to u_α by symmetry of u_α , we also have $(\Psi_1^\perp)^s$ orthogonal to u_α . Therefore, one can replace $\pi^0(\Psi_1^\perp)^s$ by $Q_\alpha^\perp \pi^0(\Psi_1^\perp)^s$ in (186) and (187). Thus, as the next step, given a constant $c_0 > 0$, we minimize

$$(188) \quad \begin{aligned} & (1 - c_0 \alpha) \|(h_\alpha + e_0)^{\frac{1}{2}} Q_\alpha^\perp \pi^0(\Psi_1^\perp)^s\|^2 - 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle Q_\alpha^\perp P A^- \Phi_\#^{u_\alpha}, Q_\alpha^\perp \pi^0(\Psi_1^\perp)^s \rangle \\ & \geq -\frac{4\alpha}{1 - c_0 \alpha} \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_\#^{u_\alpha}\|^2. \end{aligned}$$

Obviously,

$$(189) \quad \begin{aligned} & -\frac{4\alpha}{1 - c_0 \alpha} \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_\#^{u_\alpha}\|^2 \\ & \geq -\frac{4(1 + \alpha)\alpha}{1 - c_0 \alpha} \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 \\ & \quad - \frac{4(1 + \alpha^{-1})\alpha}{1 - c_0 \alpha} \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- (\Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha})\|^2. \end{aligned}$$

There exist γ_1 and γ_2 positive, independent of α , such that

$$Q_\alpha^\perp (h_\alpha + e_0)^{-1} Q_\alpha^\perp \leq (\gamma_1 P^2 + \gamma_2 \alpha^2)^{-1},$$

and thus $P Q_\alpha^\perp (h_\alpha + e_0)^{-1} Q_\alpha^\perp P$ is a bounded operator. In addition, since $\|A^- (\Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha})\|^2 = \mathcal{O}(\alpha^7 \log \alpha^{-1})$ (Lemma C.4), this shows that

$$(190) \quad \frac{4(1 + \alpha^{-1})\alpha}{1 - c_0 \alpha} \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- (\Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha})\|^2 = \mathcal{O}(\alpha^6 \log \alpha^{-1}).$$

In addition, using $\|A^- \Phi_*^{u_\alpha}\| \leq \|A^- (\Phi_\#^{u_\alpha} - \Phi_*^{u_\alpha})\| + \|A^- \Phi_\#^{u_\alpha}\| \leq c\alpha^{\frac{3}{2}}$ (Lemma C.4 and Lemma C.5), and the fact that $P Q_\alpha^\perp (h_\alpha + e_0)^{-1} Q_\alpha^\perp P$ is bounded, yields

$$(191) \quad \begin{aligned} & -\frac{(1 + \alpha)\alpha}{1 - c_0 \alpha} \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 \\ & = -4 \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 + \mathcal{O}(\alpha^5). \end{aligned}$$

Collecting (189)-(191), one gets

$$(192) \quad -\frac{4\alpha}{1 - c_0 \alpha} \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_\#^{u_\alpha}\|^2 \geq -4 \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 + \mathcal{O}(\alpha^5).$$

Finally, using $\|A^- \Phi_*^{u_\alpha}\|^2 = \mathcal{O}(\alpha^3)$, we obtain

$$(193) \quad 2\alpha |\kappa_1|^2 \|A^- \Phi_*^{u_\alpha}\|^2 \geq 2\alpha \|A^- \Phi_*^{u_\alpha}\|^2 - c |\kappa_1 - 1| \alpha^4.$$

Substituting (187), (188), (192) and (193) into (186) concludes the proof. \square

We can now collect the above results to prove Proposition 6.2.

B.3. Concluding the proof of Proposition 6.2. Collecting the results of Proposition B.1, Lemmata B.1-B.5 and Proposition B.2 yields directly the following bound,

$$\begin{aligned}
(194) \quad \langle H\Psi^\perp, \Psi^\perp \rangle &\geq \langle H\Psi_2^\perp, \Psi_2^\perp \rangle - \epsilon M[\Psi_2^\perp] - c|\kappa_2|^2 \alpha^4 \|\pi^0 \Psi_1^\perp\|^2 - c\alpha^6 \log \alpha^{-1} |\kappa_3|^2 - e_0 \|\Psi_1^\perp\|^2 \\
&\quad - 4\alpha \|(h_\alpha + e_0)^{-\frac{1}{2}} Q_\alpha^\perp P A^- \Phi_*^{u_\alpha}\|^2 + (1 - c_0 \alpha) \|(h_\alpha + e_0)^{\frac{1}{2}} \pi^0 (\Psi_1^\perp)^a\|^2 \\
&\quad + \Sigma_0 (\|\pi^0 \Psi_1^\perp\|^2 + \|\pi^1 \Psi_1^\perp\|^2) + \frac{\alpha^4}{12} \sum_{i=1}^3 |\kappa_{2,i}|^2 \|(H_f + P_f^2)^{-1} W_i\|_*^2 \\
&\quad + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle \pi^2 \Psi_2^\perp, A^+ P_f \Phi_*^{u_\alpha} \rangle + 4\alpha^{\frac{1}{2}} \operatorname{Re} \langle \pi^2 \Psi_1^\perp, A^+ P_f \Phi_\#^{u_\alpha} \rangle + |\kappa_1|^2 \|\Phi_\#^{u_\alpha}\|_\#^2 + 2\alpha \|A^- \Phi_*^{u_\alpha}\|^2 \\
&\quad + \alpha^2 |\kappa_3 + 1|^2 \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 - |\kappa_1 - 1| c \alpha^4 + o(\alpha^5 \log \alpha^{-1}).
\end{aligned}$$

Comparing this expression with the statement of the Proposition we see that it suffices to show that

$$\begin{aligned}
(195) \quad & - \epsilon M[\Psi_2^\perp] - c|\kappa_2|^2 \alpha^4 \|\pi^0 \Psi_1^\perp\|^2 - c\alpha^6 \log \alpha^{-1} |\kappa_3|^2 - e_0 \|\Psi_1^\perp\|^2 + \langle H\Psi_2^\perp, \Psi_2^\perp \rangle \\
& + \Sigma_0 (\|\pi^0 \Psi_1^\perp\|^2 + \|\pi^1 \Psi_1^\perp\|^2) + \alpha^2 |\kappa_3 + 1|^2 \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 \\
& \geq (\Sigma_0 - e_0) \|\Psi^\perp\|^2 + (1 - \epsilon) M[\Psi_2^\perp] + \frac{|\kappa_3 + 1|^2}{2} \alpha^2 \|\Phi_*^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 + \mathcal{O}(\alpha^5).
\end{aligned}$$

Using from Corollary 4.2 that $\langle H\Psi_2^\perp, \Psi_2^\perp \rangle \geq (\Sigma_0 - e_0) \|\Psi_2^\perp\|^2 + M[\Psi_2^\perp]$, we first estimate the following terms in (195),

$$\begin{aligned}
(196) \quad & \Sigma_0 (\|\pi^0 \Psi_1^\perp\|^2 + \|\pi^1 \Psi_1^\perp\|^2) - e_0 \|\Psi_1^\perp\|^2 + \langle H\Psi_2^\perp, \Psi_2^\perp \rangle - \epsilon M[\Psi_2^\perp] \\
& \geq (1 - \epsilon) M[\Psi_2^\perp] + (\Sigma_0 - e_0) (\|\Psi_1^\perp\|^2 + \|\Psi_2^\perp\|^2) - \Sigma_0 \|\pi^{n \geq 2} \Psi_1^\perp\|^2 \\
& \geq (1 - \epsilon) M[\Psi_2^\perp] + (\Sigma_0 - e_0) \|\Psi^\perp\|^2 \\
& \quad - (\Sigma_0 - e_0) (\|\Psi^\perp\|^2 - \|\Psi_1^\perp\|^2 - \|\Psi_2^\perp\|^2) - \Sigma_0 \|\pi^{n \geq 2} \Psi_1^\perp\|^2.
\end{aligned}$$

We have obviously

$$(197) \quad \|\Psi^\perp\|^2 - \|\Psi_1^\perp\|^2 - \|\Psi_2^\perp\|^2 = 2\operatorname{Re} (\langle \pi^1 \Psi_1^\perp, \pi^1 \Psi_2^\perp \rangle + \langle \pi^2 \Psi_1^\perp, \pi^2 \Psi_2^\perp \rangle + \langle \pi^3 \Psi_1^\perp, \pi^3 \Psi_2^\perp \rangle).$$

Since $|\Sigma_0 - e_0| \leq c\alpha^2$, by definition of $\pi^3 \Psi_1^\perp$ and Lemma C.6, we obtain

$$\begin{aligned}
(198) \quad & |(\Sigma_0 - e_0) 2\operatorname{Re} \langle \pi^3 \Psi_1^\perp, \pi^3 \Psi_2^\perp \rangle| \leq \epsilon \alpha^2 \|\pi^3 \Psi_2^\perp\|^2 + c\alpha^4 |\kappa_3|^2 \|(H_f + P_f^2)^{-1} A^+ A^- \Phi_\#^{u_\alpha}\|^2 \\
& \leq \epsilon \alpha^2 \|\pi^3 \Psi_2^\perp\|^2 + c|\kappa_3|^2 \alpha^7 \log \alpha^{-1}.
\end{aligned}$$

Similarly, for the two-photon sector, we find

$$\begin{aligned}
(199) \quad & |(\Sigma_0 - e_0) \langle \pi^2 \Psi_1^\perp, \pi^2 \Psi_2^\perp \rangle| \leq \epsilon \alpha^2 \|\pi^2 \Psi_2^\perp\|^2 + c\alpha^4 |\kappa_2|^2 \|\pi^0 \Psi_1^\perp\|^2 + c \sum_{i=1}^3 |\kappa_{2,i}|^2 \alpha^6 \\
& = \epsilon \alpha^2 \|\pi^2 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^6),
\end{aligned}$$

where we used Lemma 6.1. For the term $\langle \pi^1 \Psi_1^\perp, \pi^1 \Psi_2^\perp \rangle$, one gets

$$(200) \quad \begin{aligned} |(\Sigma_0 - e_0) \langle \pi^1 \Psi_2^\perp, \kappa_1 \Phi_\#^{u_\alpha} \rangle| &= |(\Sigma_0 - e_0) (|k|^{\frac{1}{2}} \pi^1 \Psi_2^\perp, |k|^{-\frac{1}{2}} \kappa_1 \Phi_\#^{u_\alpha})| \\ &\leq \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Psi_2^\perp\|^2 + c\alpha^4 |\kappa_1|^2 \| |k|^{-\frac{1}{2}} \Phi_\#^{u_\alpha} \|^2 \leq \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Psi_2^\perp\|^2 + \mathcal{O}(\alpha^5), \end{aligned}$$

since $\| |k|^{-\frac{1}{2}} \Phi_\#^{u_\alpha} \|^2 = \mathcal{O}(\alpha)$ (Lemma C.4) and $\kappa_1 = \mathcal{O}(1)$ (Lemma 6.1).

Collecting (197), (198), (199) and (200) yields

$$(201) \quad |\Sigma_0 - e_0| \left| \|\Psi^\perp\|^2 - \|\Psi_1^\perp\|^2 - \|\Psi_2^\perp\|^2 \right| \leq \epsilon M[\Psi_2^\perp] + c|\kappa_3|^2 \alpha^7 \log \alpha^{-1} + \mathcal{O}(\alpha^5).$$

Therefore, together with (196), one finds

$$(202) \quad \begin{aligned} \Sigma_0 (\|\pi^0 \Psi_1^\perp\|^2 + \|\pi^1 \Psi_1^\perp\|^2) - e_0 \|\Psi_1^\perp\|^2 + \langle H \Psi_2^\perp, \Psi_2^\perp \rangle - \epsilon M[\Psi_2^\perp] \\ \geq (1 - 2\epsilon) M[\Psi_2^\perp] + (\Sigma_0 - e_0) \|\Psi^\perp\|^2 - c|\kappa_3|^2 \alpha^7 \log \alpha^{-1} - \Sigma_0 \|\pi^{n \geq 2} \Psi_1^\perp\|^2 + \mathcal{O}(\alpha^5), \end{aligned}$$

By definition of Ψ_1^\perp and using $\Sigma_0 = \mathcal{O}(\alpha^2)$, $|\kappa_2| \|\pi^0 \Psi_1^\perp\| = \mathcal{O}(\alpha)$ (Lemma 6.1) and Inequality (223) of Lemma C.6, we straightforwardly obtain

$$\Sigma_0 \|\pi^{n \geq 2} \Psi_1^\perp\|^2 \leq c\alpha^6 + c|\kappa_3|^2 \alpha^7 \log \alpha^{-1}.$$

Substituting this in (202) yields

$$(203) \quad \begin{aligned} \Sigma_0 (\|\pi^0 \Psi_1^\perp\|^2 + \|\pi^1 \Psi_1^\perp\|^2) - e_0 \|\Psi_1^\perp\|^2 + \langle H \Psi_2^\perp, \Psi_2^\perp \rangle - \epsilon M[\Psi_2^\perp] \\ \geq (1 - 2\epsilon) M[\Psi_2^\perp] + (\Sigma_0 - e_0) \|\Psi^\perp\|^2 - 2c|\kappa_3|^2 \alpha^7 \log \alpha^{-1} + \mathcal{O}(\alpha^5). \end{aligned}$$

To conclude the proof of (195), and thus of the Proposition, we first note that according to Lemma 6.1,

$$(204) \quad -c|\kappa_2|^2 \alpha^4 \|\pi^0 \Psi_1^\perp\|^2 = \mathcal{O}(\alpha^6).$$

Similarly, taking into account that $\|\Phi_\#^{u_\alpha}\|^2 = c\alpha^3 \log \alpha^{-1}$ (see (209) in Lemma C.4), we get for some $c_2 > 0$,

$$(205) \quad \alpha^2 \frac{|\kappa_3 + 1|^2}{2} \|\Phi_*^2\|^2 \|\Phi_\#^{u_\alpha}\|^2 - c\alpha^6 \log \alpha^{-1} |\kappa_3|^2 - 2c\alpha^7 \log \alpha^{-1} |\kappa_3|^2 \geq -c_2 \alpha^6 \log \alpha^{-1}.$$

Collecting (203), (204) and (205) yields the bound (195), and thus concludes the proof of the Proposition 6.2. \square

APPENDIX C. PROOF OF THEOREM 2.1: AUXILIARY LEMMATA

The proof of Theorem 2.1 uses Lemmata C.2 \sim C.7 below.

Lemma C.1. *Let Ψ_0 be the ground state of $T(0)$, with $|\pi^0 \Psi_0| = 1$. Then we have*

$$\begin{aligned} \Psi_0 &= \Omega_f + 2\eta_1 \alpha^{\frac{3}{2}} \Phi_*^1 + \eta_2 \alpha \Phi_*^2 + 2\eta_3 \alpha^{\frac{3}{2}} \Phi_*^3 + \Delta_*^0, \\ \text{with } \langle \Delta_*^0, \Phi_*^i \rangle_* &= 0 \ (i = 1, 2, 3), \pi^0 \Delta_*^0 = 0 \end{aligned}$$

and

$$\|\Delta_*^0\|^2 = \mathcal{O}(\alpha^3).$$

Proof. It follows from a similar argument as in [9, Proposition 5.1] that

$$\|a_\lambda(k) \Delta_*^0\| \leq \frac{c\alpha}{|k|}.$$

This yields

$$\begin{aligned} \|\Delta_*^0\|^2 &\leq \int \|a_\lambda(k)\Delta_*^0\|^2 dk \leq \int_{|k|\leq\alpha} \frac{c^2\alpha^2}{|k|^2} dk + \int_{|k|>\alpha} \frac{|k| \|a_\lambda(k)\Delta_*^0\|^2 \chi_\Lambda(|k|)}{|k|} dk \\ &\leq c^2\alpha^3 + c'\alpha^{-1} \|H_f^{\frac{1}{2}}\Delta_*^0\|^2 = \mathcal{O}(\alpha^3), \end{aligned}$$

where in the last equality we used the estimate $\|\Delta_*^0\|_*^2 = \|(H_f + P_f^2)^{\frac{1}{2}}\Delta_*^0\|^2 = \mathcal{O}(\alpha^4)$ proved in [7, Theorem 3.2]. \square

Lemma C.2. *We have*

$$P.A^-u_\alpha\Phi_*^1 = 0,$$

and

$$(206) \quad \langle \Phi_*^2, \zeta(H_f, P_f^2)P_f^i\Phi_*^2 \rangle = 0, \text{ and } \langle \Phi_*^2, \zeta(H_f, P_f^2)W_i \rangle = 0,$$

for any function ζ for which the scalar products are defined. Similarly, we have

$$(207) \quad \langle \Phi_*^2\pi^0\Psi_1^\perp, A^+.P_f\Phi_*^{u_\alpha} \rangle = 0.$$

Proof. Straightforward computations using the symmetries of $A^-\Phi_*^1$ and Φ_*^2 . \square

Lemma C.3. *We have*

$$P.A^-\Phi_*^{u_\alpha} = \sqrt{\alpha}a_0\Delta u_\alpha,$$

where

$$a_0 = \int \frac{k_1^2 + k_2^2}{4\pi^2|k|^3} \frac{2}{|k|^2 + |k|} \chi_\Lambda(|k|) k_1 k_2 k_3 dk.$$

Proof. Straightforward computations using the symmetries of $A^-\Phi_*^{u_\alpha}$. \square

Lemma C.4.

$$(208) \quad \forall \varphi \in \mathfrak{F}, \quad \langle P.P_f\pi^1\varphi u_\alpha, \Phi_\#^{u_\alpha} \rangle = 0,$$

$$(209) \quad \|\Phi_\#^{u_\alpha}\|^2 = \mathcal{O}(\alpha^3 \log \alpha^{-1}),$$

$$(210) \quad \|\Phi_\#^{u_\alpha}\|_*^2 = \mathcal{O}(\alpha^3),$$

$$(211) \quad \| |k|^{-\frac{1}{2}} \Phi_\#^{u_\alpha} \|^2 = \mathcal{O}(\alpha).$$

$$(212) \quad \|P\Phi_\#^{u_\alpha}\|^2 = \mathcal{O}(\alpha^5 \log \alpha^{-1}),$$

$$(213) \quad \|P\Phi_\#^{u_\alpha}\|_*^2 = \mathcal{O}(\alpha^5),$$

$$(214) \quad \| |k|^{\frac{1}{6}} P\Phi_\#^{u_\alpha} \|^2 = \mathcal{O}(\alpha^5).$$

Proof. The proof of (208) is as follows

$$\begin{aligned} (215) \quad &\langle P.P_f\pi^1\varphi u_\alpha, \Phi_\#^{u_\alpha} \rangle \\ &= \int \sum_{i=1}^3 k^i \frac{\partial u_\alpha}{\partial x_i} \varphi(k) \frac{1}{k^2 + |k| + h_\alpha + e_0} \sum_{j=1}^3 \sum_{\lambda=1,2} \frac{\epsilon_\lambda^j(k) \chi_\Lambda(|k|)}{2\pi|k|^{\frac{1}{2}}} \frac{\overline{\partial u_\alpha}}{\partial x_j} dx dk \\ &= \sum_{i=1}^3 \int dk \left(\int dx \frac{\partial u_\alpha}{\partial x_i} \frac{1}{k^2 + |k| + h_\alpha + e_0} \frac{\overline{\partial u_\alpha}}{\partial x_i} \right) \sum_{\lambda=1,2} \frac{k^i \epsilon_\lambda^i(k) \chi_\Lambda(|k|)}{2\pi|k|^{\frac{1}{2}}} \varphi(k) dk = 0, \end{aligned}$$

using that the integral over x is independent of the value of i , and since $k \cdot \epsilon_\lambda(k) = 0$.

To prove (211), we note that

$$\begin{aligned} \| |k|^{-\frac{1}{2}} \Phi_{\sharp}^{u_{\alpha}} \|^2 &\leq c\alpha \int \left| \frac{\chi_{\Lambda}(|k|)}{|k|(|k| + k^2 + h_{\alpha} + e_0)} \nabla u_{\alpha} \right|^2 \frac{dx}{k} \\ &\leq c\alpha^3 \int \frac{\chi_{\Lambda}(|k|)^2}{|k|^2(|k| + \frac{3}{16}\alpha^2)^2} \frac{dx}{k} = \mathcal{O}(\alpha). \end{aligned}$$

The proofs of (209) and (210) are similar, but simpler.

We next prove (212).

$$(216) \quad \|P\Phi_{\sharp}^{u_{\alpha}}\|^2 = c\alpha \sum_{i=1}^3 \int \|P \frac{\chi_{\Lambda}(|k|)}{|k|^{\frac{1}{2}}(|k| + k^2 + h_{\alpha} + e_0)} \frac{\partial u_{\alpha}}{\partial x_i}\|_{L^2(\mathbb{R}^3)}^2 \frac{dx}{k}.$$

The function $\partial u_{\alpha}/\partial x_i$ is odd. On the subspace of antisymmetric functions on $L^2(\mathbb{R}^3)$, one has that $-(1 - \gamma_0)\Delta - \frac{\alpha}{|x|} > -e_0$ for some $\gamma_0 > 0$, which implies on this subspace that $h_{\alpha} + e_0 > \gamma_0 P^2$, and thus,

$$(217) \quad P^2 < \gamma_0^{-1}(h_{\alpha} + e_0).$$

The relation (217) yields, for all k

$$\begin{aligned} &\|P \frac{\chi_{\Lambda}(|k|)}{|k|^{\frac{1}{2}}(|k| + k^2 + h_{\alpha} + e_0)} \frac{\partial u_{\alpha}}{\partial x_i}\|^2 \\ (218) \quad &\leq \gamma_0^{-1} \left\| \frac{\chi_{\Lambda}(|k|)}{|k|^{\frac{1}{2}}(|k| + k^2 + h_{\alpha} + e_0)} (h_{\alpha} + e_0)^{\frac{1}{2}} \frac{\partial u_{\alpha}}{\partial x_i} \right\|^2 \\ &\leq \gamma_0^{-1} c\alpha^4 \left(\frac{\chi_{\Lambda}(|k|)}{|k|^{\frac{1}{2}}(|k| + k^2 + \frac{3}{16}\alpha^2)} \right)^2. \end{aligned}$$

Substituting (218) into (216) and integrating over k proves (212).

The proof of (214) is similar. □

Lemma C.5. *We have*

$$\begin{aligned} \|\Phi_{*}^{u_{\alpha}}\|_*^2 - \|\Phi_{\sharp}^{u_{\alpha}}\|_{\sharp}^2 &= \frac{1}{3\pi} \|(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1\|^2 \alpha^5 \log \alpha^{-1} + o(\alpha^5 \log \alpha^{-1}), \\ \|\Phi_{*}^{u_{\alpha}} - \Phi_{\sharp}^{u_{\alpha}}\|_*^2 &= \mathcal{O}(\alpha^5), \\ \|A^{-}(\Phi_{\sharp}^{u_{\alpha}} - \Phi_{*}^{u_{\alpha}})\|^2 &= \mathcal{O}(\alpha^7 \log \alpha^{-1}). \end{aligned}$$

Proof. We have

$$(219) \quad \begin{aligned} &\|\Phi_{*}^{u_{\alpha}}\|_*^2 - \|\Phi_{\sharp}^{u_{\alpha}}\|_{\sharp}^2 \\ &= \frac{\alpha^5}{(2\pi)^2} \frac{4}{3} \left\| \frac{\chi_{\Lambda}(|k|)}{|k|^{\frac{1}{2}}(|k| + k^2)^{\frac{1}{2}}(|k| + k^2 + \alpha^2(h_1 + \frac{1}{4}))^{\frac{1}{2}}} (h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1 \right\|^2, \end{aligned}$$

where the norm here is obviously taken on $L^2(\mathbb{R}^3, \frac{dx}{x}) \otimes L^2(\mathbb{R}^3, \frac{dx}{k})$. Since $(h_1 + \frac{1}{4})\nabla u_1 \in L^2$, it implies that for sufficiently large $c > 0$ independent of α ,

$$\|\chi(h_1 > c)(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1\|^2 < \epsilon,$$

and

$$(220) \quad \left\| \frac{\chi_\Lambda(|k|)}{|k|^{\frac{1}{2}}(|k|+k^2)^{\frac{1}{2}}(|k|+k^2+\alpha^2(h_1+\frac{1}{4}))^{\frac{1}{2}}} \chi(h_1 > c)(h_1+\frac{1}{4})^{\frac{1}{2}} \nabla u_1 \right\|^2 \\ \leq \epsilon \int_0^\infty \frac{\chi_\Lambda^2(t)}{t+c\alpha^2} dt = \epsilon \log \alpha^{-1} + \mathcal{O}(1).$$

For the contribution of $\chi(h_1 \leq c)(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1$ in (219), the following inequalities hold,

$$(221) \quad \left(\frac{1}{\pi} \log \alpha^{-1} + \mathcal{O}(1) \right) \frac{1}{3} \|\chi(h_1 < c)(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1\|^2 \\ = \frac{1}{(2\pi)^2} \frac{1}{3} \left\| \frac{\chi_\Lambda(|k|)}{|k|^{\frac{1}{2}}(|k|+k^2)^{\frac{1}{2}}(|k|+k^2+(c+\frac{1}{4})\alpha^2))^{\frac{1}{2}}} \chi(h_1 < c)(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1 \right\|^2 \\ \leq \frac{1}{(2\pi)^2} \frac{1}{3} \left\| \frac{\chi_\Lambda(|k|)}{|k|^{\frac{1}{2}}(|k|+k^2)^{\frac{1}{2}}(|k|+k^2+(h_1+\frac{1}{4})\alpha^2))^{\frac{1}{2}}} \chi(h_1 < c)(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1 \right\|^2 \\ \leq \frac{1}{(2\pi)^2} \frac{1}{3} \left\| \frac{\chi_\Lambda(|k|)}{|k|^{\frac{1}{2}}(|k|+k^2)^{\frac{1}{2}}(|k|+k^2+\frac{3}{16}\alpha^2))^{\frac{1}{2}}} \chi(h_1 < c)(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1 \right\|^2 \\ \leq \left(\frac{1}{\pi} \log \alpha^{-1} + \mathcal{O}(1) \right) \frac{1}{3} \|\chi(h_1 < c)(h_1 + \frac{1}{4})^{\frac{1}{2}} \nabla u_1\|^2.$$

The inequalities (220) and (221) prove the first equality of the Lemma.

The proofs of the last two equalities are similar but simpler. \square

Lemma C.6.

$$(222) \quad \|(H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha}\|_*^2 - \|\Phi_\#^2\|_*^2 \|\Phi_\#^{u_\alpha}\|^2 = o(\alpha^3 \log \alpha^{-1}),$$

$$(223) \quad \|(H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha}\|^2 = \mathcal{O}(\alpha^3 \log \alpha^{-1})$$

$$(224) \quad \| |k_1|^{\frac{1}{6}} |k_2|^{\frac{1}{6}} |k_3|^{\frac{1}{6}} (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha} \|^2 = \mathcal{O}(\alpha^3)$$

$$(225) \quad \|(h_\alpha + e_0)(H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha}\|^2 = \mathcal{O}(\alpha^7 \log \alpha^{-1}).$$

Proof. Denoting by σ_n the set of all permutations of $\{1, 2, \dots, n\}$, we have

$$\|(H_f + P_f^2)^{-1} A^+ . A^+ \Phi_\#^{u_\alpha}\|_*^2 \\ = \frac{4\alpha}{(2\pi)^3} \left\| \frac{1}{\sqrt{6}} \sum_{(i,j,n) \in \sigma_n} \frac{\sum_{\lambda=1,2} \varepsilon_\lambda(k_i) \cdot \sum_{\nu=1,2} \varepsilon_\nu(k_j)}{\left(\sum_{p=1}^3 |k_p| + (\sum_{p=1}^3 k_p)^2 \right)^{\frac{1}{2}}} \right. \\ \left. \times \frac{\sum_{\kappa=1,2} \varepsilon_\kappa(k_n) \chi_\Lambda(|k_1|) \chi_\Lambda(|k_2|) \chi_\Lambda(|k_3|)}{|k_i|^{\frac{1}{2}} |k_j|^{\frac{1}{2}} |k_n|^{\frac{1}{2}} (|k_n| + k_n^2 + (h_\alpha + e_0))} \nabla u_\alpha \right\|^2$$

If we pick two triples (i, j, n) and (i', j', n') , such that $n \neq n'$, then we get a product which is integrable at $k_1 = k_2 = k_3 = 0$, even without the term $(h_\alpha + e_0)$. The contribution of such terms is $c\alpha \|\nabla u_\alpha\|^2 = \mathcal{O}(\alpha^3)$. Moreover, using the symmetry

in k_1, k_2, k_3 , the twelve remaining terms give the same contribution. This yields

$$\begin{aligned}
 (226) \quad & \| (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_{\sharp}^{u_\alpha} \|_*^2 \\
 &= \frac{8\alpha}{(2\pi)^3} \left\| \frac{\sum_{\lambda} \varepsilon_{\lambda}(k_3) \cdot \sum_{\nu} \varepsilon_{\nu}(k_2) \sum_{\kappa} \varepsilon_{\kappa}(k_1) \chi_{\Lambda}(|k_1|) \chi_{\Lambda}(|k_2|) \chi_{\Lambda}(|k_3|)}{\left(\sum_{i=1}^3 |k_i| + (\sum_{i=1}^3 k_i)^2 \right)^{\frac{1}{2}} |k_3|^{\frac{1}{2}} |k_2|^{\frac{1}{2}} |k_1|^{\frac{1}{2}} (|k_1| + k_1^2 + (h_{\alpha} + e_0))} \cdot \nabla u_{\alpha} \right\|^2 \\
 &+ \mathcal{O}(\alpha^3).
 \end{aligned}$$

On the other hand, we have

$$\begin{aligned}
 & \| \Phi_*^2 \|_*^2 \| \Phi_{\sharp}^{u_\alpha} \|_*^2 \\
 &= \frac{8\alpha}{(2\pi)^3} \left\| \frac{\sum_{\lambda} \varepsilon_{\lambda}(k_3) \cdot \sum_{\nu} \varepsilon_{\nu}(k_2) \sum_{\kappa} \varepsilon_{\kappa}(k_1) \chi_{\Lambda}(|k_1|) \chi_{\Lambda}(|k_2|) \chi_{\Lambda}(|k_3|)}{(|k_2| + |k_3| + (k_2 + k_3)^2)^{\frac{1}{2}} |k_3|^{\frac{1}{2}} |k_2|^{\frac{1}{2}} |k_1|^{\frac{1}{2}} (|k_1| + k_1^2 + (h_{\alpha} + e_0))} \cdot \nabla u_{\alpha} \right\|^2.
 \end{aligned}$$

Therefore, we obtain

$$\begin{aligned}
 (227) \quad & \| (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_{\sharp}^{u_\alpha} \|_*^2 - \| \Phi_*^2 \|_*^2 \| \Phi_{\sharp}^{u_\alpha} \|_*^2 \\
 &= -\frac{8\alpha}{(2\pi)^3} \int \left(\frac{(|k_1|^2 + 2|k_1| |k_3| + 2|k_1| |k_2|)}{|k_2| |k_3| (|k_1| + |k_2| + |k_3| + |k_1 + k_2 + k_3|^2) (|k_2| + |k_3| + |k_2 + k_3|^2)} \right. \\
 &\times \sum_{\lambda} \varepsilon_{\lambda}(k_3) \cdot \sum_{\nu} \varepsilon_{\nu}(k_2) \chi_{\Lambda}(|k_1|) \chi_{\Lambda}(|k_2|)^2 \chi_{\Lambda}(|k_3|)^2 |u(k_1, x)|^2 k_1 k_2 k_3 x \Big),
 \end{aligned}$$

where

$$u(k_1, x) = \frac{\sum_{\kappa} \varepsilon_{\kappa}(k_1) \chi_{\Lambda}(|k_1|)^{\frac{1}{2}}}{|k_1|^{\frac{1}{2}} (|k_1| + k_1^2 + (h_{\alpha} + e_0))} \cdot \nabla u_{\alpha}.$$

For fixed δ , we first compute the integral in (227) over the regions $|k_1| > \delta$. This yields a term $c_{\delta} \alpha^3$, where c_{δ} is independent on α .

Next, integrating (227) over the regions $|k_1| \leq \delta$ yields a bound $\mathcal{O}(\delta) \alpha^3 \log \alpha$, with $\mathcal{O}(\delta)$ independent of α .

This concludes the proof of (222).

The proof of (223) is a straightforward computation showing

$$\begin{aligned}
 & \| (H_f + P_f^2)^{-1} A^+ . A^+ \Phi_{\sharp}^{u_\alpha} \|_*^2 \\
 &= \frac{8\alpha}{(2\pi)^3} \left\| \frac{\sum_{\lambda} \varepsilon_{\lambda}(k_3) \cdot \sum_{\nu} \varepsilon_{\nu}(k_2) \sum_{\kappa} \varepsilon_{\kappa}(k_1) \chi_{\Lambda}(|k_1|) \chi_{\Lambda}(|k_2|) \chi_{\Lambda}(|k_3|)}{\left(\sum_{i=1}^3 |k_i| + (\sum_{i=1}^3 k_i)^2 \right) |k_3|^{\frac{1}{2}} |k_2|^{\frac{1}{2}} |k_1|^{\frac{1}{2}} (|k_1| + k_1^2 + (h_{\alpha} + e_0))} \cdot \nabla u_{\alpha} \right\|^2 \\
 &= \mathcal{O}(\alpha^3 \log \alpha^{-1}).
 \end{aligned}$$

The proofs of (224) and (225) are similar to the above. \square

Lemma C.7. For Ψ_1^{\perp} and Ψ_2^{\perp} given in Definition 6.1, we have

$$(228) \quad |2\alpha \operatorname{Re} \langle A^- . A^- \pi^3 \Psi_1^{\perp}, \pi^1 \Psi_2^{\perp} \rangle| \leq c\alpha^5 + \epsilon \|H_f^{\frac{1}{2}} \Psi_2^{\perp}\|^2.$$

Proof. Using $2\alpha \operatorname{Re} \langle A^- . A^- \pi^3 \Psi_1^{\perp}, \pi^1 \Psi_2^{\perp} \rangle = 2\alpha \operatorname{Re} \langle \pi^3 \Psi_1^{\perp}, A^+ . A^+ \pi^1 \Psi_2^{\perp} \rangle$, we obtain that $2\alpha \operatorname{Re} \langle A^- . A^- \pi^3 \Psi_1^{\perp}, \pi^1 \Psi_2^{\perp} \rangle$ can be rewritten as a linear combination of the

following two integrals I_1 and I_2

$$(229) \quad \begin{aligned} I_1 &= \alpha^{\frac{5}{2}} \int \frac{\mathbb{k} \mathbb{k}' \mathbb{k}'' \mathbb{x}}{|k|+|k'|+|k''|+|k+k'+k''|^2} \sum_{\kappa, \nu} \frac{\epsilon_\kappa(k') \cdot \epsilon_\nu(k'')}{|k'|^{\frac{1}{2}} |k''|^{\frac{1}{2}}} \\ &\times \left(\frac{1}{|k|+|k|^2+(h_\alpha+e_0)} \sum_{\lambda} \frac{\epsilon_\lambda(k)}{|k|^{\frac{1}{2}}} \nabla u_\alpha \right) \sum_{\kappa, \eta} \frac{\epsilon_\kappa(k') \cdot \epsilon_\eta(k'')}{|k'|^{\frac{1}{2}} |k''|^{\frac{1}{2}}} \overline{(\pi^1 \Psi_2^\perp)(k, x)}, \end{aligned}$$

and I_2 is defined as I_1 , except that in the last sum, we reverse the role of k and k'' , namely

$$(230) \quad \begin{aligned} I_2 &= \alpha^{\frac{5}{2}} \int \frac{\mathbb{k} \mathbb{k}' \mathbb{k}'' \mathbb{x} \chi_\Lambda(|k|) \chi_\Lambda(|k'|) \chi_\Lambda(|k''|)}{|k|+|k'|+|k''|+|k+k'+k''|^2} \sum_{\kappa, \nu} \frac{\epsilon_\kappa(k') \cdot \epsilon_\nu(k'')}{|k'|^{\frac{1}{2}} |k''|^{\frac{1}{2}}} \\ &\times \left(\frac{1}{|k|+|k|^2+(h_\alpha+e_0)} \sum_{\lambda} \frac{\epsilon_\lambda(k)}{|k|^{\frac{1}{2}}} \nabla u_\alpha \right) \sum_{\kappa, \eta} \frac{\epsilon_\kappa(k') \cdot \epsilon_\eta(k'')}{|k'|^{\frac{1}{2}} |k|^{\frac{1}{2}}} \overline{(\pi^1 \Psi_2^\perp)(k'', x)}. \end{aligned}$$

To bound I_2 , we use the Schwarz inequality, $|ab| \leq \frac{1}{2\epsilon} a^2 + \frac{\epsilon}{2} b^2$. This yields

$$\begin{aligned} |I_2| &\leq \alpha^{\frac{5}{2}} \left(\int \frac{\mathbb{k} \mathbb{k}' \mathbb{k}'' \mathbb{x} \chi_\Lambda^2(|k|) \chi_\Lambda^2(|k'|) \chi_\Lambda^2(|k''|)}{(|k|+|k'|+|k''|+|k+k'+k''|^2)^2} \frac{2}{|k'| |k''|} \right. \\ &\times \left. \left| \frac{1}{|k|+|k|^2+(h_\alpha+e_0)} \frac{1}{|k|^{\frac{1}{2}}} \nabla u_\alpha \right|^2 |k|^{2-\epsilon} |k'|^{2-\epsilon} \frac{1}{|k''|} \right)^{\frac{1}{2}} \\ &\times \left(2 \int \frac{\mathbb{k} \mathbb{k}' \mathbb{k}'' \mathbb{x}}{|k'| |k''|} |(\pi^1 \Psi_2^\perp)(k'', x)|^2 |k''| \frac{\chi_\Lambda^2(|k|) \chi_\Lambda^2(|k'|) \chi_\Lambda^2(|k''|)}{|k|^{2-\epsilon} |k'|^{2-\epsilon}} \right)^{\frac{1}{2}} \\ &\leq c\alpha^7 + \epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2, \end{aligned}$$

Similarly, we bound I_1 as follows,

$$\begin{aligned} |I_1| &\leq \alpha^{\frac{5}{2}} \left(\int \frac{\mathbb{k} \mathbb{k}' \mathbb{k}'' \mathbb{x}}{(|k|+|k'|+|k''|+|k+k'+k''|^2)^2} \frac{2}{|k'| |k''|} \right. \\ &\times \left. \left| \frac{1}{|k|+|k|^2+(h_\alpha+e_0)} \frac{1}{|k|^{\frac{1}{2}}} \nabla u_\alpha \right|^2 |k'|^{2-\epsilon} |k''|^{2-\epsilon} \frac{\chi_\Lambda^2(|k|) \chi_\Lambda^2(|k'|) \chi_\Lambda^2(|k''|)}{|k|} \right)^{\frac{1}{2}} \\ &\times \left(2 \int \frac{\mathbb{k} \mathbb{k}' \mathbb{k}'' \mathbb{x}}{|k'| |k''|} |(\pi^1 \Psi_2^\perp)(k, x)|^2 |k| \frac{\chi_\Lambda^2(|k|) \chi_\Lambda^2(|k'|) \chi_\Lambda^2(|k''|)}{|k'|^{2-\epsilon} |k''|^{2-\epsilon}} \right)^{\frac{1}{2}} \\ &\leq c\alpha^5 \int \mathbb{k} \mathbb{k}'' |k'|^{-\epsilon} |k''|^{-\epsilon} \chi_\Lambda^2(|k'|) \chi_\Lambda^2(|k''|) \int \mathbb{k} \frac{1}{|k|^2} \frac{\chi_\Lambda^2(|k|)}{(|k|+\frac{1}{16}\alpha^2)^2} \|\nabla u_\alpha\|^2 \\ &+ \epsilon \|H_f^{\frac{1}{2}} \pi^1 \Psi_2^\perp\|^2 \leq c\alpha^5 + \epsilon \|H_f^{\frac{1}{2}} \Psi_2^\perp\|^2, \end{aligned}$$

where we took into account that $\frac{\partial u_\alpha}{\partial x_i}$ is orthogonal to u_α , and on the subspace of such functions we have $(h - \alpha + e_0) \geq \frac{1}{16}\alpha^2$. \square

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